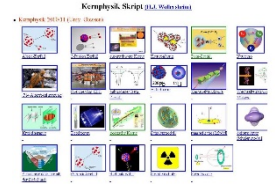


Outline: Nuclear surface vibration

Lecturer: Hans-Jürgen Wollersheim

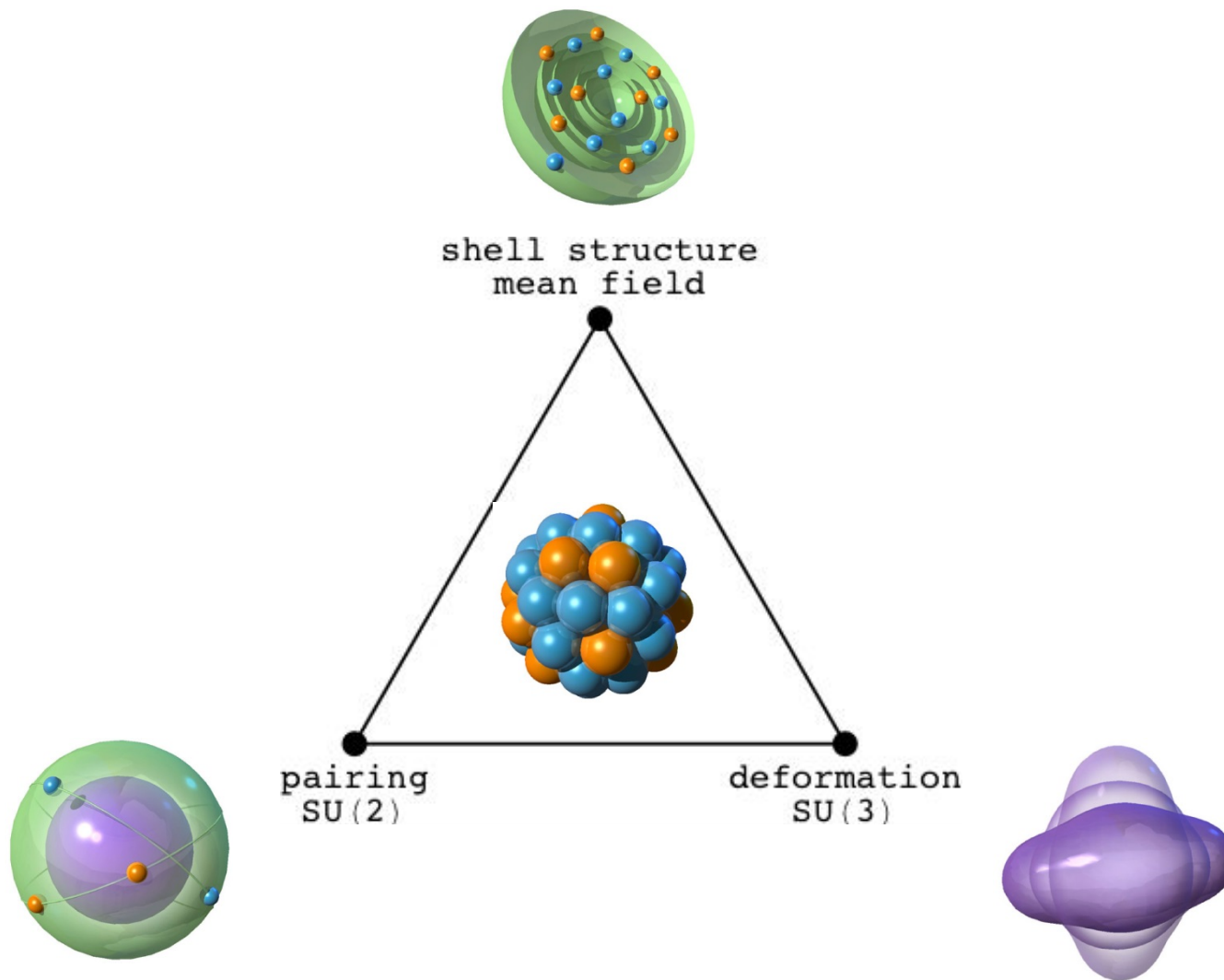
e-mail: h.j.wollersheim@gsi.de

web-page: <https://web-docs.gsi.de/~wolle/> and click on

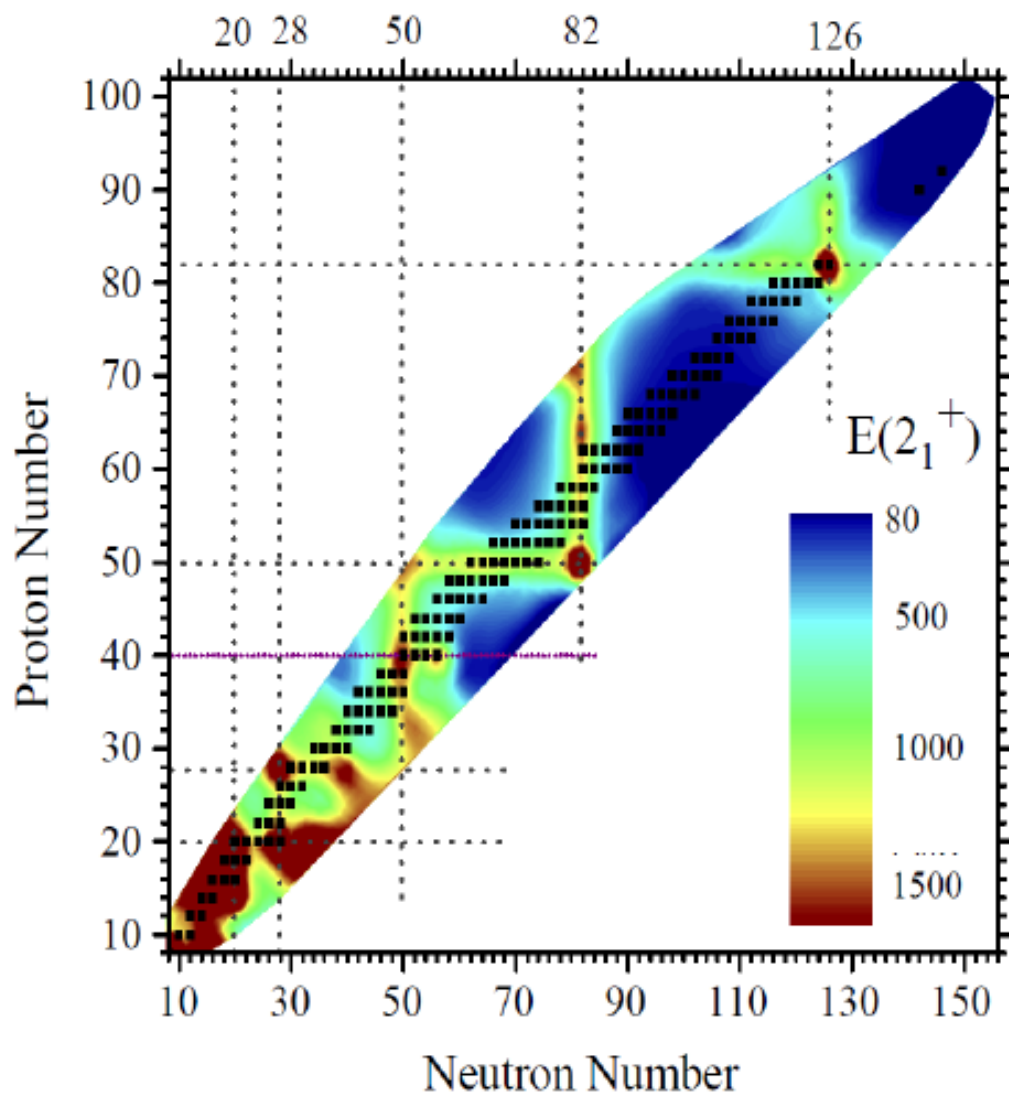


