

3. Particle accelerators

3.1 Relativistic particles

3.2 Electrostatic accelerators

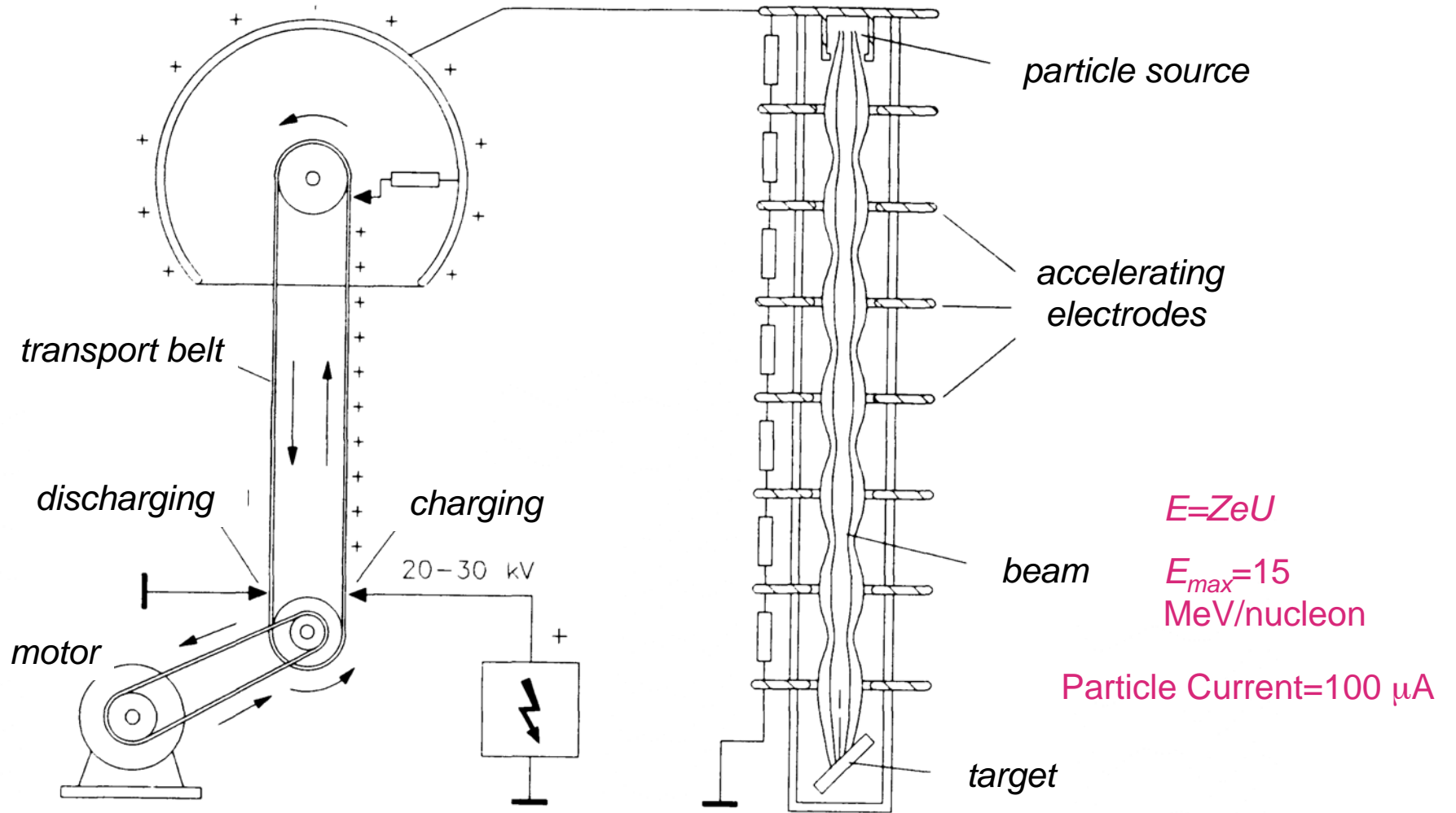
3.3 Ring accelerators

Betatron // Cyclotron // Synchrotron

3.4 Linear accelerators

3.5 Collider

Van-de-Graaf accelerator



Tandem-accelerator

(Povh et al., N & P)

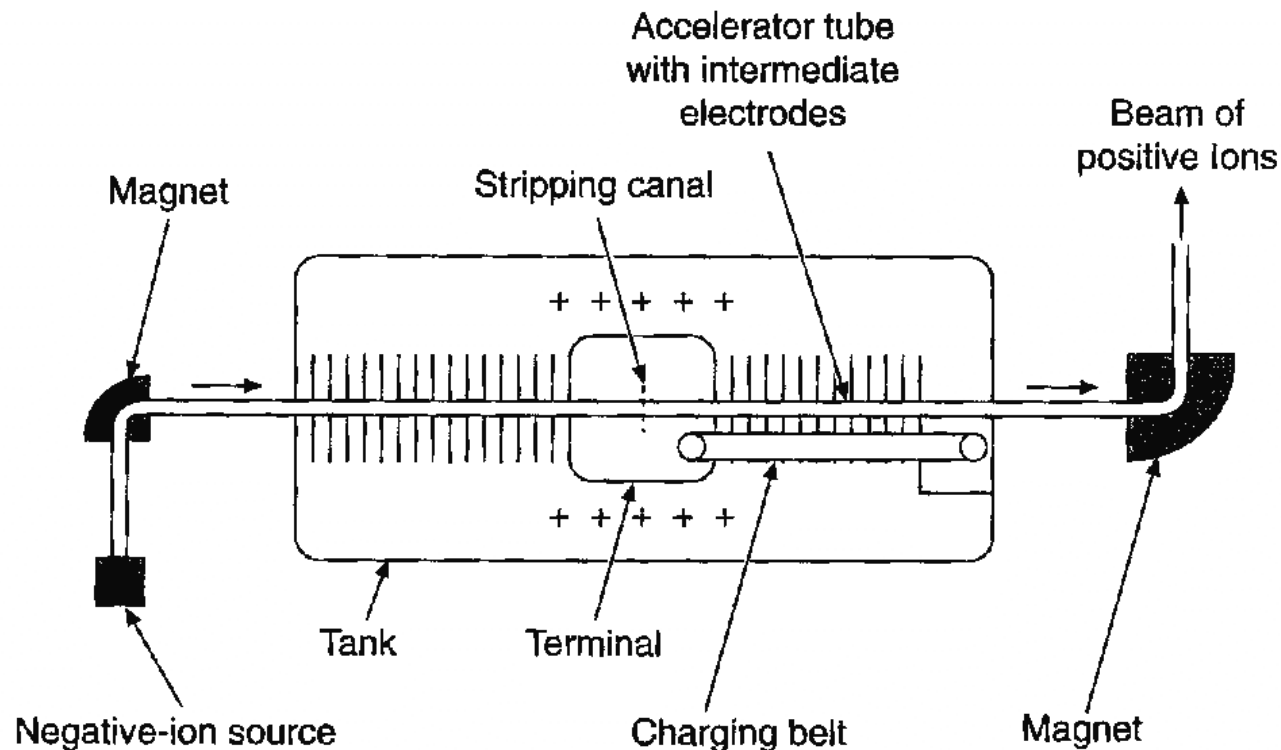


Fig. A.1. Sketch of a tandem Van de Graaff accelerator. Negative ions are accelerated from the left towards the terminal where some of their electrons are stripped off and they become positively charged. This causes them to now be accelerated away from the terminal and the potential difference between the terminal and the tank is traversed for a second time.

RING ACCELERATORS

1) The Betatron

$$q\mathbf{v} \times \mathbf{B}_0 = -m\mathbf{v} \times \boldsymbol{\omega}$$

$$\omega_c \text{ Cyclotron Frequency such that: } q\mathbf{B}_0 = -m\boldsymbol{\omega}_c = -m \mathbf{v}/r = -|\mathbf{p}|/r$$

If we have a time dependent \mathbf{B} field, this induces an electric field that can be used for acceleration.

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}, \quad \text{Stoke Theorem:}$$

