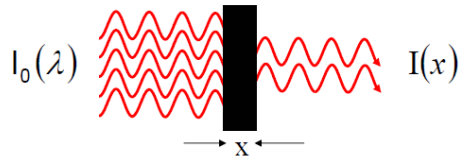


# Interaction of gamma rays with matter

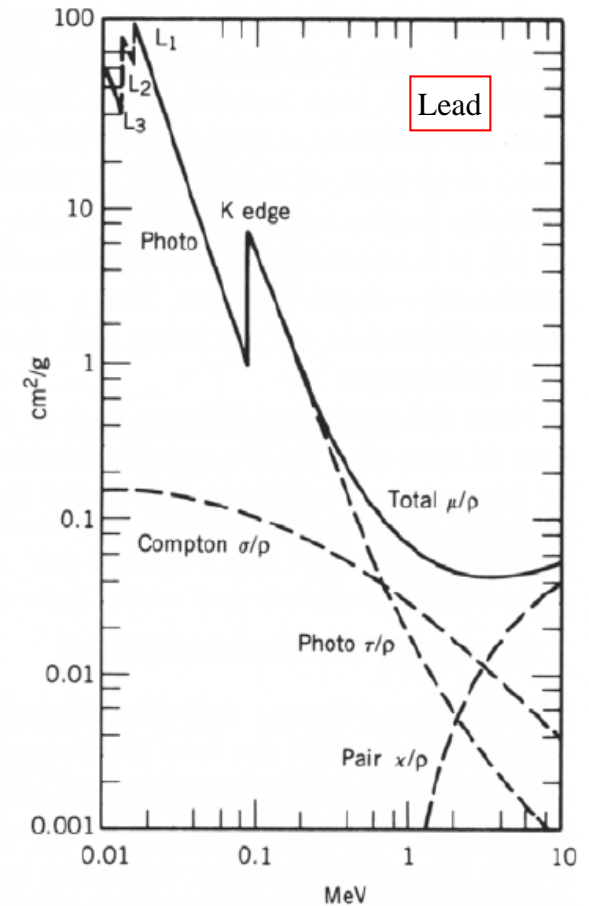
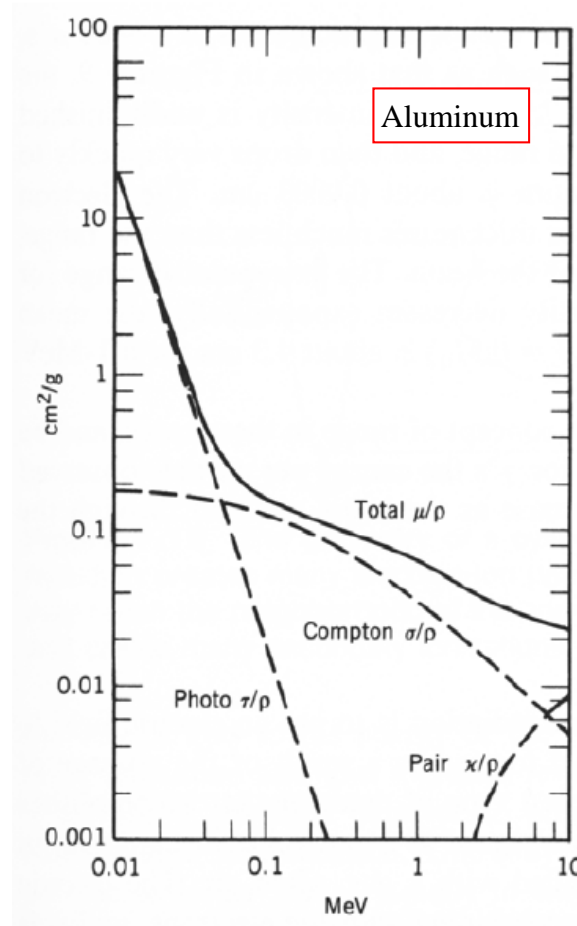


$$I(x) = I_0(\lambda) \cdot e^{-\frac{\mu(\lambda, Z)}{\rho} \rho \cdot x}$$

total absorption coefficient:  $\mu/\rho$  [ $\text{cm}^2/\text{g}$ ]

$$\frac{\mu_{total}}{\rho} = \sum_{i=1}^3 \sigma_i$$

- i=1 photoelectric effect
- i=2 Compton scattering
- i=3 pair production

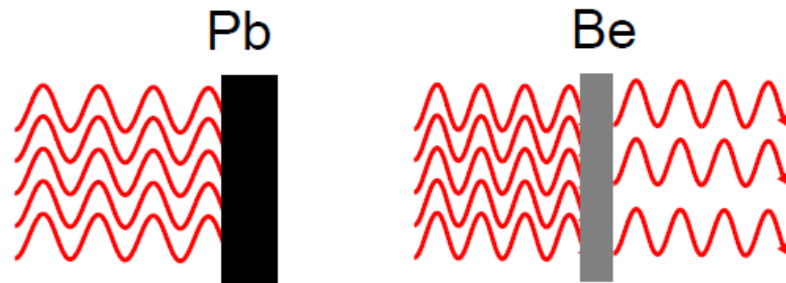


# Mass dependence of X-ray absorption

For X-ray radiation the **photoelectric effect** is the most important interaction.