1. 2^+ , 4^+ systematics
2. collective vibration
3. second quantization
4. reduced transition probabilities

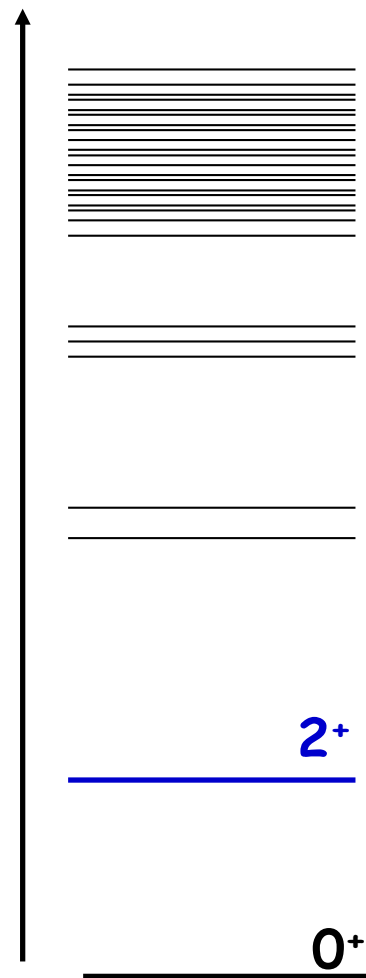
Nuclear surface vibration



2⁺ Systematics

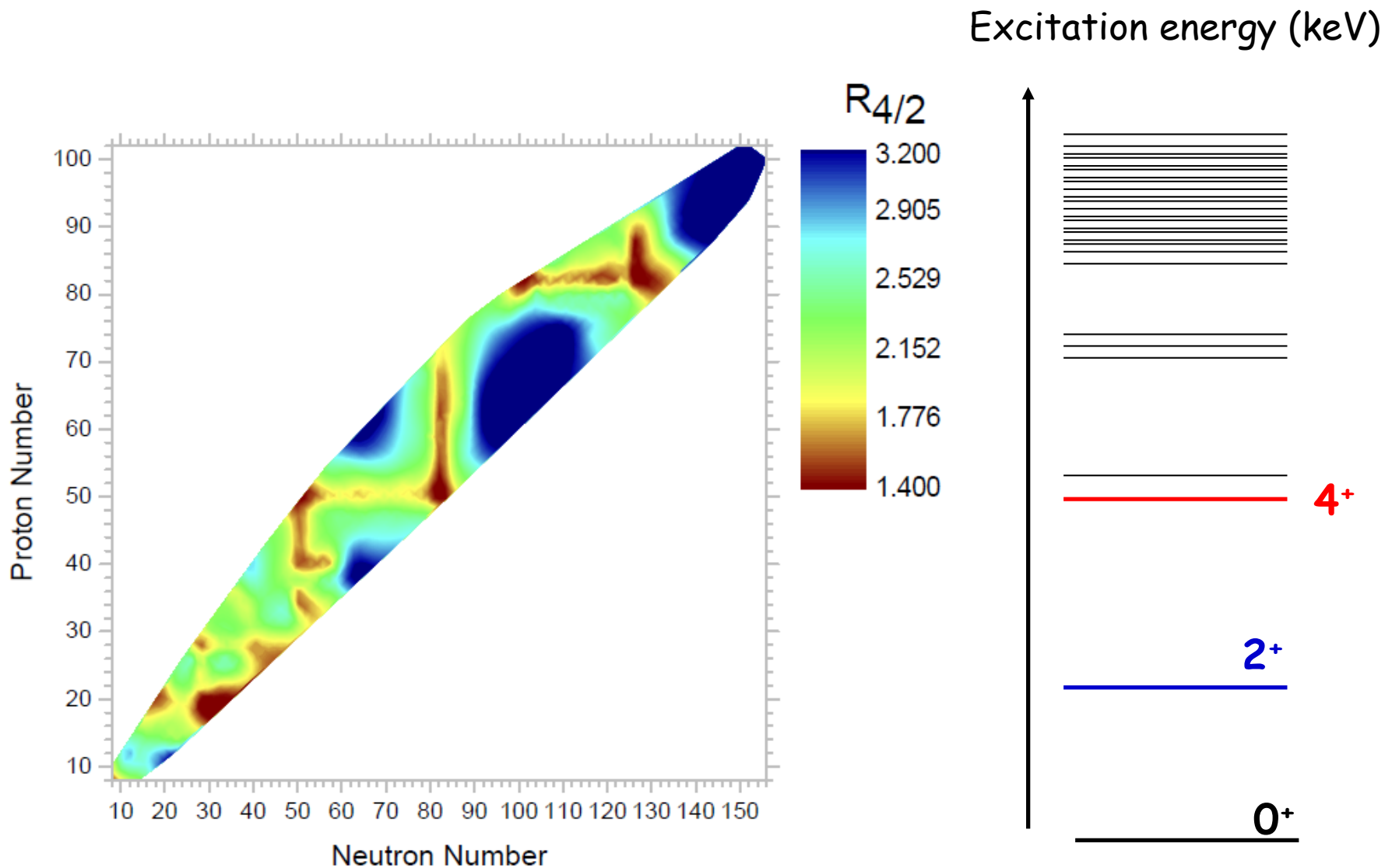


Excitation energy (keV)