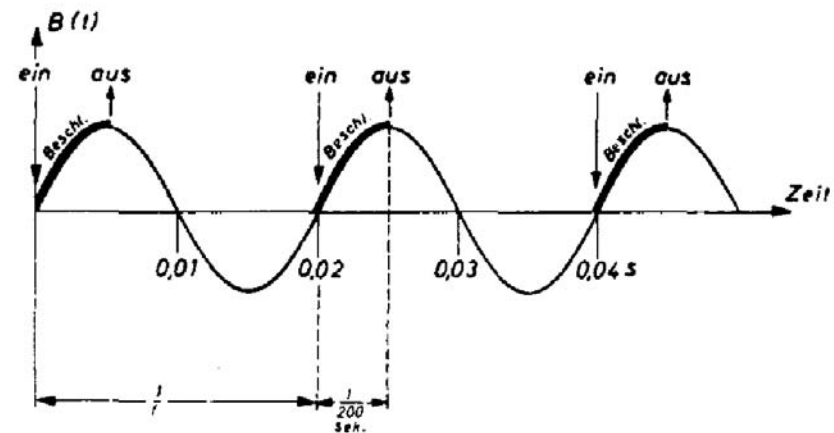
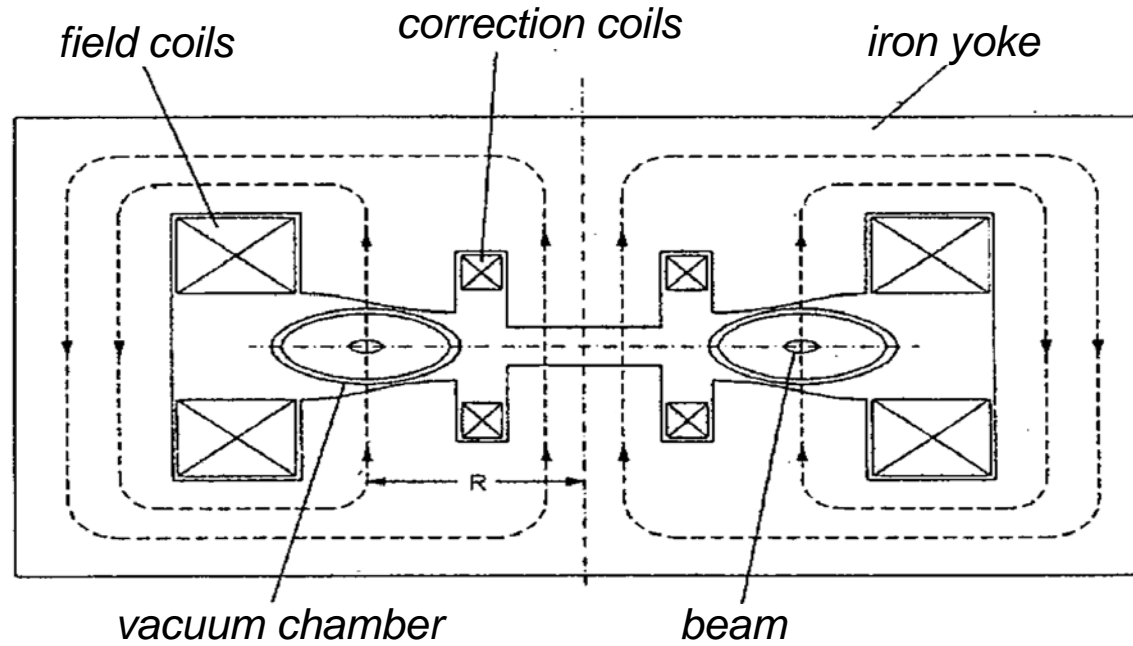
$$\Rightarrow 2\pi r E = -\frac{d}{dt}(\bar{B}\pi r^2), \quad r = r_0 \rightarrow F = \frac{d|p|}{dt} = qE = -\frac{qr_0}{2} \frac{d\bar{B}}{dt}$$

$$\Rightarrow \Delta|p| = -\frac{qr_0}{2} \Delta\bar{B} \quad \text{and} \quad \Delta|p| = -qr_0 \Delta B_0$$

$$\Rightarrow \Delta\bar{B} = 2\Delta B_0$$

The particle can be accelerated only once until the field reaches its maximum value B_0

How it looks like:



Axial stability

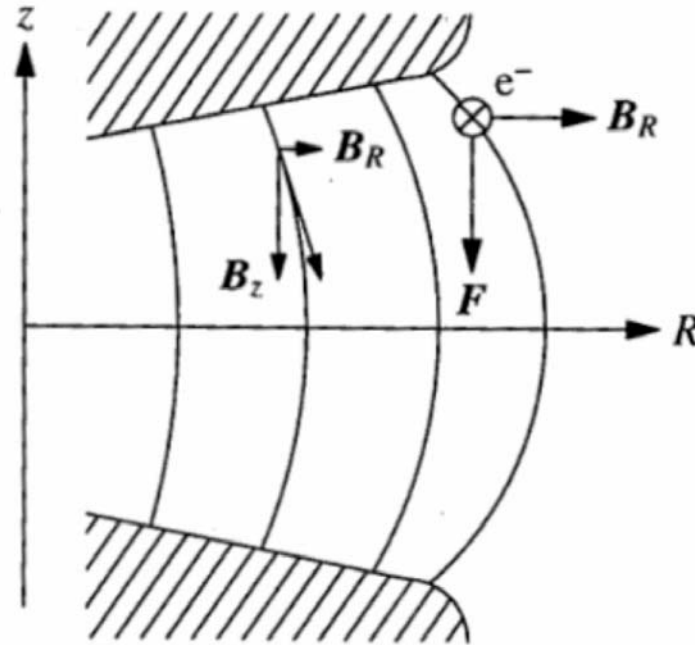


Abb. 2.10: Axiale Stabilitätsbedingung

The magnetic field forces the particle back to the medium plane.

The restoring force is provided by the magnetic field gradient.

$E_{\max} = 300 \text{ MeV}$ for e^-

2) Cyclotron

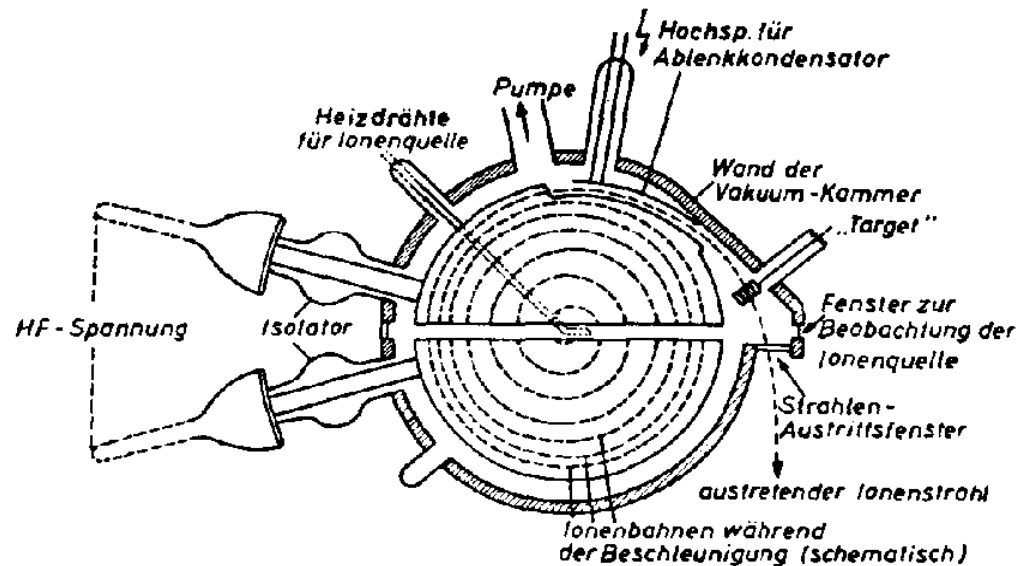
Constant Magnetic Field, Acceleration happens via a oscillating electric field between the dees
 $\omega_{HF} = |\omega_z|$ angular velocity of the particle (10 MHz)

$$m \cdot \frac{v^2}{r} = q \cdot (v \times B) \rightarrow \frac{v}{r} = \omega_z = \frac{q}{m} B$$

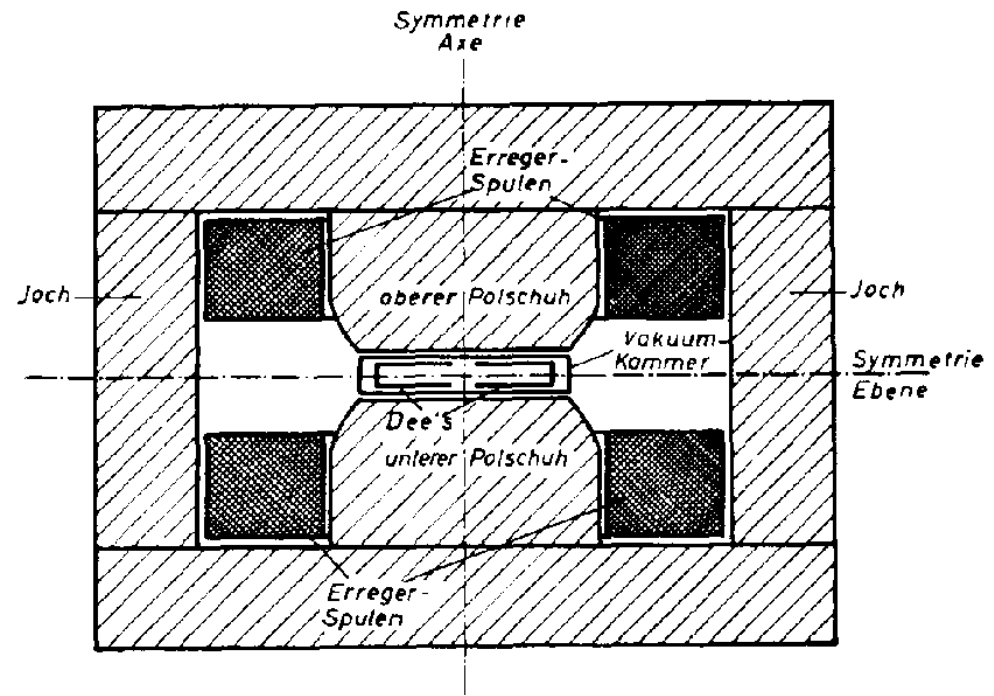
- Independent of the radius!!!
- Maximal Energy does not depend on E!!

$E_{\text{MAX proton}} = 20 \text{ MeV}$

Schematic from
the top



Side View



When the particles become relativistic $m \rightarrow m\gamma$ with

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}, \quad \frac{v}{c} = \beta$$

Hence the particle becomes heavier and the ω_c diminishes.

One can overcome this problem reducing the HF frequency while the particle travels (**Synchrocyclotron**, only possible in bunch mode) or one can increase the magnetic field such that the radius stays constant

(**Isocyclotron**, possible in continuous mode)

$$\omega_z = \frac{q}{m(E)} B(r(E))$$

Not suited for the acceleration of electrons!

One of the first Cyclotrons...