$$(\mu / \rho)_{\text{Photo}} \approx \lambda^3 \cdot Z^5$$

**Lead absorbs more than Beryllium!**

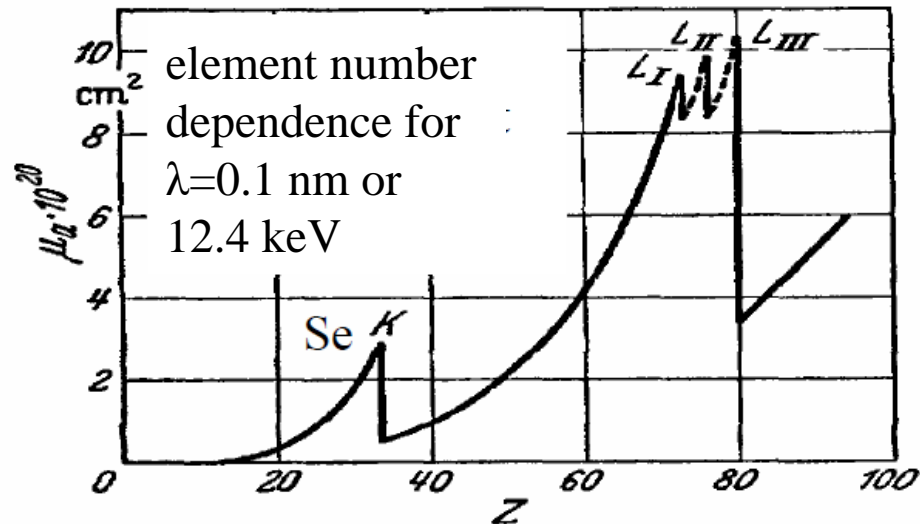
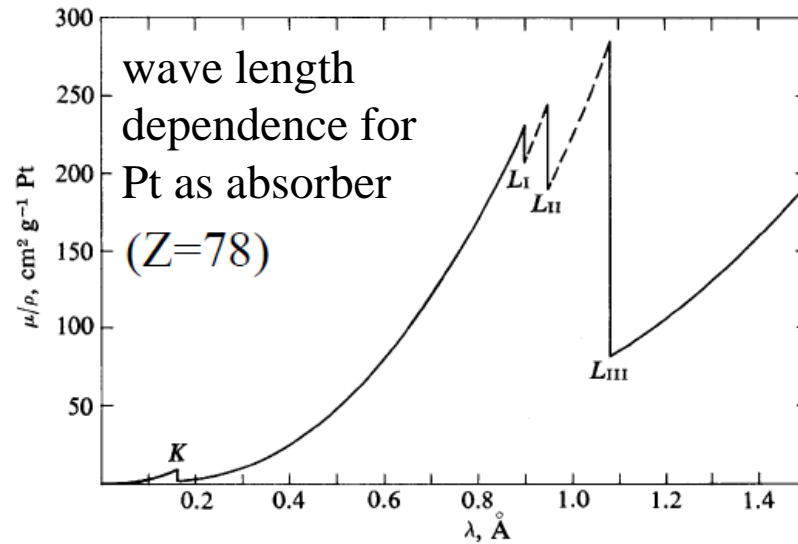


$^{82}\text{Pb}$  serves as shielding for X-ray and  $\gamma$ -ray radiation; lead vests are used by medical staff people who are exposed to X-ray radiation. Co-sources are transported in thick lead container.

On the contrary:

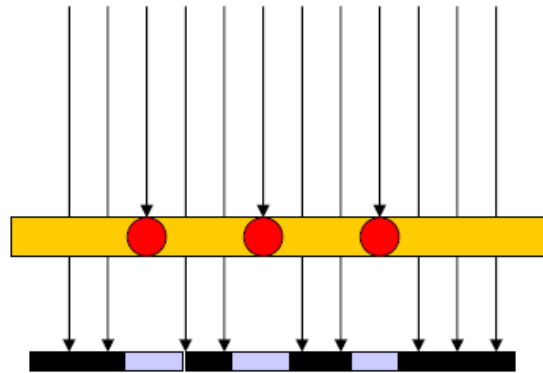
$^4\text{Be}$  is often used as windows in X-ray tubes to allow for almost undisturbed transmission of X-ray radiation.

# Mass dependence $\mu/\rho$ of X-ray absorption



# X-ray image shows the effect of different absorptions

Bones absorb more radiation as tissues because of their higher  $^{20}\text{Ca}$  content



# Interaction of gamma rays with matter

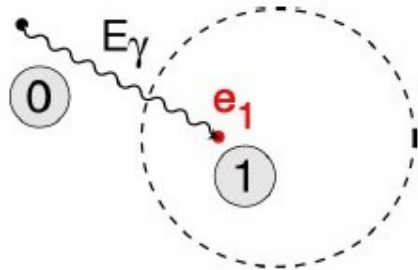
~ 100 keV

~1 MeV

~ 10 MeV

$\gamma$ -ray energy

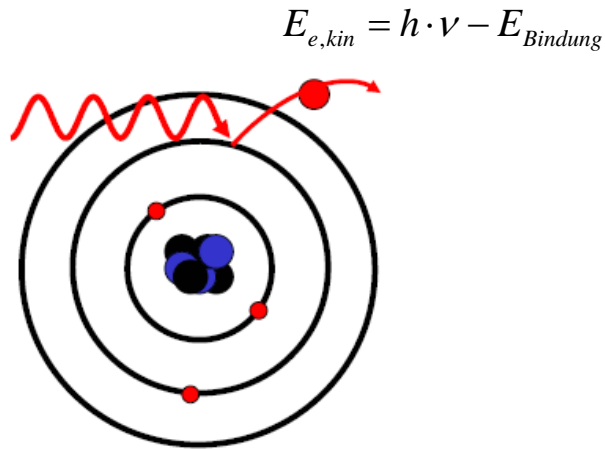
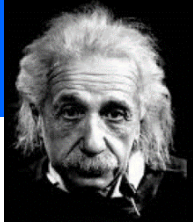
## Photoelectric