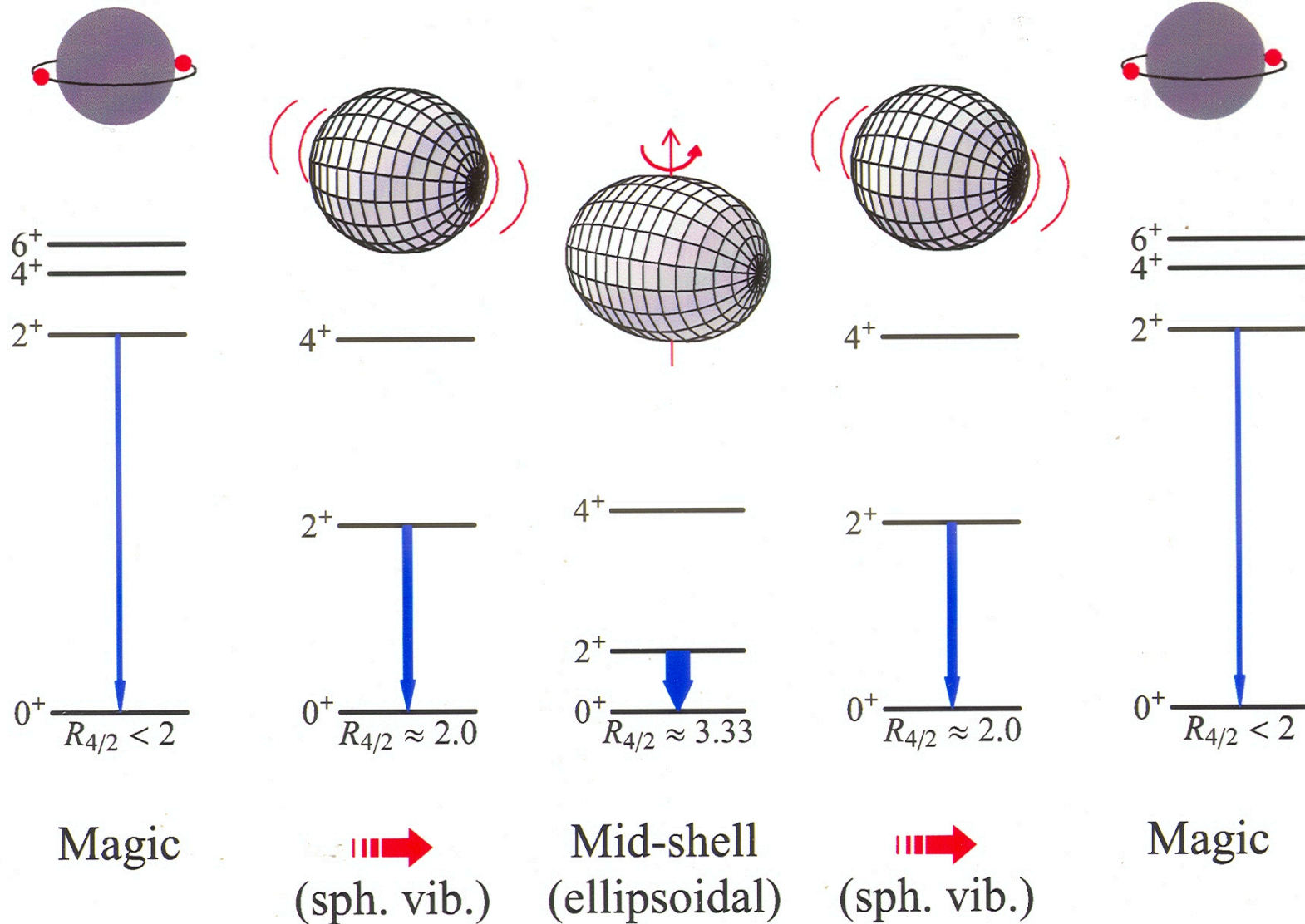
ground state configuration
spin/parity $I^\pi=0^+$; $E_x = 0$ keV

$4^+/2^+$ energy ratio: mirrors 2^+ systematics



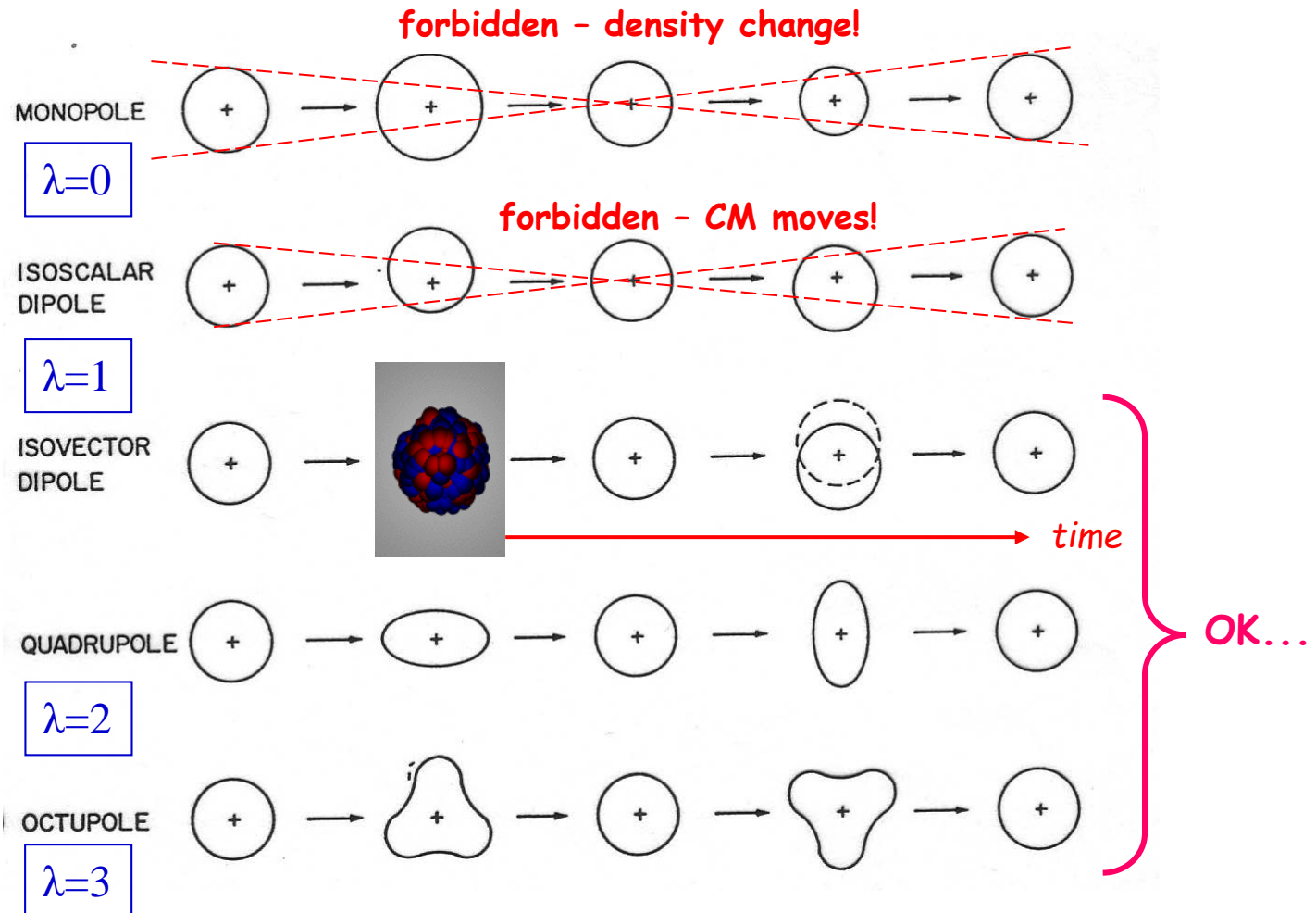
ground state configuration
spin/parity $I^\pi=0^+$; $E_x = 0$ keV