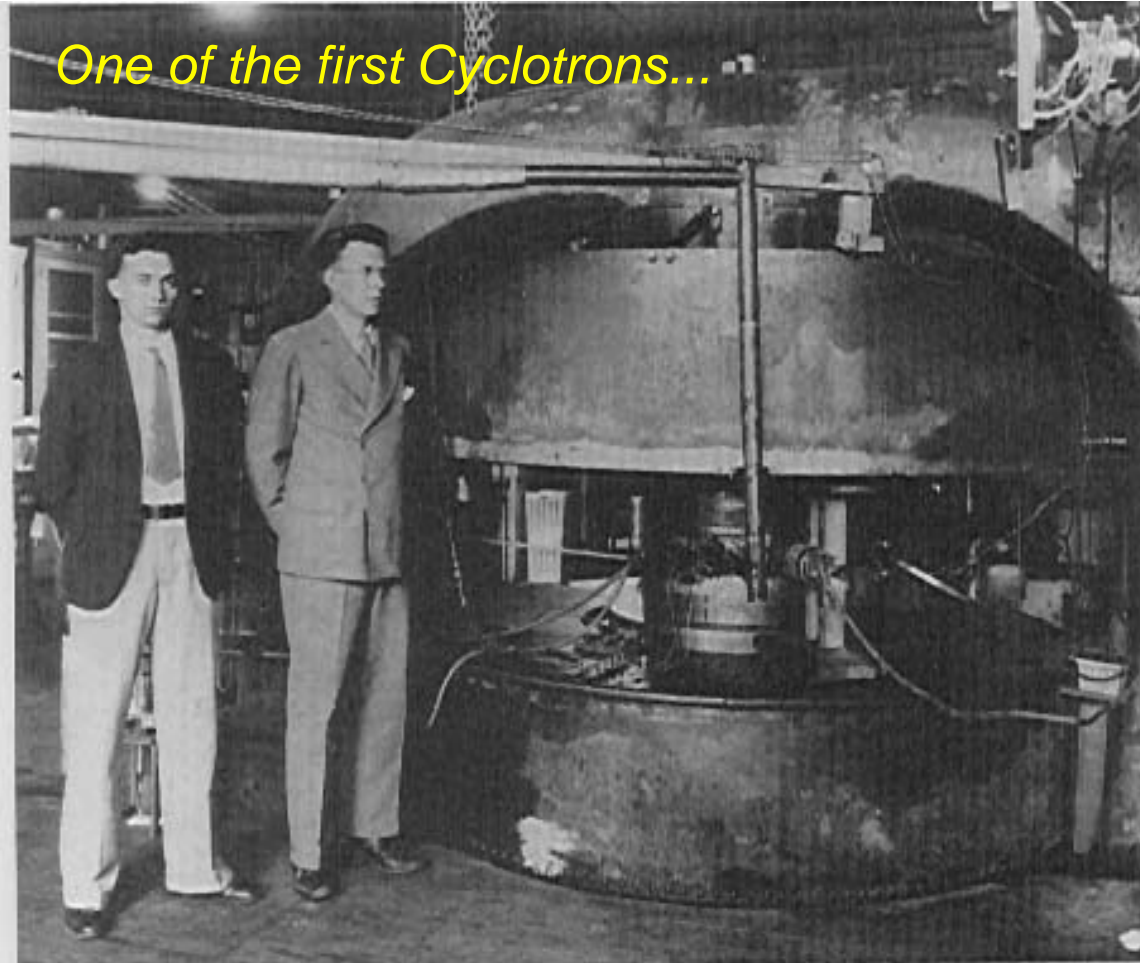


Abb. 11.7 M. S. Livingston und E. O. Lawrence im ersten Strahlenlabor der *University of California* in Berkeley neben dem 37-inch-Zyklotron. Ursprünglich maß das Zyklotron 27 inch, es wurde aber auf 37 inch vergrößert und zur Messung des magnetischen Moments von Neutronen sowie zur Herstellung des ersten künstlichen Elements, Technetium, eingesetzt, (Lawrence Berkeley Laboratory)

... and a little bit later



Abb. 11.9 Das nach dem Krieg gebaute 184-inch-Synchrozyklotron mit einem Teil der an seinem Bau beteiligten Belegschaft. Mit dieser Maschine wurden die ersten künstlichen Mesonen erzeugt. (Lawrence Berkeley Laboratory)

The isochrone- cyclotron at PSI



Synchrotron

For relativistic particles ($v \approx c$):

$$\frac{v}{r} = \frac{q}{m} B \rightarrow r = \frac{mv}{qB} = \frac{p}{qB} \approx \frac{E}{qcB}$$

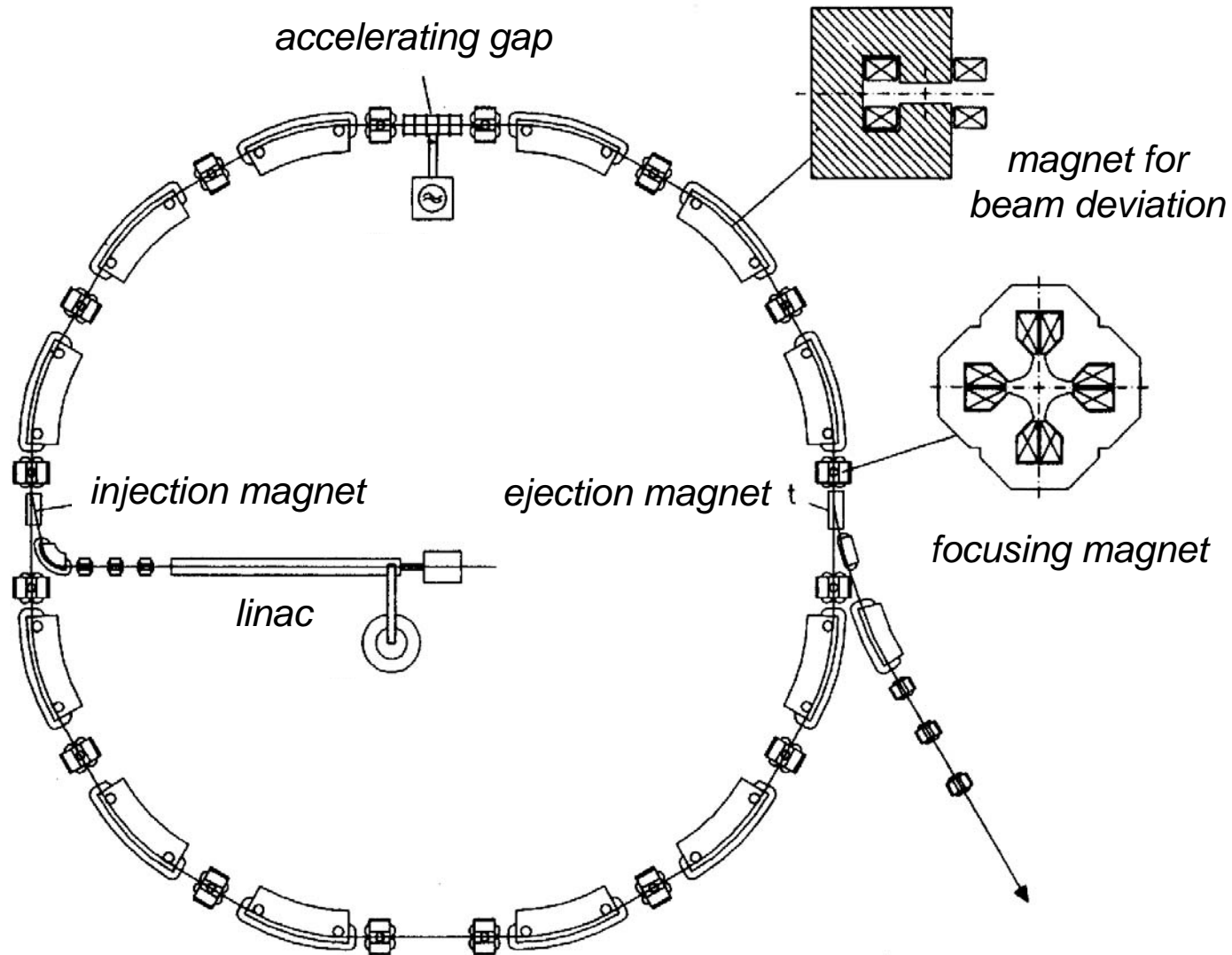
$$p = eBr \text{ bzw. } p(\text{GeV}/c) = 0.3 \cdot B(\text{T}) \cdot r(\text{m})$$

The orbit radius increases with the Energy and this can be compensated only by higher magnetic fields. Maximal $B = 5-10 \text{ T}!!!!$ Moreover big jokes are very expensive.