## Isolated hits

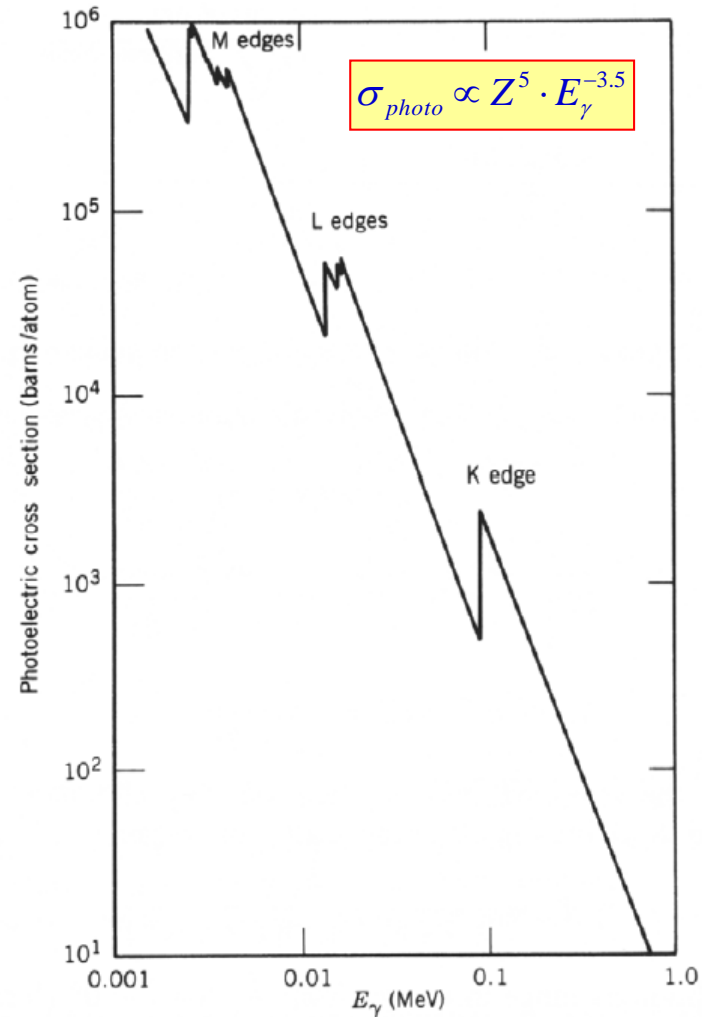
Probability of  
interaction depth

# Interaction of gamma rays with matter

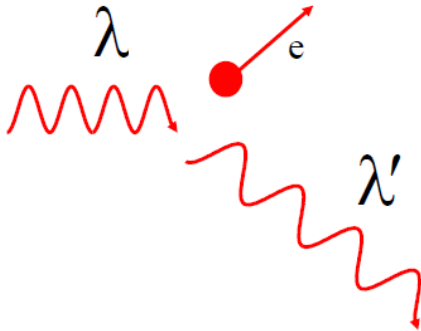
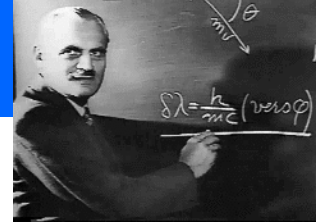


## Photo effect:

Absorption of a photon by a bound electron and conversion of the  $\gamma$ -energy in potential and kinetical energy of the ejected electron. (Nucleus preserves the momentum conservation.)



# Interaction of gamma rays with matter



relativistic  $E^2 = (pc)^2 + (m_0c^2)^2$  photons:  $m_0 = m_\gamma = 0$

$$\rightarrow E_\gamma = p_\gamma c$$

Momentum balance:

$$\vec{p}_e = \vec{p}_\gamma - \vec{p}'_\gamma \rightarrow |\vec{p}_e c|^2 = |(\vec{p}_\gamma - \vec{p}'_\gamma) c|^2$$

$$p_e^2 c^2 = E_\gamma^2 + E_{\gamma'}^2 - 2E_\gamma E_{\gamma'} \cdot \cos\theta$$

Energy balance:

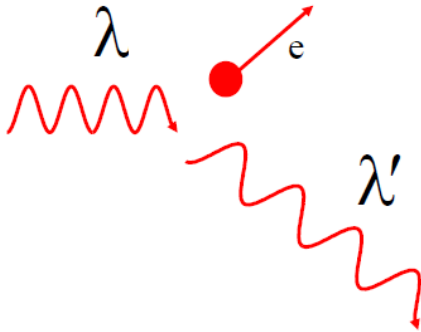
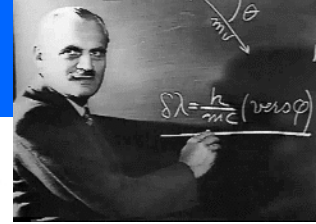
$$E_\gamma + m_e c^2 = E_{\gamma'} + \sqrt{(p_e c)^2 + (m_e c^2)^2}$$

$$E_{\gamma'} = \frac{E_\gamma}{1 + (E_\gamma/m_e c^2)(1 - \cos\theta)}$$

## Compton scattering:

Elastic scattering of a  $\gamma$ -ray on a free electron. A fraction of the  $\gamma$ -ray energy is transferred to the Compton electron. The wave length of the scattered  $\gamma$ -ray is increased:  $\lambda' > \lambda$ .

# Interaction of gamma rays with matter



Maximum energy of the scattered electron:

$$T(e^-)_{\max} = E_\gamma \cdot \frac{2 \cdot E_\gamma}{m_e c^2 + 2 \cdot E_\gamma}$$

Energy of the scattered  $\gamma$ -photon:

$$E'_\gamma = \frac{E_\gamma \cdot m_e c^2}{m_e c^2 + E_\gamma \cdot (1 - \cos \theta)}$$

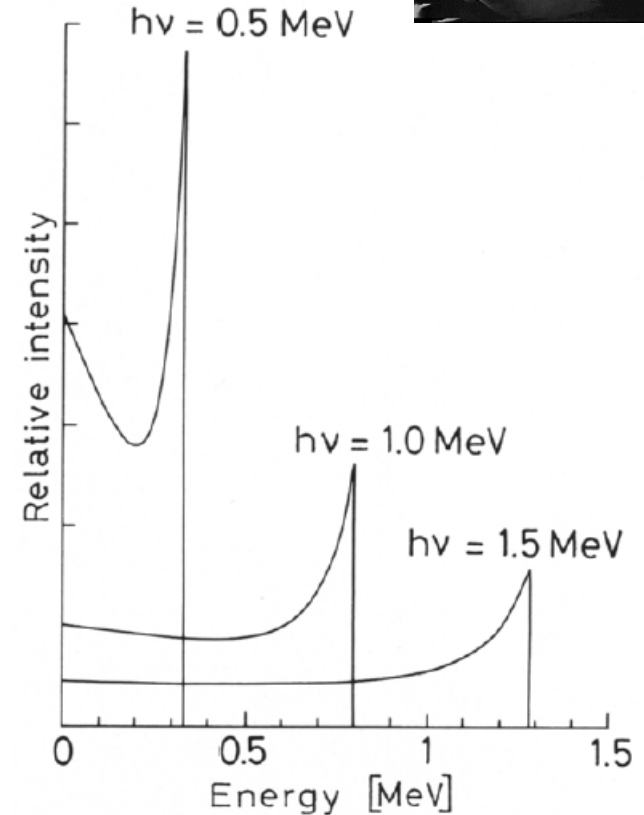
$$\cos \theta = 1 + \frac{m_e c^2}{E_\gamma} - \frac{m_e c^2}{E'_\gamma}$$