Evolution of nuclear structure as a function of nucleon number



Collective vibration

In general,
$$R(\theta, \phi, t) = R_0 \cdot \left[1 + \sum_{\lambda} \sum_{\mu=-\lambda}^{+\lambda} \alpha_{\lambda\mu}(t) \cdot Y_{\lambda\mu}^*(\theta, \phi) \right]$$

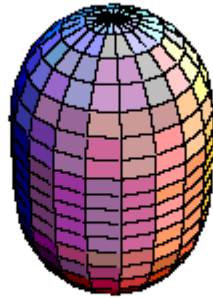


Collective vibration

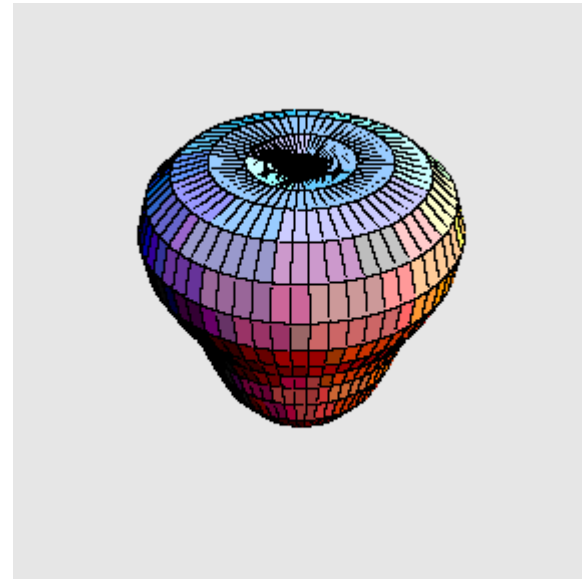


In general,
$$R(\theta, \phi, t) = R_0 \cdot \left[1 + \sum_{\lambda} \sum_{\mu=-\lambda}^{+\lambda} \alpha_{\lambda\mu}(t) \cdot Y_{\lambda\mu}^*(\theta, \phi) \right]$$

$$V = \frac{1}{2} \sum_{\lambda, \mu} C_{\lambda} \cdot |\alpha_{\lambda\mu}|^2 \quad \text{harmonic vibration}$$



$\lambda=2$: quadrupole vibration



$\lambda=3$: octupole vibration

Collective vibration

classical Hamiltonian

$$E = \frac{1}{2} \cdot \sum_{\lambda\mu} B_{\lambda} \cdot |\dot{\alpha}_{\lambda\mu}|^2 + \frac{1}{2} \cdot \sum_{\lambda\mu} C_{\lambda} \cdot |\alpha_{\lambda\mu}|^2$$

constants:

$$B_{\lambda} = \frac{3}{4 \cdot \pi} \frac{M \cdot A \cdot R_0^2}{\lambda}$$

$$C_{\lambda} = \frac{a_s \cdot A^{2/3}}{4 \cdot \pi} \cdot (\lambda - 1) \cdot (\lambda + 2) \cdot \left\{ 1 - \frac{6 \cdot (Z \cdot e)^2}{a_s \cdot r_0 \cdot A} \cdot \frac{1}{(2 \cdot \lambda + 1) \cdot (\lambda + 2)} \right\}$$

Binding energy of a nucleus:

$$B(A, Z) = a_V A - a_S A^{2/3} - a_C \frac{Z^2}{A^{1/3}} - a_A \frac{(N - Z)^2}{A} \pm \delta + \text{shell corr.}$$



$$a_V = 15.560 \text{ MeV} \quad a_S = 17.230 \text{ MeV} \quad a_C = 0.6970 \text{ MeV} \quad a_A = 23.385 \text{ MeV} \quad a_P = 12.000 \text{ MeV}$$

Collective vibration

classical Hamiltonian

$$E = \frac{1}{2} \cdot \sum_{\lambda\mu} B_\lambda \cdot |\dot{\alpha}_{\lambda\mu}|^2 + \frac{1}{2} \cdot \sum_{\lambda\mu} C_\lambda \cdot |\alpha_{\lambda\mu}|^2$$

$$H = \frac{1}{2} \cdot \sum_{\lambda\mu} \frac{1}{B_\lambda} \cdot |\pi_{\alpha_{\lambda\mu}}|^2 + C_\lambda \cdot |\alpha_{\lambda\mu}|^2$$