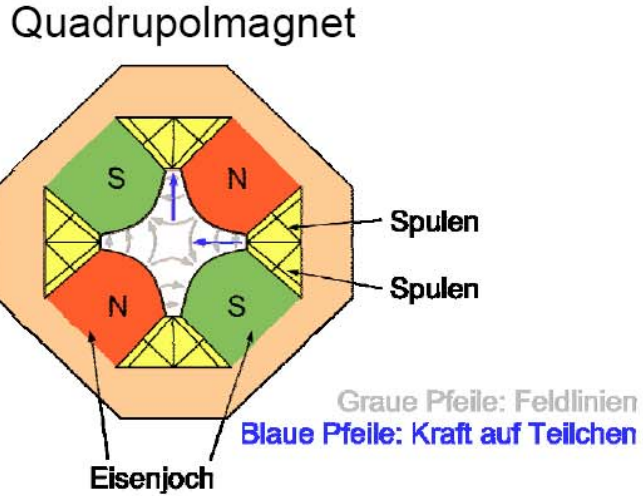
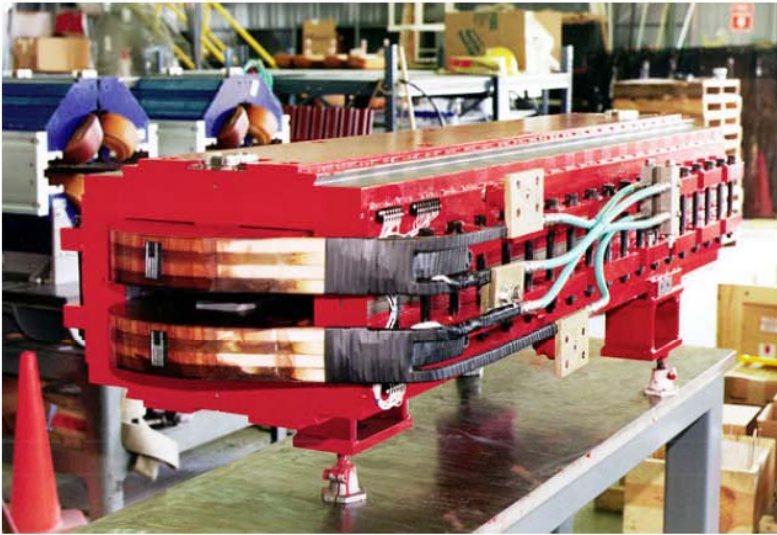
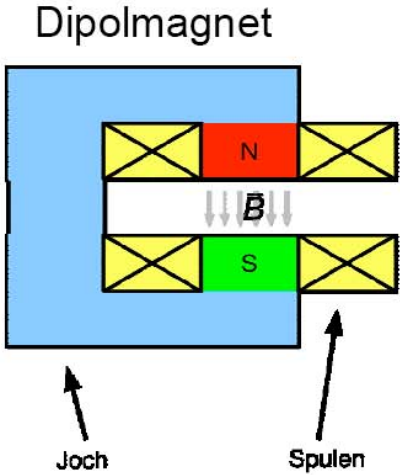
The new idea is to keep the orbit constant and oblige the particle to run along the circle via dipole magnets. Along the path there are different acceleration gaps such that E/B stays constant. This means that the magnetic field has to be risen synchronic to the E field.

Synchrotron

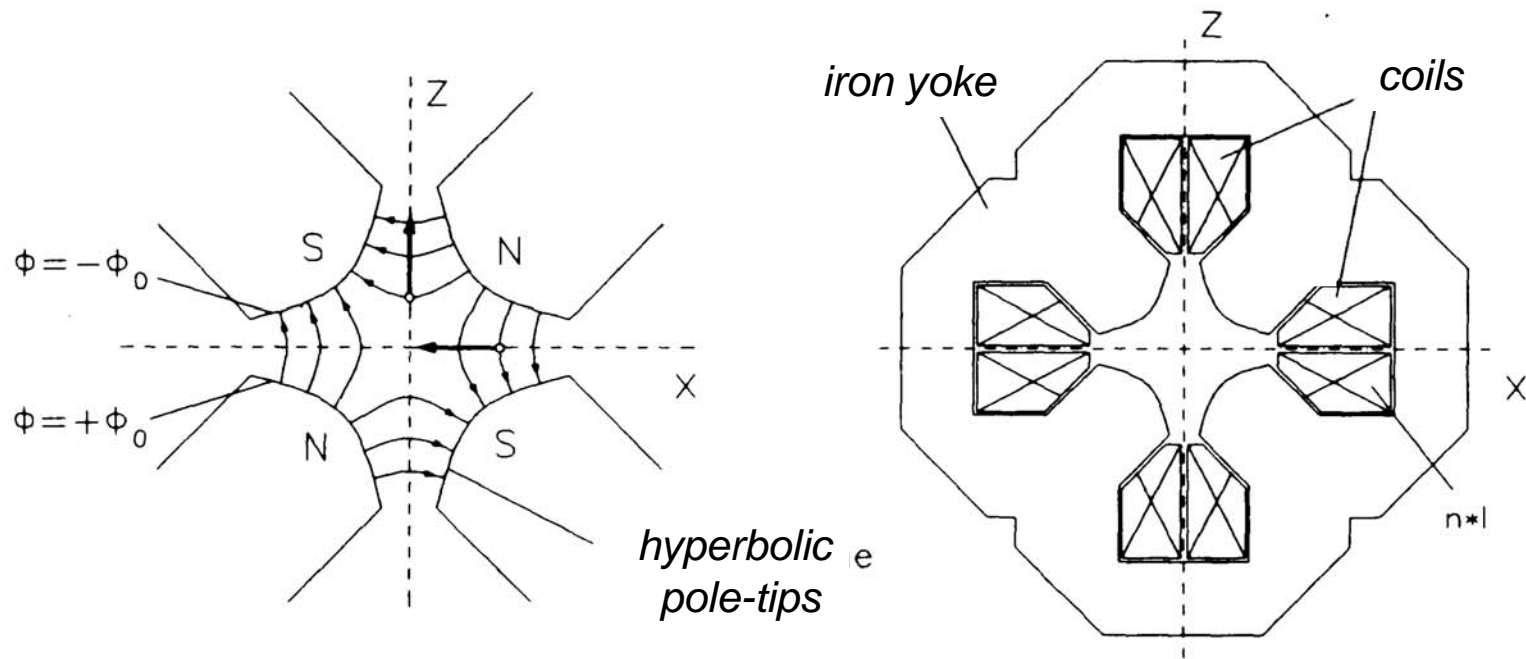
(Wille, Teilchenbeschleuniger)



Dipole and Quadropole



Focusing



Since the quadrupole is focusing the beam in one direction and defocusing in the other, there are placed couple-wise after one another and turned of 90 deg.

Synchrotron Radiation

Since the particles are accelerated on a circular orbit, they radiate energy. For each circle we have the following energy loss:

$$\Delta E_{synchr} = \frac{e^2}{3\epsilon_0(m_0c^2)} \frac{E^4}{R}$$

Such energy loss is 10^{13} times larger for electrons than for protons.

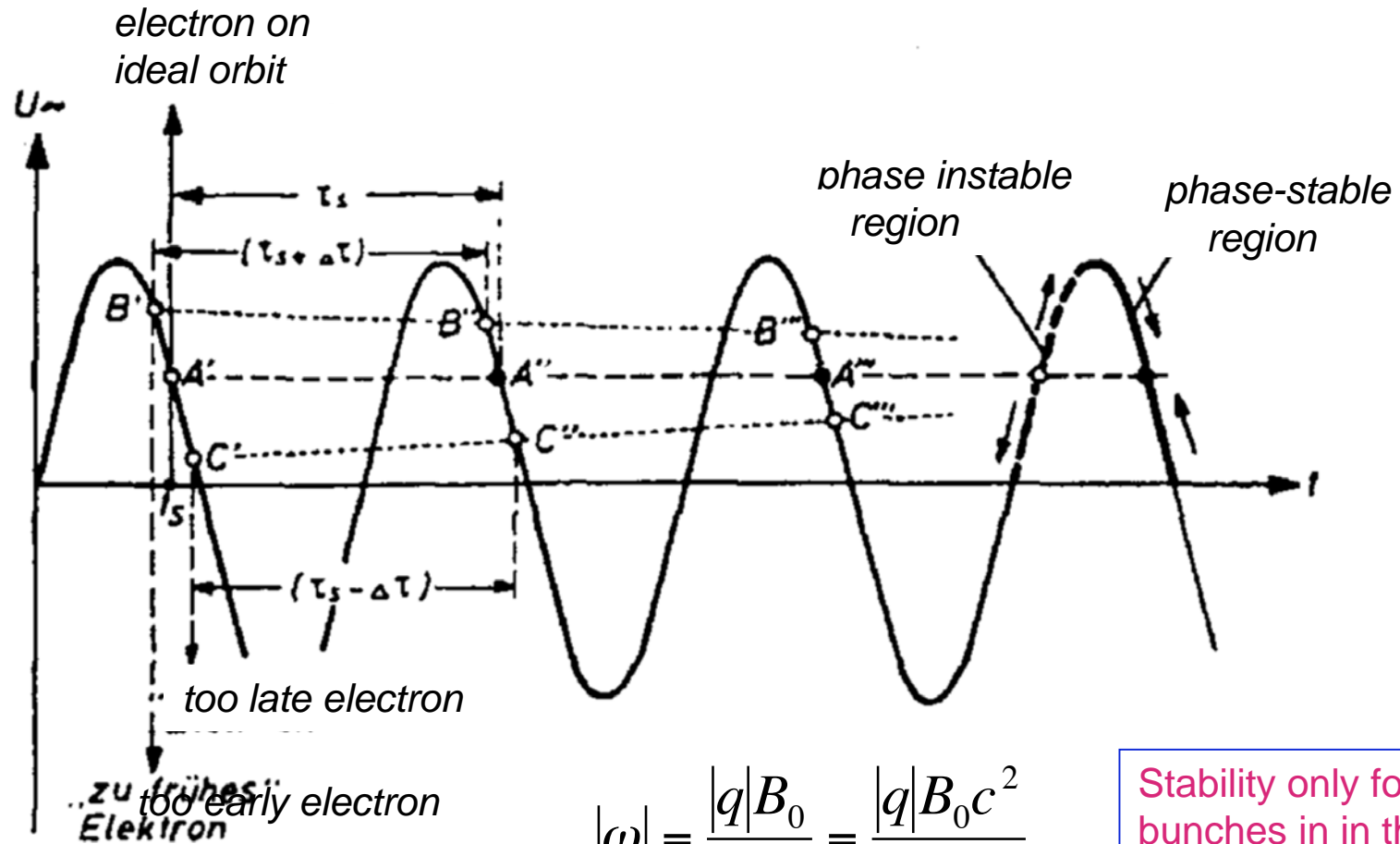
Despite of the fact that large radii can reduce such loss this implies a maximal reachable energy for electron of 100 GeV

$$\Delta E_{synch,e}(keV) = 88,5 \cdot \frac{E^4(GeV^4)}{R(m)}$$

$$\Delta E_{synch,p}(eV) = 7,79 \cdot 10^{-9} \frac{E^4(GeV^4)}{R(m)}$$

The limit in the acceleration of the protons is given by the steering magnets. Furthermore particles have to be pre-accelerated before entering the synchrotron, since the magnets cannot deflect particles with energy close to 0.