Special case for  $E \gg m_e c^2$ :  
 $\gamma$ -ray energy after  $180^\circ$  scatter is approximately

$$E'_\gamma = \frac{m_e c^2}{2} = 256 \text{ keV}$$

## Compton scattering:

Elastic scattering of a  $\gamma$ -ray on a free electron. A fraction of the  $\gamma$ -ray energy is transferred to the Compton electron. The wave length of the scattered  $\gamma$ -ray is increased:  $\lambda' > \lambda$ .

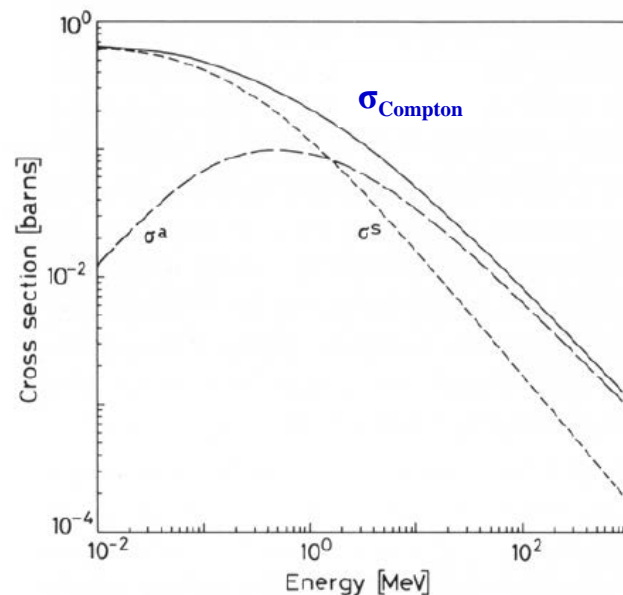
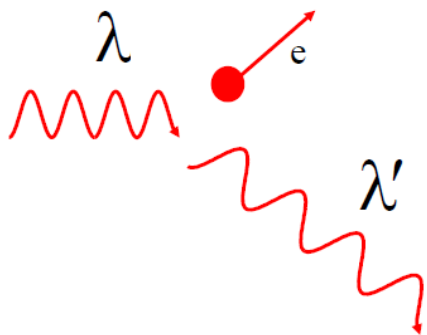


Gap between the incoming  $\gamma$ -ray and the maximum electron energy.

$$E_{\text{kin}}^{\max} = E_\gamma - E'_\gamma = E_\gamma \cdot \frac{2 \cdot E_\gamma / m_e c^2}{1 + 2 \cdot E_\gamma / m_e c^2}$$



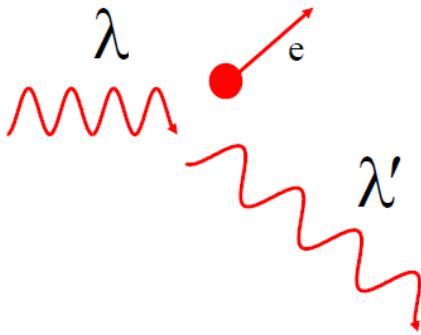
# Interaction of gamma rays with matter



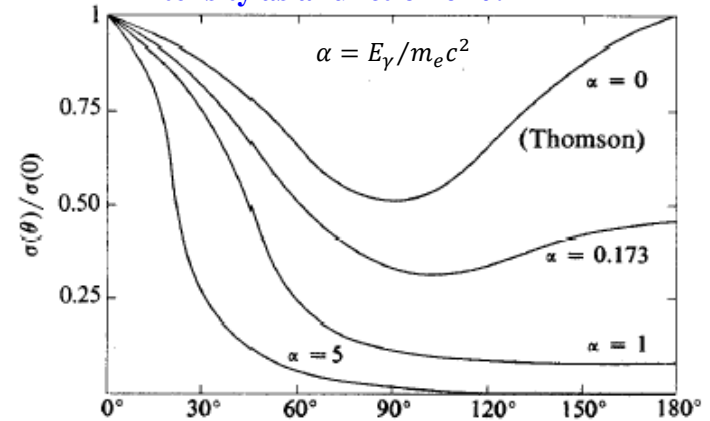
## Compton scattering:

Elastic scattering of a  $\gamma$ -ray on a free electron. A fraction of the  $\gamma$ -ray energy is transferred to the Compton electron. The wave length of the scattered  $\gamma$ -ray is increased:  $\lambda' > \lambda$ .

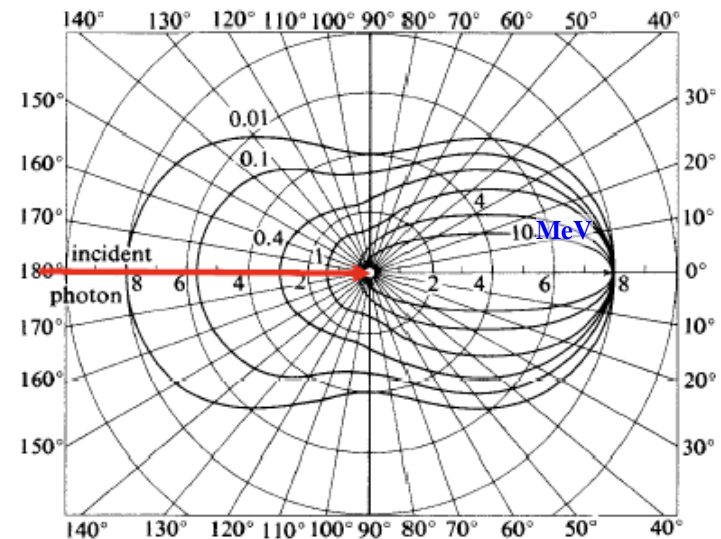
# Interaction of gamma rays with matter



Intensity as a function of  $\theta$ :