→ quantization

$$[\pi_{\alpha_{\lambda\mu}}, \alpha_{\lambda'\mu'}] = -i \cdot \hbar \cdot \delta_{\lambda\lambda'} \cdot \delta_{\mu\mu'}$$

$$H = \frac{1}{2} \cdot \sum_{\lambda\mu} -\frac{\hbar^2}{B_\lambda} \frac{\partial^2}{\partial x_{\lambda\mu}^2} + C_\lambda \cdot x_{\lambda\mu}^2$$

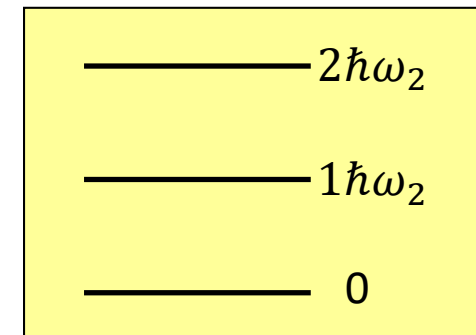
energy eigenvalue

$$E = \sum_{\lambda\mu} \hbar \cdot \omega_\lambda \cdot \left(n_{\lambda\mu} + \frac{1}{2} \right)$$

wave function

$$\psi_n = N_n \cdot H_n(x) \cdot e^{-x^2/2}$$

Hermite polynomials



Second quantization

Hamilton operator

$$H = \frac{1}{2} \cdot \sum_{\lambda\mu} \frac{1}{B_\lambda} \cdot |\pi_{\alpha_{\lambda\mu}}|^2 + C_\lambda \cdot |\alpha_{\lambda\mu}|^2$$

$$H = \sum_{\lambda\mu} \hbar \cdot \omega_\lambda \cdot \left(\beta_{\lambda\mu}^+ \beta_{\lambda\mu} + \frac{1}{2} \right)$$

$$\alpha_{\lambda\mu} = \sqrt{\frac{\hbar}{2B_\lambda\omega_\lambda}} \cdot (\beta_{\lambda\mu}^+ + (-)^\mu \beta_{\lambda\mu})$$

$$\pi_{\alpha_{\lambda\mu}} = \sqrt{\frac{\hbar B_\lambda\omega_\lambda}{2}} \cdot ((-)^\mu \beta_{\lambda-\mu}^+ - \beta_{\lambda\mu})$$

rule for boson operators: $\beta_{\lambda\mu} \cdot \beta_{\lambda'\mu'}^+ - \beta_{\lambda\mu}^+ \cdot \beta_{\lambda'\mu'} = \delta_{\lambda\lambda'} \cdot \delta_{\mu\mu'}$

Second quantization

Hamilton operator $H = \sum_{\lambda\mu} \hbar \cdot \omega_{\lambda} \cdot \left(\beta_{\lambda\mu}^+ \beta_{\lambda\mu} + \frac{1}{2} \right)$

rule for boson operators: $\beta_{\lambda\mu} \cdot \beta_{\lambda'\mu'}^+ - \beta_{\lambda\mu}^+ \cdot \beta_{\lambda'\mu'} = \delta_{\lambda\lambda'} \cdot \delta_{\mu\mu'}$

creation & annihilation operators

increase or decrease the number of phonons in a wave function

ground state $|0\rangle$ (vacuum)

1-phonon state $\beta_{\lambda\mu}^+ |0\rangle = |1\rangle$ or $\langle 1| = \langle 0| \beta_{\lambda\mu}$

2-phonon state $\sim \beta_{\lambda\mu}^+ \beta_{\lambda'\mu'}^+ |0\rangle = |2\rangle$ or $\langle 2| \sim \langle 0| \beta_{\lambda\mu} \beta_{\lambda'\mu'}$

Remember for the following calculations: $\beta_{\lambda\mu} |0\rangle = 0$

Second quantization

Hamilton operator $H = \sum_{\lambda\mu} \hbar \cdot \omega_{\lambda} \cdot \left(\beta_{\lambda\mu}^{\dagger} \beta_{\lambda\mu} + \frac{1}{2} \right)$

rule for boson operators: $\beta_{\lambda\mu} \cdot \beta_{\lambda'\mu'}^{\dagger} - \beta_{\lambda\mu}^{\dagger} \cdot \beta_{\lambda'\mu'} = \delta_{\lambda\lambda'} \cdot \delta_{\mu\mu'}$

1-phonon energy:

$$\langle 1 | H | 1 \rangle = \hbar \cdot \omega \cdot \left(\langle 0 | \beta \beta^{\dagger} \beta \beta^{\dagger} | 0 \rangle + \frac{1}{2} \langle 0 | \beta \beta^{\dagger} | 0 \rangle \right)$$

$$\langle 1 | H | 1 \rangle = \hbar \cdot \omega \cdot \left(\langle 0 | \beta \beta^{\dagger} (1 + \beta^{\dagger} \beta) | 0 \rangle + \frac{1}{2} \langle 0 | (1 + \beta^{\dagger} \beta) | 0 \rangle \right)$$

$$\langle 1 | H | 1 \rangle = \hbar \cdot \omega \cdot \left(\langle 0 | \beta \beta^{\dagger} | 0 \rangle + \frac{1}{2} \langle 0 | 1 | 0 \rangle \right)$$

$$\langle 1 | H | 1 \rangle = \hbar \cdot \omega \cdot \left(\langle 0 | (1 + \beta^{\dagger} \beta) | 0 \rangle + \frac{1}{2} \right)$$

$$\langle 1 | H | 1 \rangle = \hbar \cdot \omega \cdot \left(1 + \frac{1}{2} \right)$$

Second quantization

Hamilton operator $H = \sum_{\lambda\mu} \hbar \cdot \omega_{\lambda} \cdot \left(\beta_{\lambda\mu}^{\dagger} \beta_{\lambda\mu} + \frac{1}{2} \right)$

rule for boson operators: $\beta_{\lambda\mu} \cdot \beta_{\lambda'\mu'}^{\dagger} - \beta_{\lambda\mu}^{\dagger} \cdot \beta_{\lambda'\mu'} = \delta_{\lambda\lambda'} \cdot \delta_{\mu\mu'}$

2-phonon energy:

$$\begin{aligned} \langle 2 | H | 2 \rangle &\sim \hbar \cdot \omega \cdot \left(\langle 0 | \beta\beta\beta^{\dagger} \beta\beta^{\dagger} \beta^{\dagger} | 0 \rangle + \frac{1}{2} \langle 0 | \beta\beta\beta^{\dagger} \beta^{\dagger} | 0 \rangle \right) \\ \langle 2 | H | 2 \rangle &\sim \hbar \cdot \omega \cdot \left(\langle 0 | \beta\beta\beta^{\dagger} (1 + \beta^{\dagger} \beta) \beta^{\dagger} | 0 \rangle + \frac{1}{2} \langle 0 | \beta (1 + \beta^{\dagger} \beta) \beta^{\dagger} | 0 \rangle \right) \\ \langle 2 | H | 2 \rangle &\sim \hbar \cdot \omega \cdot \left(\langle 0 | \beta\beta\beta^{\dagger} \beta^{\dagger} + \beta\beta\beta^{\dagger} \beta^{\dagger} \beta\beta^{\dagger} | 0 \rangle + \frac{1}{2} \langle 0 | \beta\beta^{\dagger} + \beta\beta^{\dagger} \beta\beta^{\dagger} | 0 \rangle \right) \end{aligned}$$

etc.