Phase Diagram of the Synchrotron



$$|\omega| = \frac{|q|B_0}{\gamma m} = \frac{|q|B_0 c^2}{E}$$

Stability only for bunches in in the orbit!



Lear (CERN)

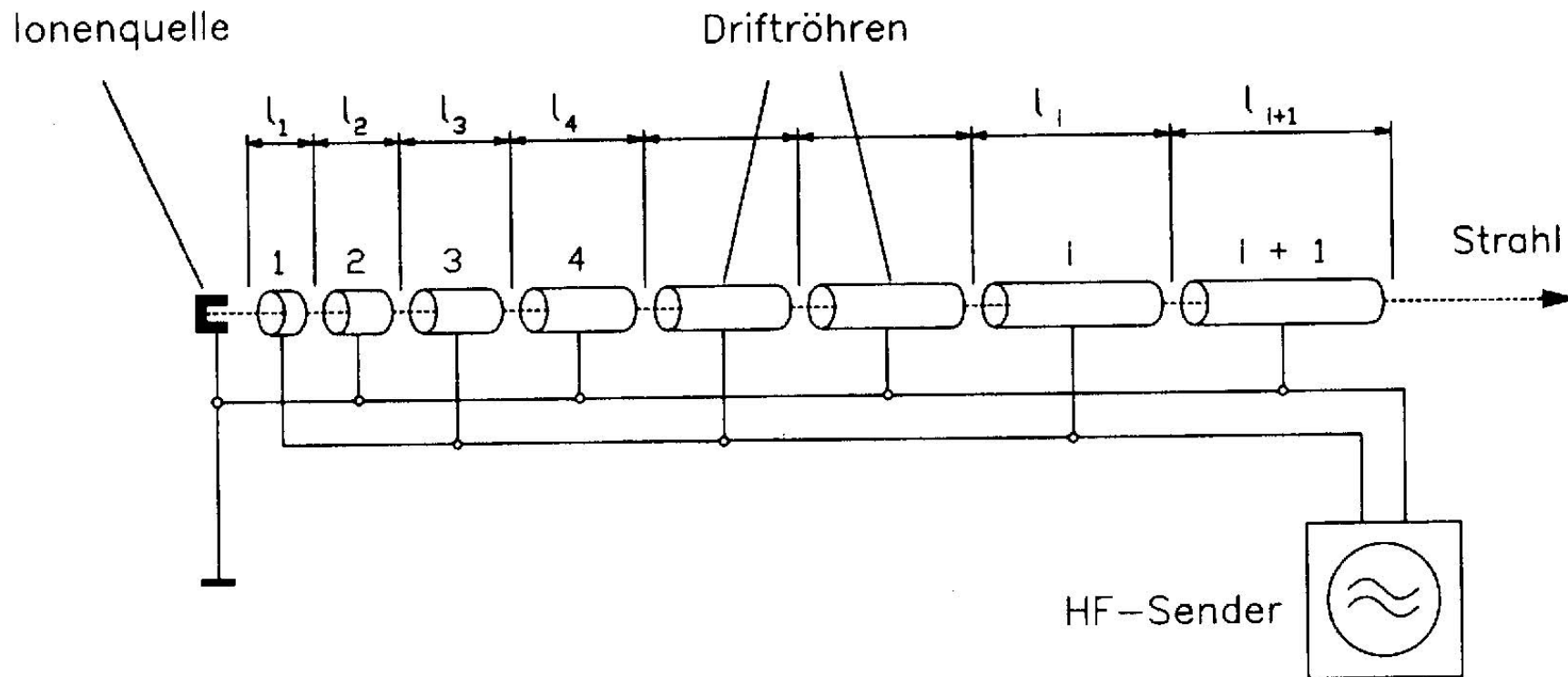
Proton-Linac

(Wille, Teilchenbeschleuniger)

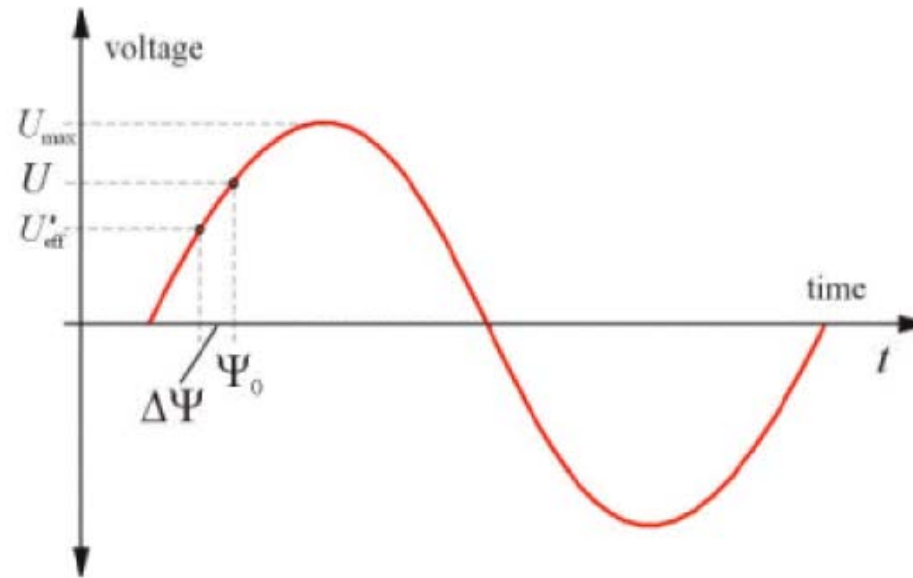
Good also for electron acceleration!

$E_{\text{MAX proton}} = 100 \text{ MeV}$, used as injector for ring accelerators

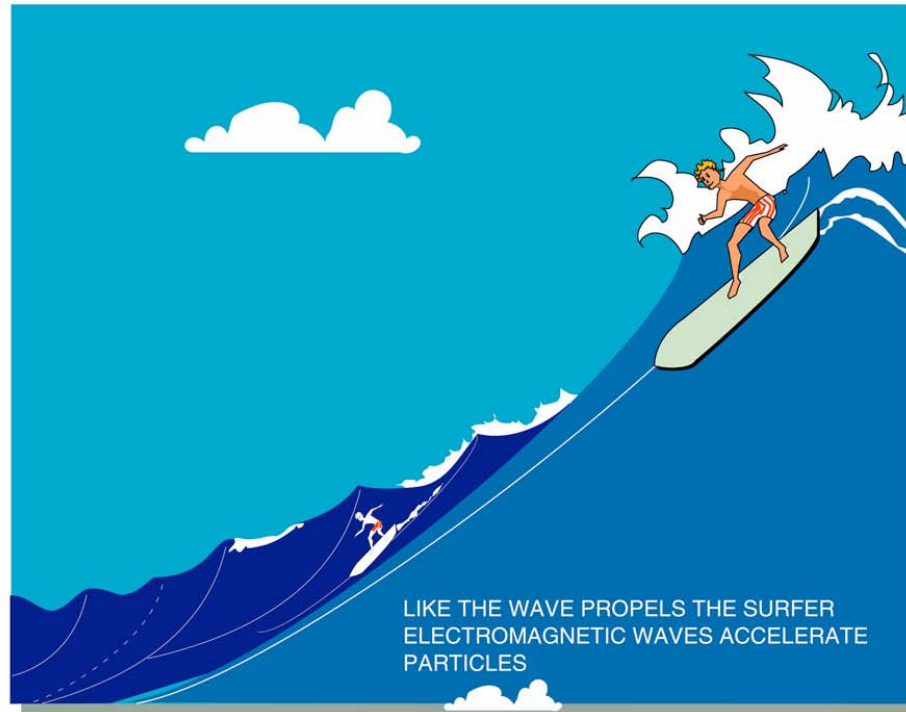
TESLA: 30 km electron LINAC for 500 GeV electrons