Angular distribution:



## Compton scattering:

Elastic scattering of a  $\gamma$ -ray on a free electron.  
The angle dependence is expressed by the

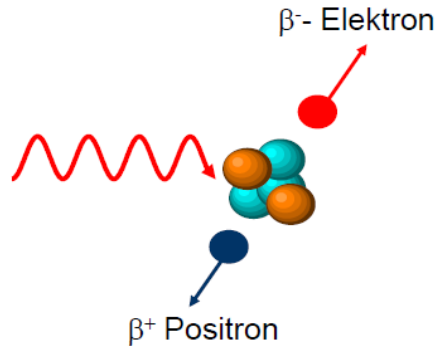
### Klein-Nishina-Formula:

$$\frac{d\sigma_c}{d\Omega} = \frac{r_0^2}{2} \left( \frac{E_{\gamma'}}{E_\gamma} \right)^2 \cdot \left\{ \frac{E_\gamma}{E_{\gamma'}} + \frac{E_{\gamma'}}{E_\gamma} - 2 \sin^2 \theta \cdot \cos^2 \phi \right\}$$

As shown in the plot **forward scattering** ( $\theta$  small) is dominant for  $E_\gamma > 100$  keV.

$r_0 = 2.818$  fm (classical electron radius)

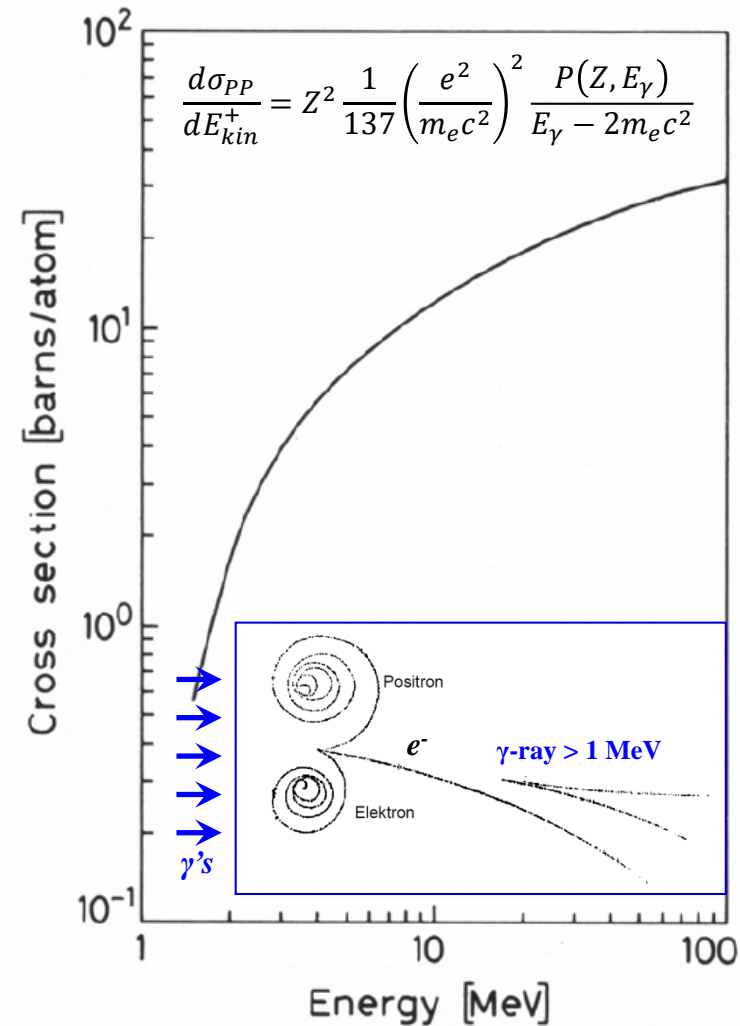
# Interaction of gamma rays with matter



## Pair production:

If  $\gamma$ -ray energy is  $\gg 2m_0c^2$  (electron rest mass 511 keV), a positron-electron pair can be formed in the strong Coulomb field of a nucleus. This pair carries the  $\gamma$ -ray energy minus  $2m_0c^2$ .