$\langle 2 | H | 2 \rangle \sim 2 \cdot \hbar \cdot \omega \cdot \left(2 + \frac{1}{2} \right)$ normalization of wave function !!!

Second quantization

Hamilton operator $H = \sum_{\lambda\mu} \hbar \cdot \omega_{\lambda} \cdot \left(\beta_{\lambda\mu}^+ \beta_{\lambda\mu} + \frac{1}{2} \right)$

rule for boson operators: $\beta_{\lambda\mu} \cdot \beta_{\lambda'\mu'}^+ - \beta_{\lambda\mu}^+ \cdot \beta_{\lambda'\mu'} = \delta_{\lambda\lambda'} \cdot \delta_{\mu\mu'}$

2-phonon state: $\Psi_{IM} = N \sum_{m_1 m_2} (1 \ 1 \ 2 \ m_1 m_2 \ | \ IM) \cdot \beta_{1 \ m_1}^+ \cdot \beta_{2 \ m_2}^+ \ |0\rangle$

normalization (approximation):

$$1 = N^2 \langle 0 | \beta \beta \beta^+ \beta^+ | 0 \rangle$$

$$1 = N^2 \langle 0 | \beta (1 + \beta^+ \beta) \beta^+ | 0 \rangle$$

$$1 = N^2 \langle 0 | \beta \beta^+ + \beta \beta^+ \beta \beta^+ | 0 \rangle$$

$$1 = N^2 \langle 0 | (1 + \beta^+ \beta) + (1 + \beta^+ \beta)(1 + \beta^+ \beta) | 0 \rangle = N^2 \cdot 2$$

Second quantization

Hamilton operator $H = \sum_{\lambda\mu} \hbar \cdot \omega_{\lambda} \cdot \left(\beta_{\lambda\mu}^+ \beta_{\lambda\mu} + \frac{1}{2} \right)$

rule for boson operators: $\beta_{\lambda\mu} \cdot \beta_{\lambda'\mu'}^+ - \beta_{\lambda\mu}^+ \cdot \beta_{\lambda'\mu'} = \delta_{\lambda\lambda'} \cdot \delta_{\mu\mu'}$

2-phonon state: $\Psi_{IM} = N \sum_{m_1 m_2} (1 \ 1 | \ 2 \ m_1 m_2 | IM) \cdot \beta_{1 \ m_1}^+ \cdot \beta_{1 \ 2 m_2}^+ | 0 \rangle$

normalization:

$$1 = N^2 \sum_{m_1 m_2 m'_1 m'_2} (1 \ 1 | \ 2 \ m_1 m_2 | IM) \cdot (1 \ 1 | \ 2 \ m'_1 m'_2 | IM) \cdot \langle 0 | \beta_{1 \ m_1} \beta_{1 \ 2 m_2} \beta_{1 \ m'_1}^+ \beta_{1 \ 2 m'_2}^+ | 0 \rangle$$

$$1 = N^2 \sum_{m_1 m_2 m'_1 m'_2} (1 \ 1 | \ 2 \ m_1 m_2 | IM) \cdot (1 \ 1 | \ 2 \ m'_1 m'_2 | IM) \cdot \langle 0 | \beta_{1 \ m_1} [\delta_{1 \ 1 \ 2} \delta_{m_2 m'_1} + \beta_{1 \ m'_1} \beta_{1 \ 2 m_2}] \beta_{1 \ 2 m'_2}^+ | 0 \rangle$$

$$1 = N^2 \sum_{m_1 m_2 m'_1 m'_2} (1 \ 1 | \ 2 \ m_1 m_2 | IM) \cdot (1 \ 1 | \ 2 \ m'_1 m'_2 | IM) \cdot [\delta_{1 \ 1 \ 2} \delta_{m_2 m'_1} \delta_{m_1 m'_2} + \delta_{m_1 m'_1} \delta_{m_2 m'_2}]$$

$$1 = N^2 \sum_{m_1 m_2} (1 \ 1 | \ 2 \ m_1 m_2 | IM) \cdot \left\{ (1 \ 1 | \ 2 \ m_1 m_2 | IM) + \delta_{1 \ 1 \ 2} (1 \ 1 | \ 2 \ m_2 m_1 | IM) \right\}$$

$$1 = N^2 \sum_{m_1 m_2} (1 \ 1 | \ 2 \ m_1 m_2 | IM) \cdot (1 \ 1 | \ 2 \ m_1 m_2 | IM) \left\{ 1 + (-)^l \delta_{1 \ 1 \ 2} \right\} = N^2 \left\{ 1 + (-)^l \delta_{1 \ 1 \ 2} \right\}$$