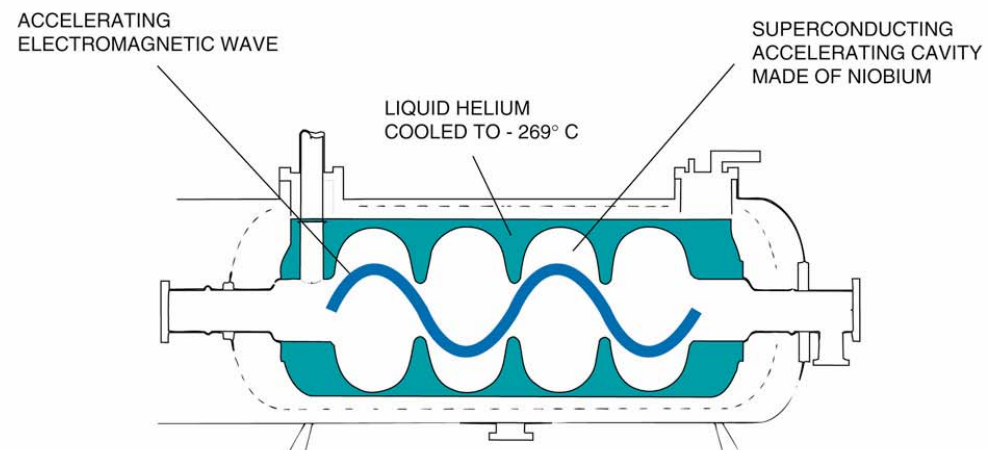
Phase Focusing



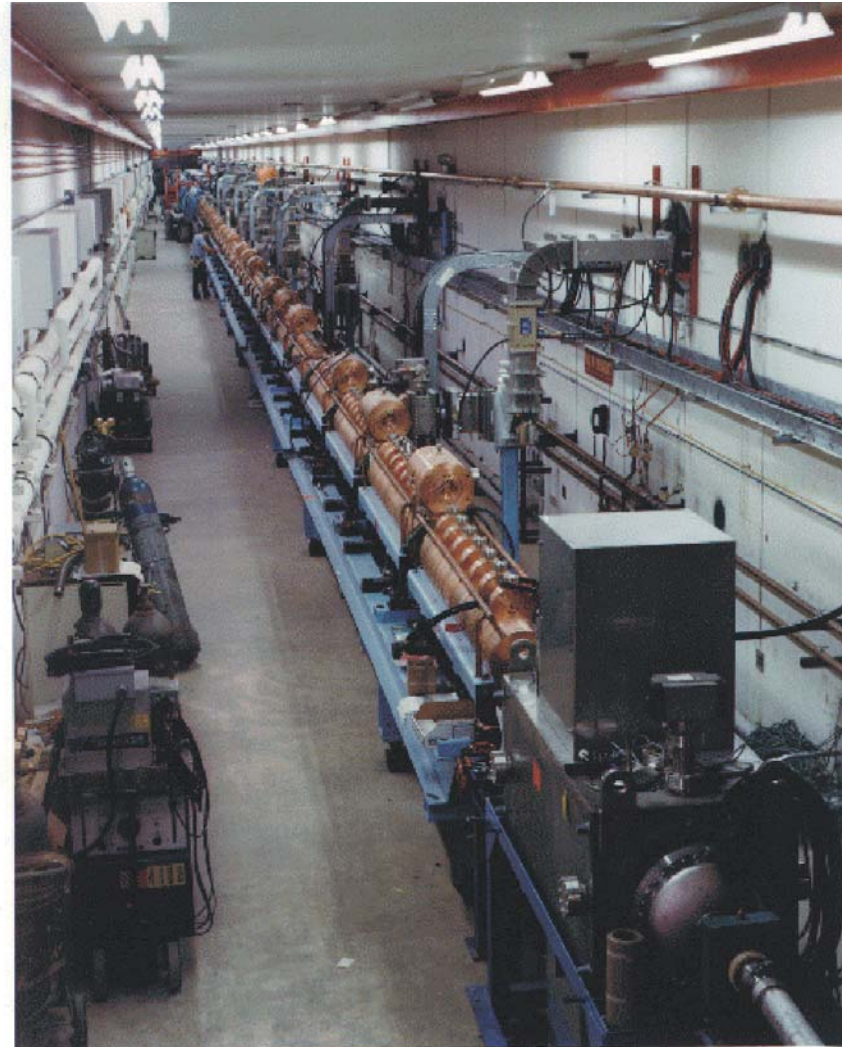
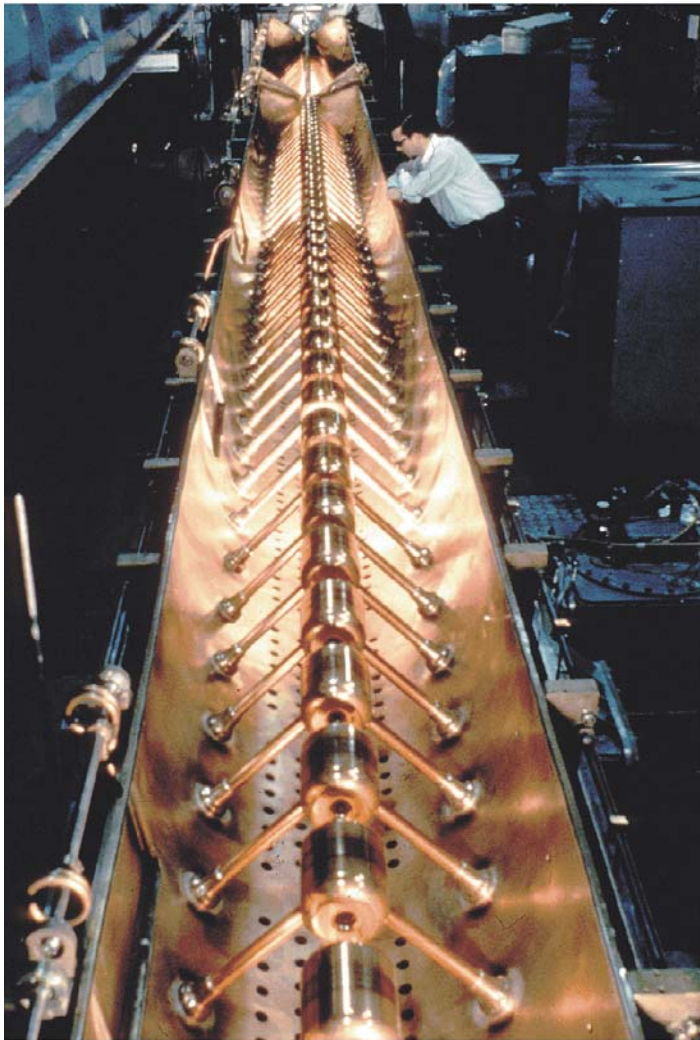
A particle that is faster and arrives earlier sees a smaller V and hence will be slowed down in the next cycle. This is again only possible for a BUNCHED beam.



THE USE OF SUPRACONDUCTIVITY TO INCREASE
PERFORMANCES AND CONSIDERABLY REDUCE
ELECTRICITY CONSUMPTION



Linacs at CERN

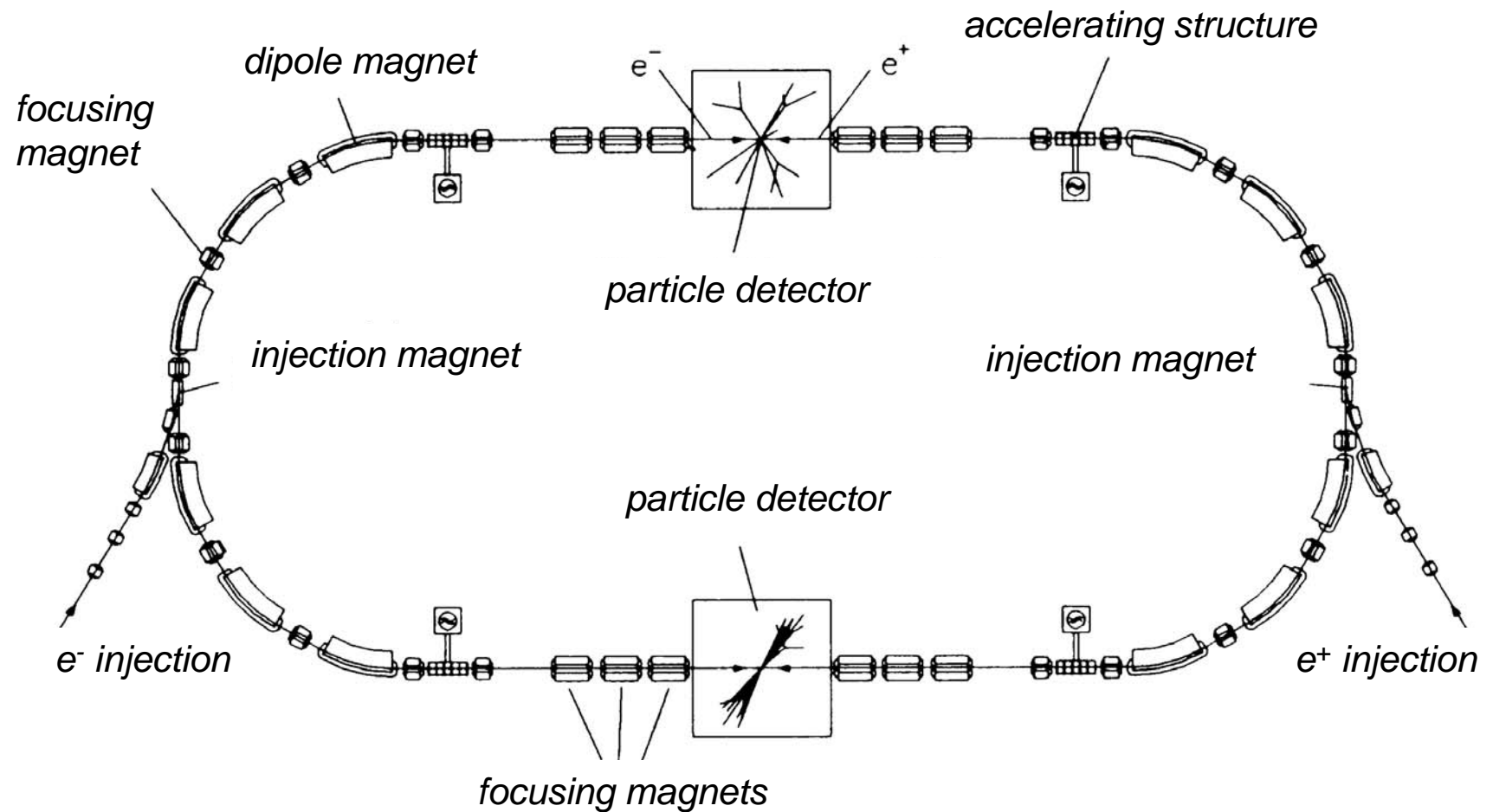


Largest LINAC at SLAC (Stanford Linear Accelerator Center)

$L=3\text{km}$, $E_{\text{MAX electron}}=50\text{ GeV}$

Collider

(Wille, Teilchenbeschleuniger)