Pair production for  $E_\gamma > 2m_e c^2 = 1.022 \text{ MeV}$



picture of a bubble chamber

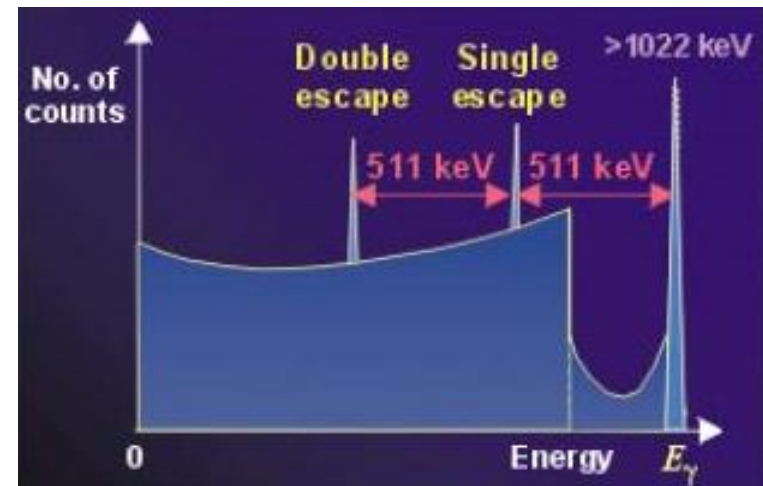
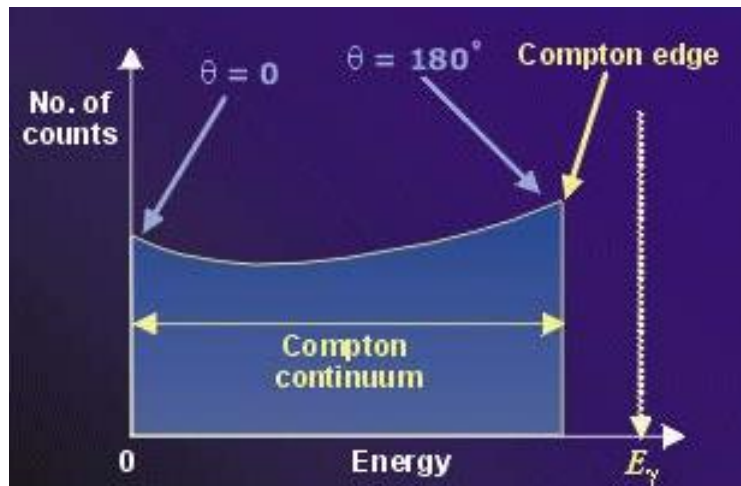
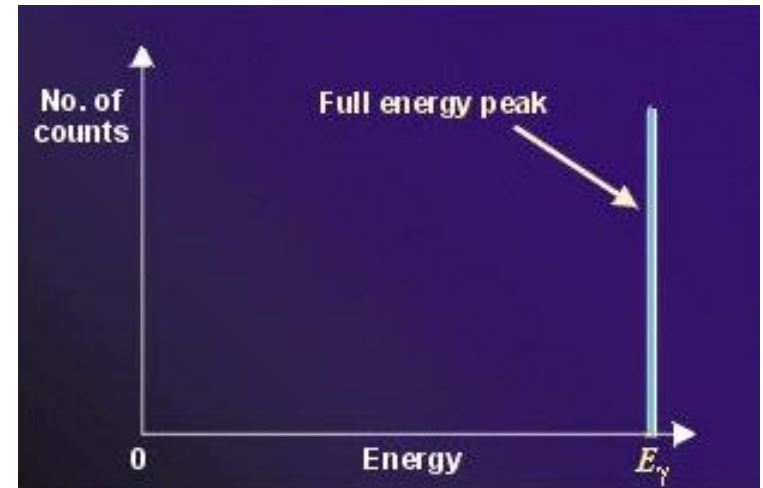
# Interaction of gamma rays with matter

$\gamma$ -rays interaction with matter via three main reaction mechanisms:

Photoelectric absorption

Compton scattering

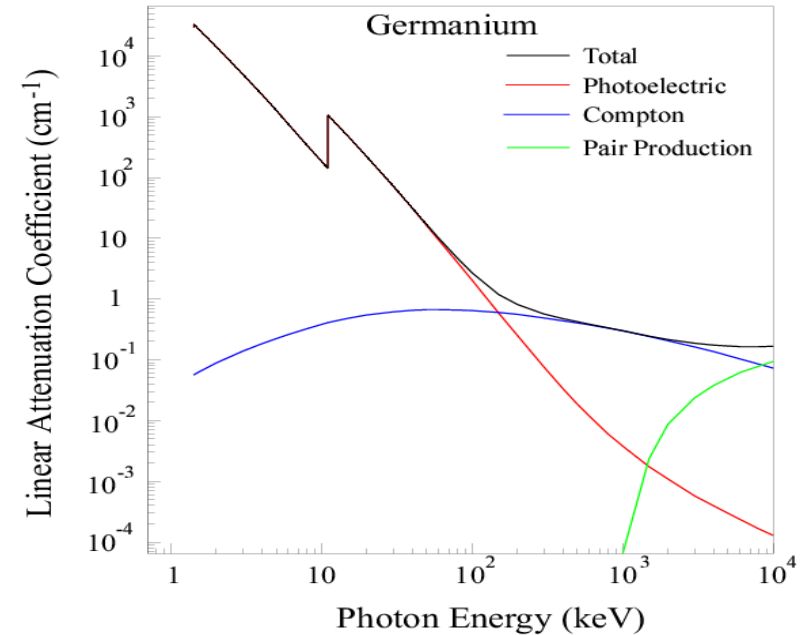
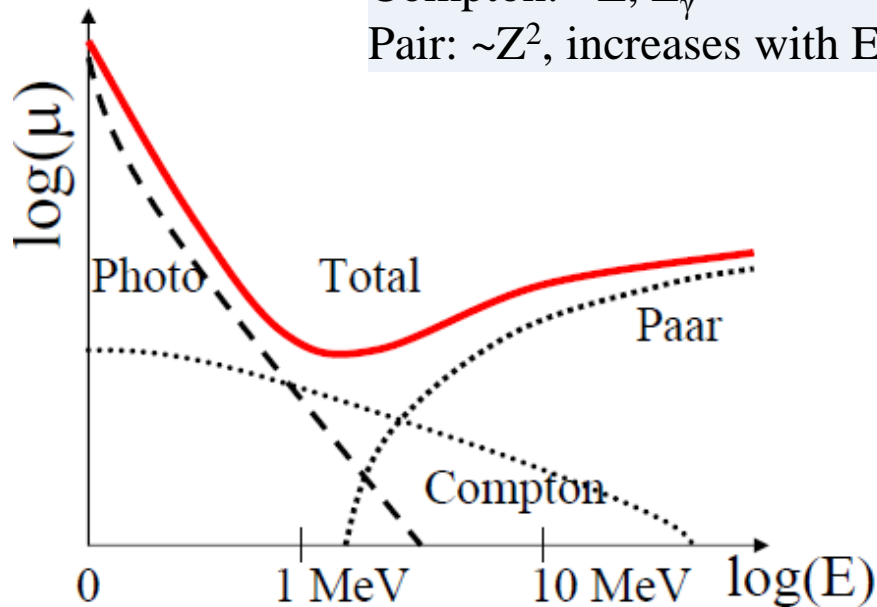
Pair production