Collective vibration

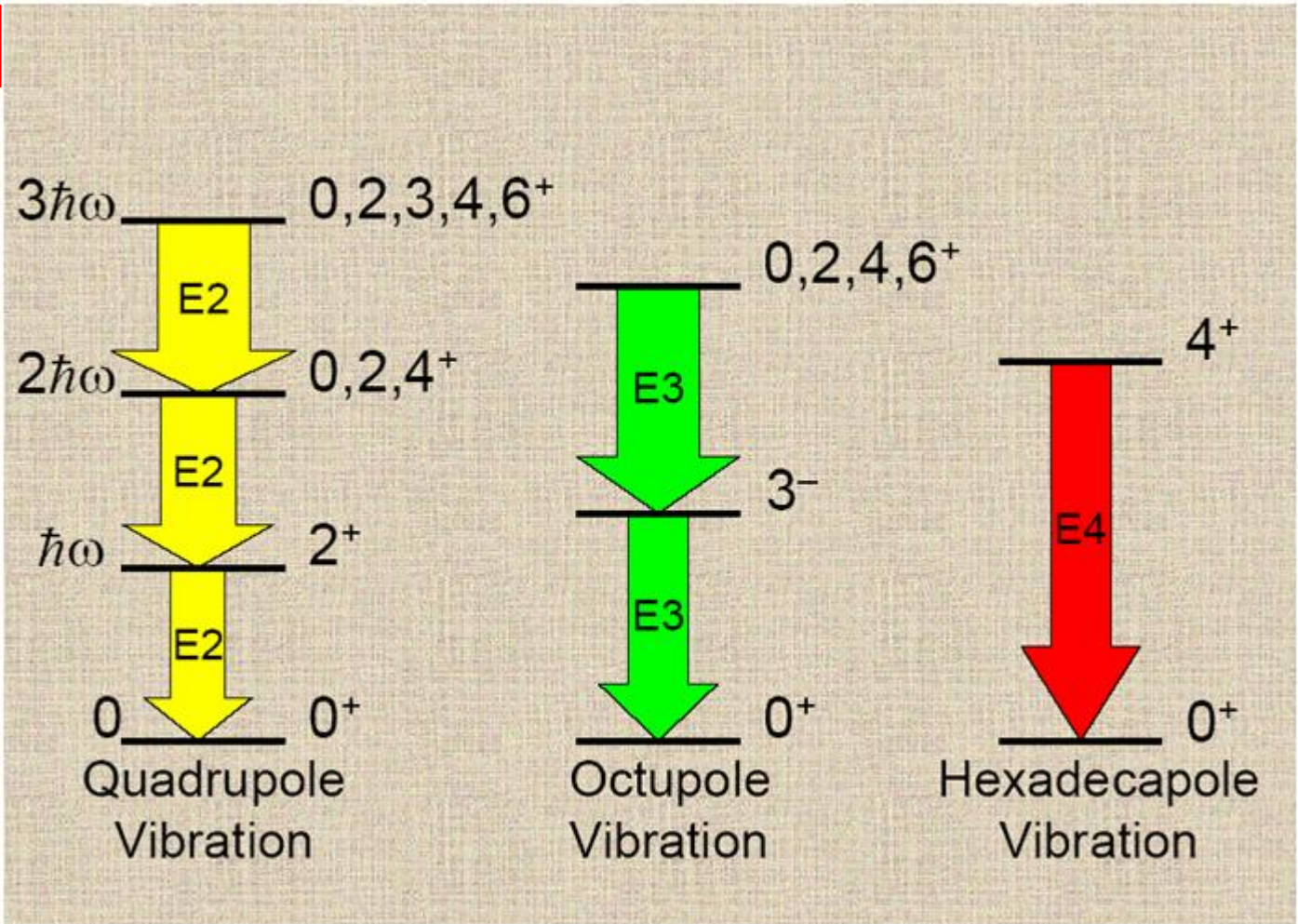
$$E_N = \hbar\omega \cdot \left(N + \frac{5}{2}\right)$$

$$E_3 = \hbar\omega \cdot \left(3 + \frac{5}{2}\right)$$

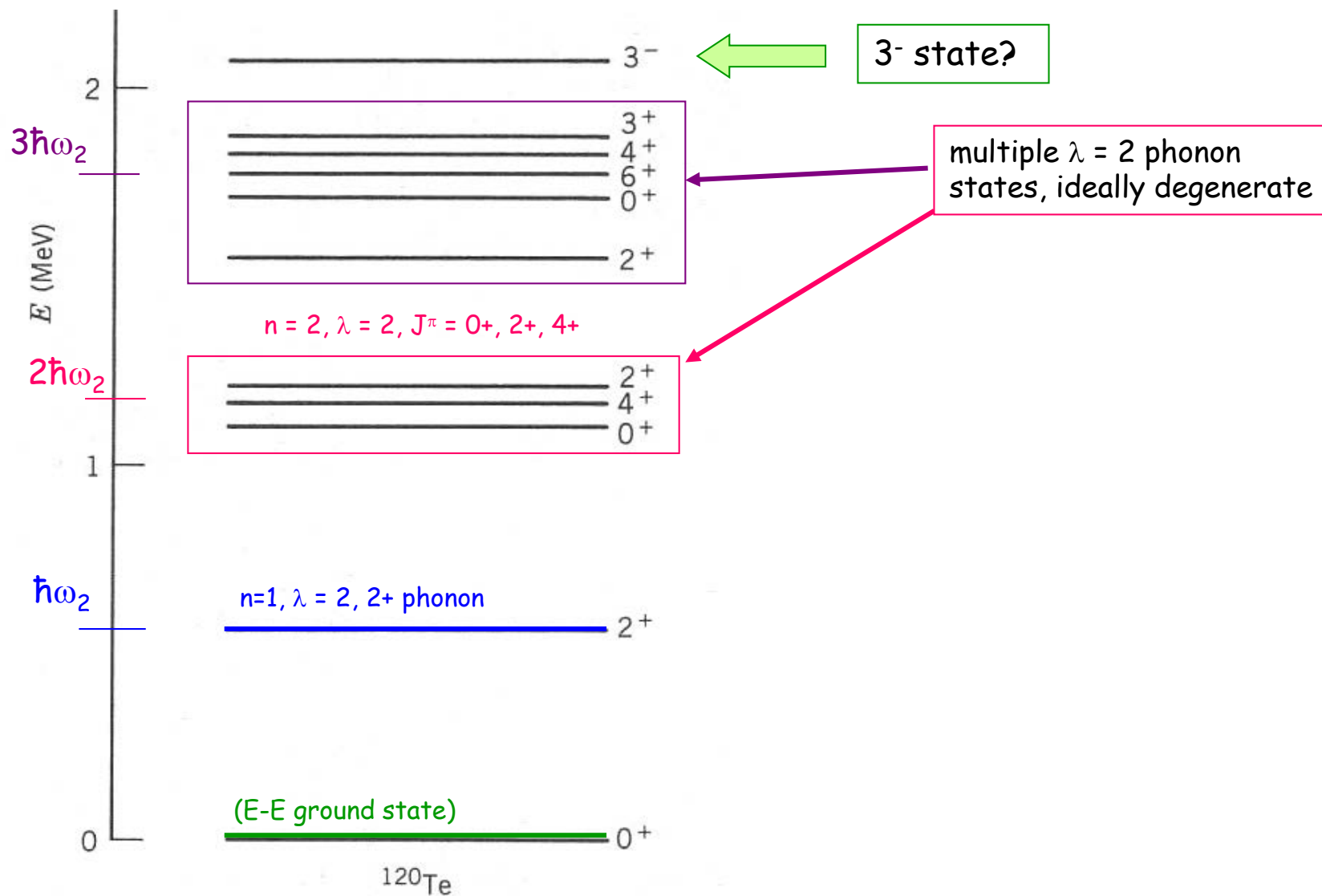
$$E_2 = \hbar\omega \cdot \left(2 + \frac{5}{2}\right)$$