1. Energy:

in the cm system in terms of momentum 4-vectors $s = (p_1+p_2)^2$

fixed target: $s = (m_1c^2)^2 + 2\gamma_1m_1c^2m_2c^2 + (m_2c^2)^2$

special case of equal masses $\sqrt{s} = mc^2 \sqrt{2+2\gamma}$

high energy limit $\sqrt{s} = mc^2 \sqrt{2\gamma}$

colliding beams: $s = (m_1c^2)^2 + (m_2c^2)^2 + 2\gamma_1\gamma_2m_1c^2m_2c^2(1 + \beta_1\beta_2)$

high energy limit

$$s = 4E_1E_2$$

special case of equal mass and energy

$$\sqrt{s} = 2 E = mc^2 2\gamma$$

note: linear in γ

Luminosity

2. Luminosity:

fixed target: $\mathcal{L} = N_b [1/s] N_t [1/cm^2]$ beam rate times target thickness

$$N_t = \rho t N_A / M \quad \text{e.g. for 1m liquid hydrogen } N_t = 2 \cdot 10^{24}/cm^2$$

$$\text{typical for protons } N_b = 10^{13}/s \rightarrow \mathcal{L} = 2 \cdot 10^{37}/cm^2s$$

colliding beams: $\mathcal{L} = f n N_1 N_2 / A$

f frequency

n number of bunches in either beam around ring

$N_{1,2}$ particles per bunch

A cross sectional area of beam

$$\text{typical } e^+e^- \text{ collider } \mathcal{L} = 10^{31}/cm^2s$$

$$\text{ppbar collider } \mathcal{L} = 10^{30}/cm^2s$$

$$\text{pp collider } \mathcal{L} = 10^{33}/cm^2s$$

Large Electron-Positron Collider (LEP) am CERN

Betrieben von 1989-2000

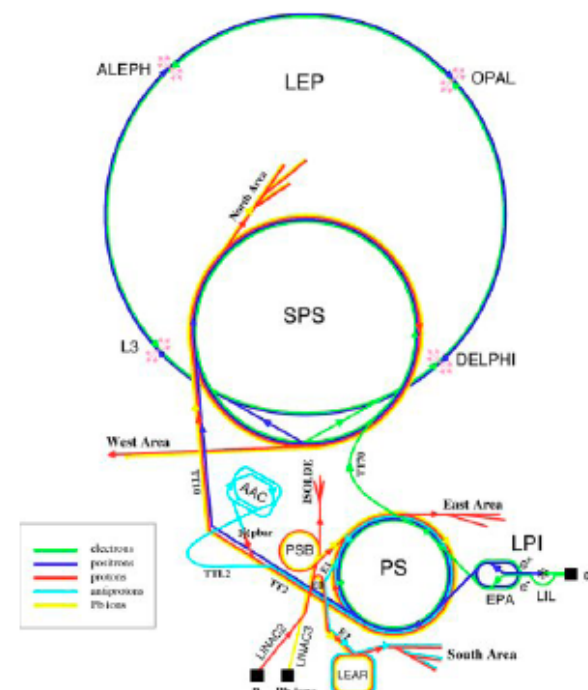
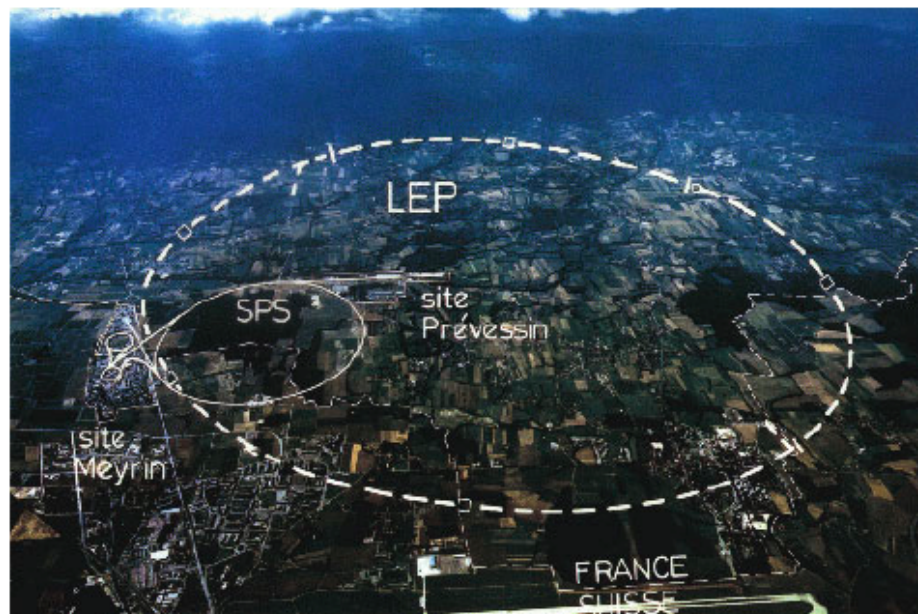
Maximalenergie 100 GeV $\rightarrow \sqrt{s} = 200$ GeV