# Gamma-ray interaction cross section

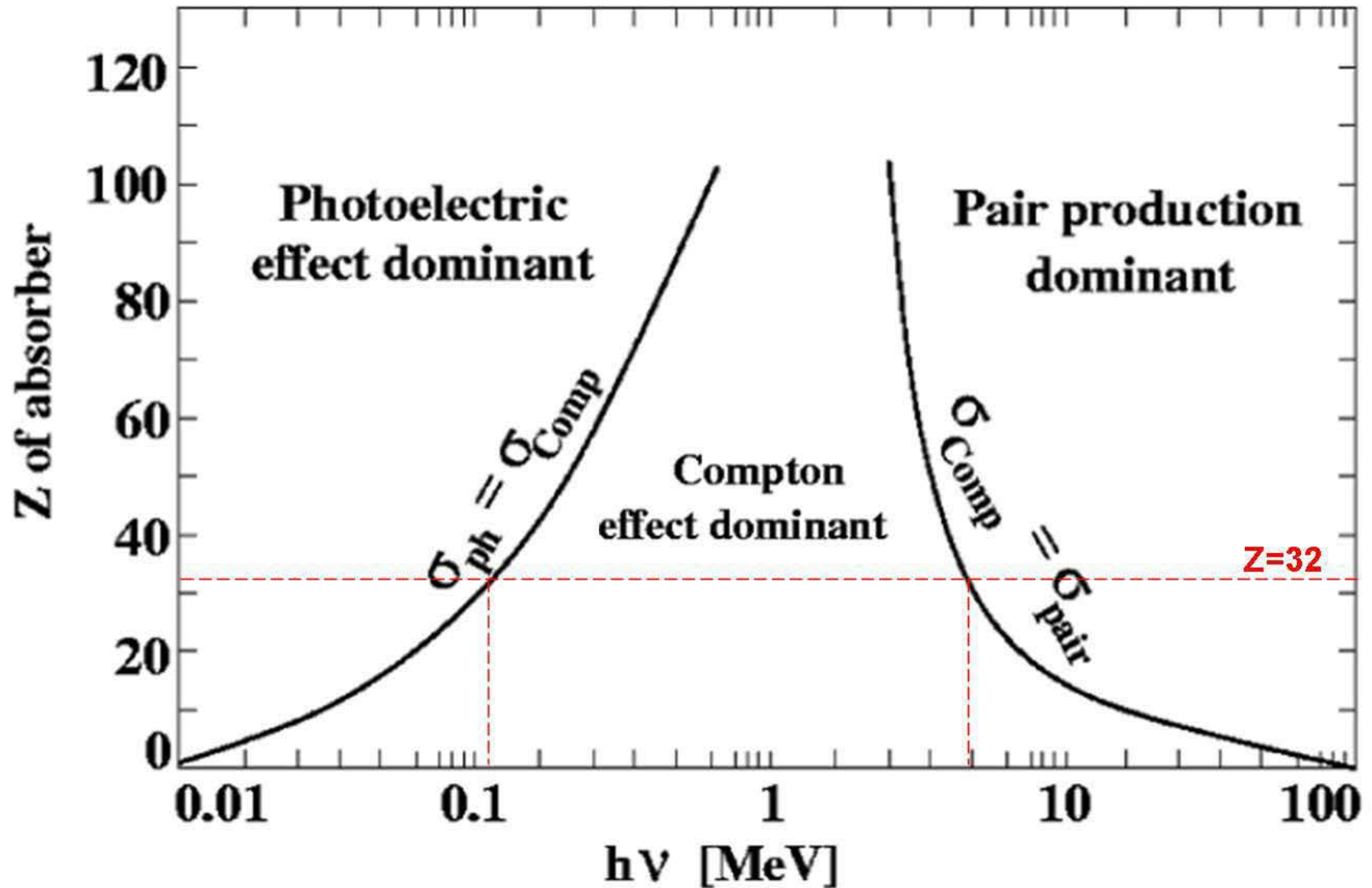
All three interaction (photo effect, Compton scattering and pair production) lead to an attenuation of the  $\gamma$ -ray or X-ray radiation when passing through matter. The particular contribution depends on the  $\gamma$ -ray energy:

Photo effect:  $\sim Z^{4-5}, E_{\gamma}^{-3.5}$   
Compton:  $\sim Z, E_{\gamma}^{-1}$   
Pair:  $\sim Z^2$ , increases with  $E_{\gamma}$



The absorption attenuates the intensity, but the energy and the frequency of the  $\gamma$ -ray and X-ray radiation is preserved!

# Z dependence of interaction probabilities



# Detector types

## Solid state semiconductor detectors: Ge

Electron-hole pairs are collected as charge

knock-on effect → an avalanche arrives at the electrode

lots of electrons → good energy resolution

cooled to liquid N<sub>2</sub> temperature (77K) to reduce noise

**Advantage:** good energy resolution (~0.15% FWHM at 1.3 MeV)

**Disadvantage:** relative low efficiency, cryogenic operation, limited size of crystal/detector

## Scintillation detectors: e.g. NaI, BGO, LaBr<sub>3</sub>(Ce)

Recoiling electrons excite atoms, which then de-excite by emitting visible light

Light is collected in photomultiplier tubes (PMT) where it generates a pulse proportional to the light collected

**Advantage:** good time resolution

detector can be made relative large e.g. NaI detector 14"Ø x 10"