$$E_1 = \hbar\omega \cdot \left(1 + \frac{5}{2}\right)$$

$$E_0 = \hbar\omega \cdot \frac{5}{2}$$



Example of vibrational excitations



Second quantization

Hamilton operator $H = \sum_{\lambda\mu} \hbar \cdot \omega_{\lambda} \cdot \left(\beta_{\lambda\mu}^{\dagger} \beta_{\lambda\mu} + \frac{1}{2} \right)$

rule for boson operators: $\beta_{\lambda\mu} \cdot \beta_{\lambda'\mu'}^{\dagger} - \beta_{\lambda\mu}^{\dagger} \cdot \beta_{\lambda'\mu'} = \delta_{\lambda\lambda'} \cdot \delta_{\mu\mu'}$

2-phonon state: $\Psi_{IM} = N \sum_{m_1 m_2} (l_1 l_2 m_1 m_2 | IM) \cdot \beta_{l_1 m_1}^{\dagger} \cdot \beta_{l_2 m_2}^{\dagger} | 0 \rangle$

reduced transition probability: $B(EI; i \rightarrow f) = \sum_{M_f m} \left| \langle f | M(l, m) | i \rangle \right|^2$

$$M(l, m) = \int \rho_p \cdot r^l \cdot Y_{lm}^*(\theta, \phi) d\tau \cong \frac{3 \cdot Z}{4 \cdot \pi \cdot R_0^3} \cdot R_0^{l+3} \cdot \alpha_{lm}^*$$

$$B(E2; 2 \rightarrow 0) = \left[\frac{3 \cdot Z \cdot e \cdot R_0^2}{4 \cdot \pi} \right]^2 \cdot \sum_{M_f m} \left| \left\langle 0 \left| \left\{ \sqrt{\frac{\hbar}{2 \cdot B_2 \cdot \omega_2}} \cdot (-1)^m \cdot [\beta_{2-m}^{\dagger} + (-1)^m \beta_{2m}] \right\} \beta_{2M_i}^{\dagger} \right| 0 \right\rangle \right|^2$$

$$B(E2; 2 \rightarrow 0) = \left[\frac{3 \cdot Z \cdot e \cdot R_0^2}{4 \cdot \pi} \right]^2 \cdot \frac{\hbar}{2 \cdot B_2 \cdot \omega_2} \cdot \boxed{\sim \frac{1}{E_{2^+}}}$$

Reduced transition probabilities

2-phonon state $B(E2; n_2 = 2 \rightarrow n_2 = 1) = 2 \cdot B(E2; n_2 = 1 \rightarrow n_2 = 0)$

3-phonon state $B(E2; n_2 = 3 \rightarrow n_2 = 2) = 3 \cdot B(E2; n_2 = 1 \rightarrow n_2 = 0)$

$$B(E2; I_i \rightarrow I_f) = \frac{1}{2 \cdot I_i + 1} \left| \langle I_f \| M(E2) \| I_i \rangle \right|^2$$

$$Q_{vib} = \frac{3 \cdot Z \cdot R_0^2}{4 \cdot \pi} \cdot \sqrt{\frac{\hbar}{2 \cdot B_2 \cdot \omega_2}}$$

Appendix: Reduced transition probabilities

1-phonon state $\langle I = 2, n_2 = 1 \| M(E2) \| I = 0, n_2 = 0 \rangle = \sqrt{5} \cdot Q_{vib} \cdot e$

2-phonon state $\langle I = 4, n_2 = 2 \| M(E2) \| I = 2, n_2 = 1 \rangle = \sqrt{18} \cdot Q_{vib} \cdot e$

$$\langle I = 2, n_2 = 2 \| M(E2) \| I = 2, n_2 = 1 \rangle = \sqrt{10} \cdot Q_{vib} \cdot e$$

$$\langle I = 0, n_2 = 2 \| M(E2) \| I = 2, n_2 = 1 \rangle = \sqrt{2} \cdot Q_{vib} \cdot e$$

3-phonon state $\langle I = 6, n_2 = 3 \| M(E2) \| I = 4, n_2 = 2 \rangle = \sqrt{39} \cdot Q_{vib} \cdot e$

$$\langle I = 4, n_2 = 3 \| M(E2) \| I = 4, n_2 = 2 \rangle = \sqrt{\frac{90}{7}} \cdot Q_{vib} \cdot e$$

$$\langle I = 4, n_2 = 3 \| M(E2) \| I = 2, n_2 = 2 \rangle = \sqrt{\frac{99}{7}} \cdot Q_{vib} \cdot e$$

$$\langle I = 3, n_2 = 3 \| M(E2) \| I = 4, n_2 = 2 \rangle = \sqrt{6} \cdot Q_{vib} \cdot e$$

$$\langle I = 3, n_2 = 3 \| M(E2) \| I = 2, n_2 = 2 \rangle = -\sqrt{15} \cdot Q_{vib} \cdot e$$

$$\langle I = 2, n_2 = 3 \| M(E2) \| I = 2, n_2 = 2 \rangle = \sqrt{\frac{20}{7}} \cdot Q_{vib} \cdot e$$

$$\langle I = 2, n_2 = 3 \| M(E2) \| I = 0, n_2 = 2 \rangle = \sqrt{7} \cdot Q_{vib} \cdot e$$

$$\langle I = 0, n_2 = 3 \| M(E2) \| I = 0, n_2 = 2 \rangle = \sqrt{3} \cdot Q_{vib} \cdot e$$

$$B(E2; I_i \rightarrow I_f) = \frac{1}{2 \cdot I_i + 1} \langle I_f \| M(E2) \| I_i \rangle^2$$

$$Q_{vib} = \frac{3 \cdot Z \cdot R_0^2}{4 \cdot \pi} \cdot \sqrt{\frac{\hbar}{2 \cdot B_2 \cdot \omega_2}}$$