Umfang 26.7 km, zwischen 40 und 150 m unter der Erde, 1,4% Neigung

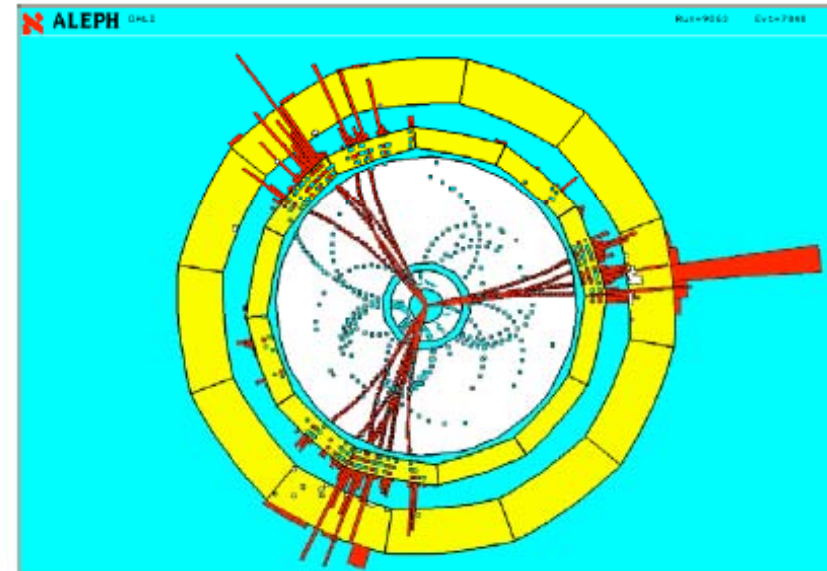
3368 Magnete

272 Beschleunigerkavitäten

4 Kollisionspunkte mit Experimenten

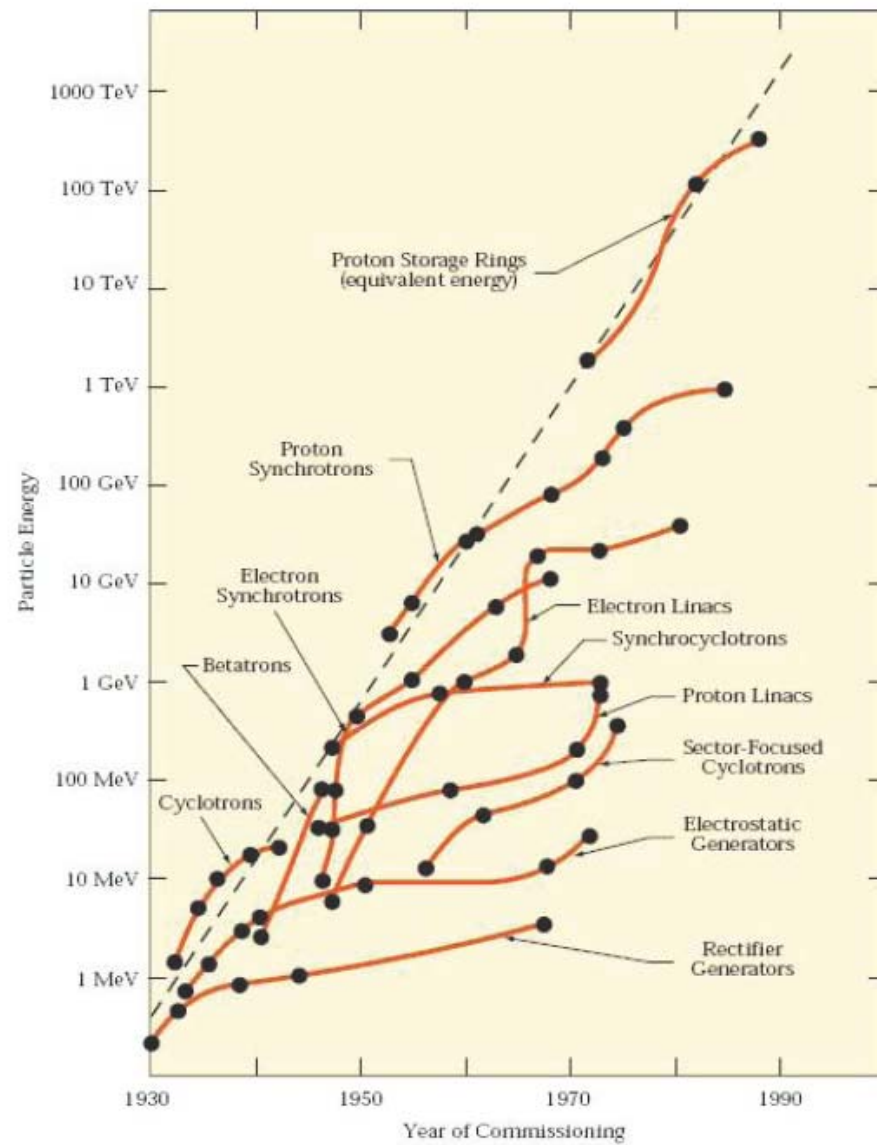


Aleph Detector at CERN LEP

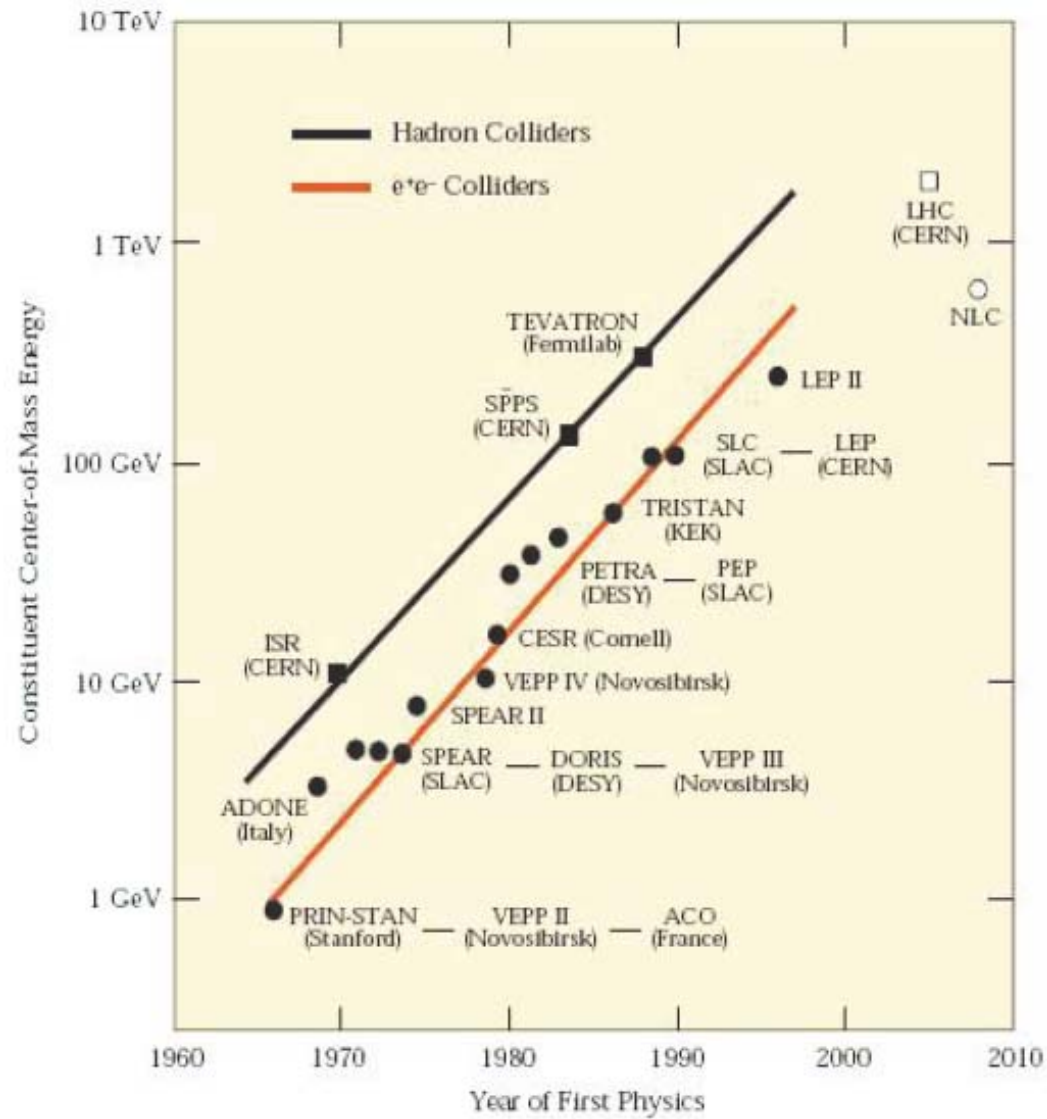


$$e^+ + e^- \rightarrow Z^0 \rightarrow q + \bar{q} + g \rightarrow \text{Hadronen}$$

Accelerator Evolution: Fixed target Experiment

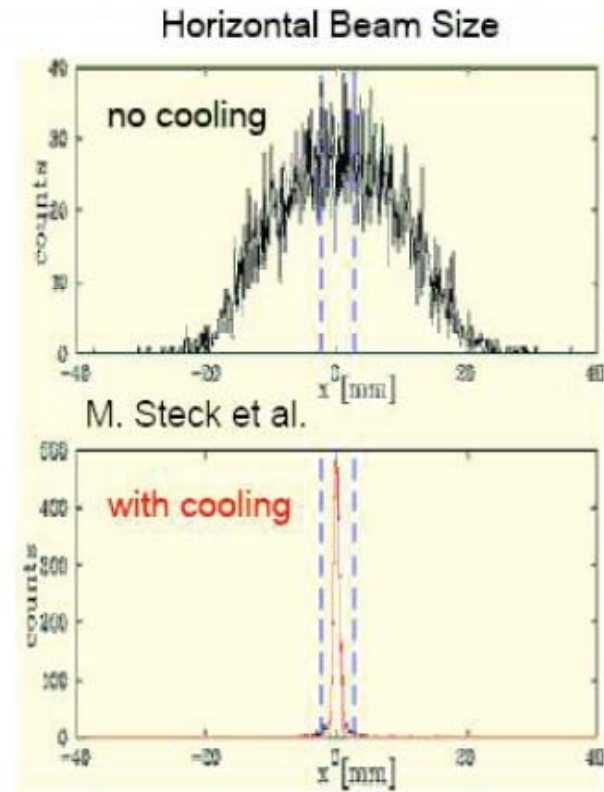
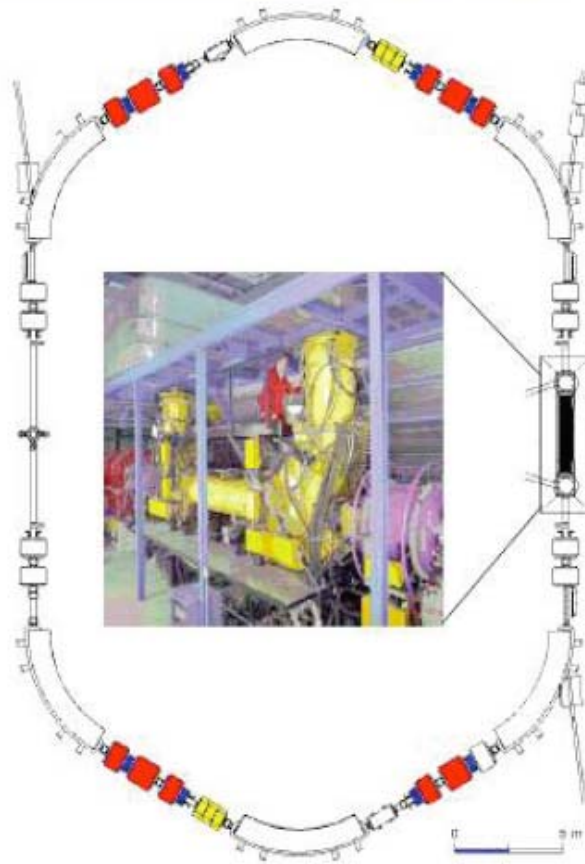


Accelerator Evolution: Colliders



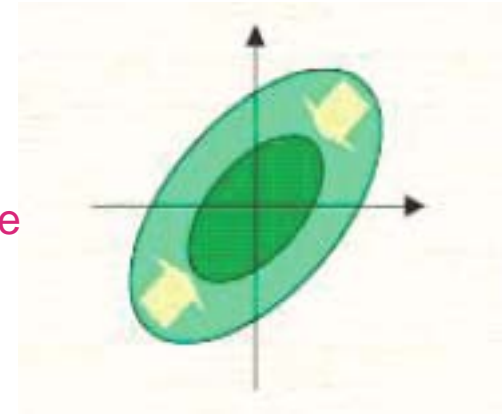
<i>Accelerator</i>	<i>Energy, GeV</i>
Proton Synchrotrons	
<i>CERNS PS</i>	28
<i>BNL AGS</i>	32
<i>KEK</i>	12
<i>SPS</i>	450
<i>Tevatron II</i>	1000
Electron Accelerators	
<i>SLAC linac</i>	25-50
<i>Desy Synchrotron</i>	7
Colliding-beam machines	
<i>PETRA</i>	$e^+e^- 22+22$
<i>LEP II</i>	$e^+e^- 100+100$
<i>HERA</i>	$ep 30e+820p$
<i>LHC</i>	$pp 7000+7000$

Electron Beam Cooling at ESR GSI



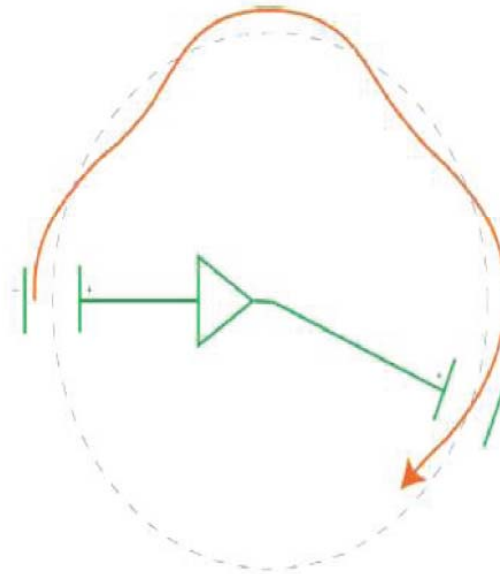
Cooling in the ESR counteracts the beam heating effect of the experimental insertions

Momentum spread due to the thermal motion.
Cooling should reduce the spread and hence increase the phase-space density



Principle of the stochastic Cooling

- *Measure beam center by pick-ups*
- *Correction signal to opposite kicker*



S. van der Meer included
in Nobel prize for W,Z