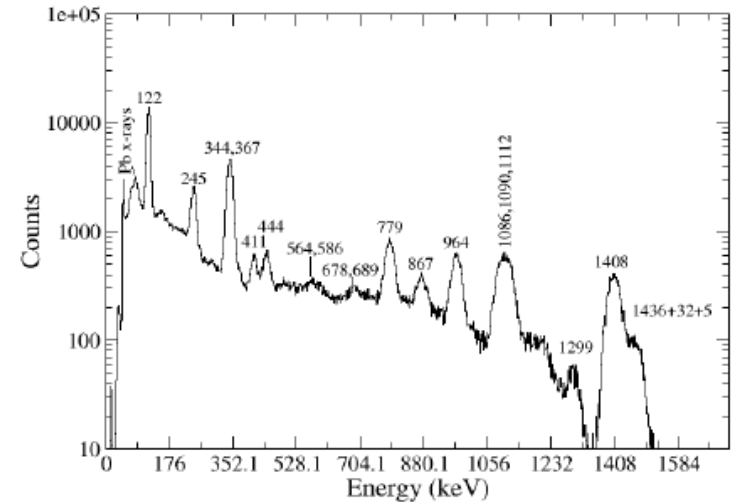
no need for cryogenics

**Disadvantage:** poor energy resolution (~5% FWHM at 1.3 MeV)

# Scintillation detectors

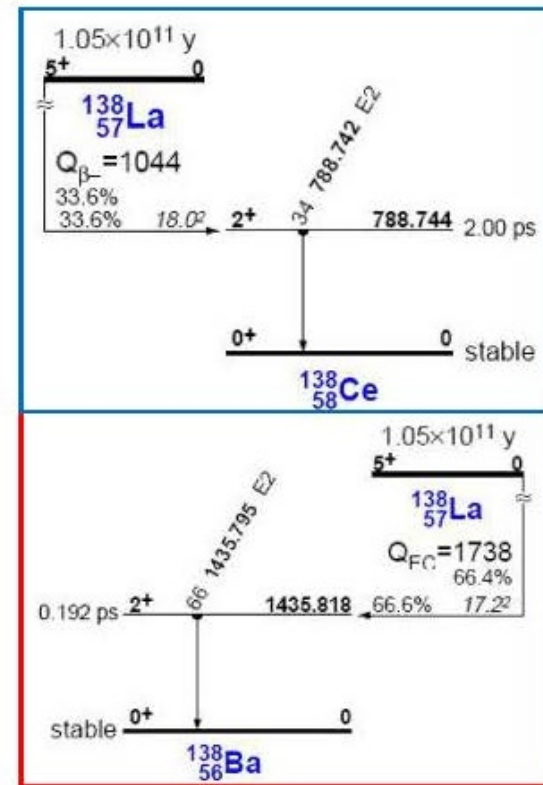
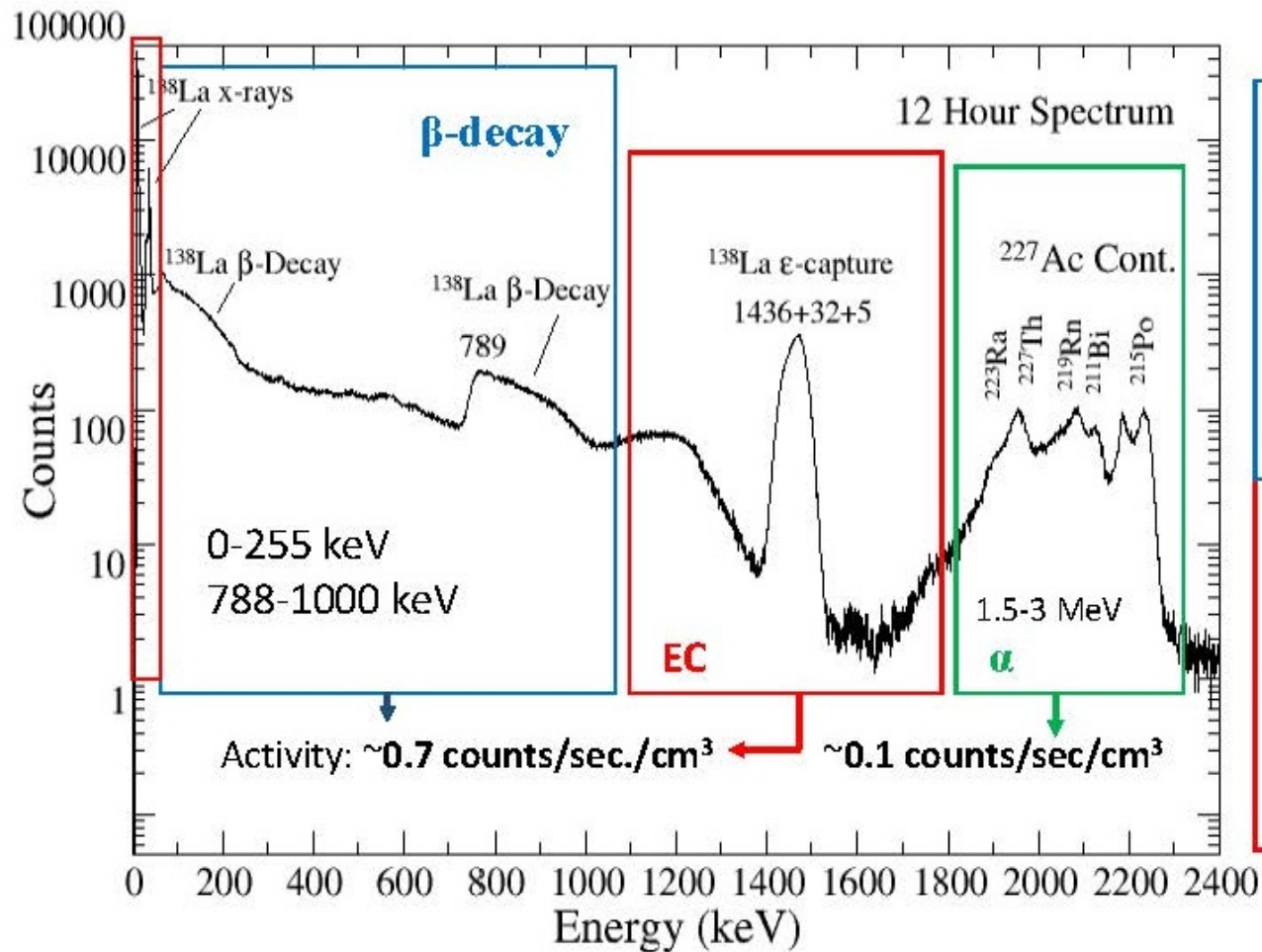
## LaBr<sub>3</sub>(Ce)

- ▶ LaBr<sub>3</sub>(Ce) timing properties:
  - ~ 25 ns decay time
  - Timing Resolution FWHM of 130-150 ps with <sup>60</sup>Co for a Ø1"x1" crystal.
- ▶ High energy resolution, 3 % FWHM at 662 keV.
- ▶ Peak Emission wavelength in Blue/UV part of EM spectrum (380 nm), compatible with PMTs.





# Detector characterization



J. McIntyre et al., NIM A 652, 1, 2011, 201-204

# Gamma-ray spectrum of a radioactive decay

