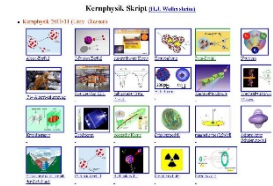


Outline: Experimental Nuclear Astrophysics

Lecturer: Hans-Jürgen Wollersheim

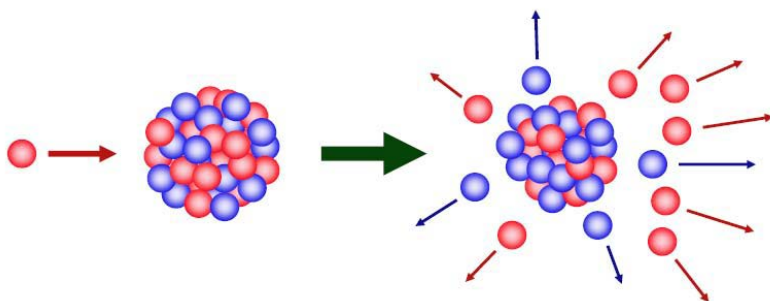
e-mail: h.j.wollersheim@gsi.de

web-page: <https://web-docs.gsi.de/~wolle/> and click on



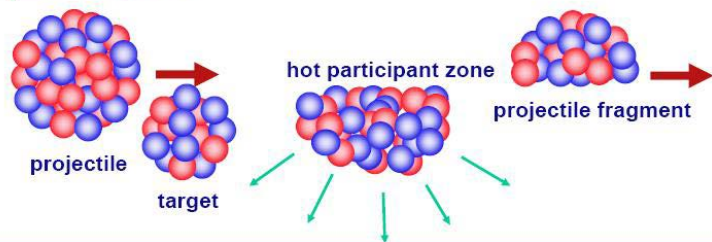
1. Projectile fragmentation
2. Isotope separation on line
3. Waiting point nuclei
4. reactions, masses, radii

Radioactive Ion Beams production methods

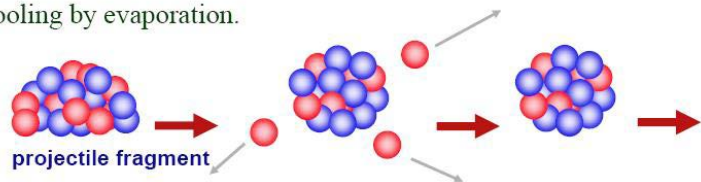


Target fragmentation

Random removal of protons and neutrons from heavy projectile in peripheral collisions



Cooling by evaporation.



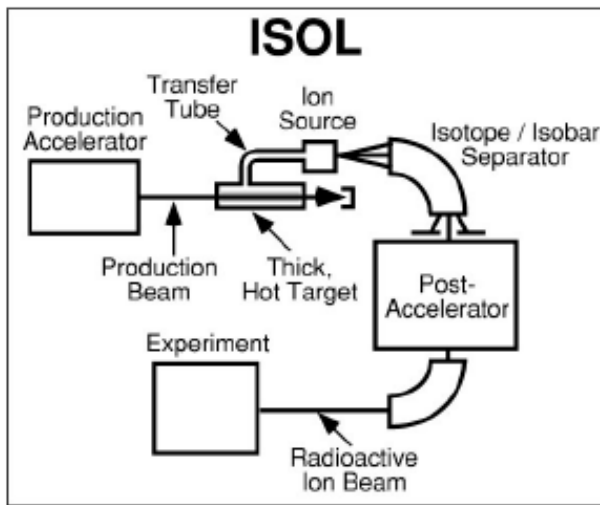
Projectile fragmentation

fragmentation invented at LBNL in the 1980's



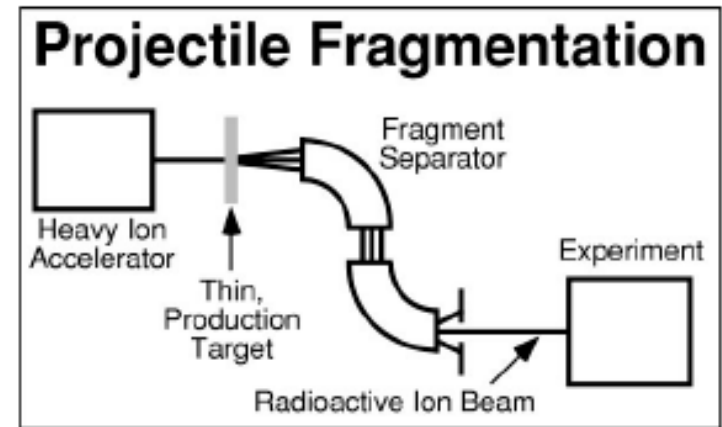
Radioactive Ion Beams production methods

- Isotope Separation on Line (ISOL) (CERN, LLN, ORNL, TRIUMF)
- Projectile Fragmentation (PF) (GANIL, GSI, MSU, RIKEN)
- in-flight production (ANL, Notre Dame, TAMU)
- batch mode production (suitable for long-lived species)
- ...



- 😊 excellent quality
- high purity
- high intensities

- 😞 limited number of species
- different production for different species
- limited to nuclei with $t_{1/2} \geq 1s$ (allow for diffusion)



- 😊 independent from chemical properties
- no limitations on $t_{1/2}$ (fast separation)

- 😞 typical beam energies too high for NA
- poorer beam quality (energy, size)
- possible beam contaminations

M.S. Smith and K.E. Rehm, Ann. Rev. Nucl. Part. Sci, 51 (2001) 91-130

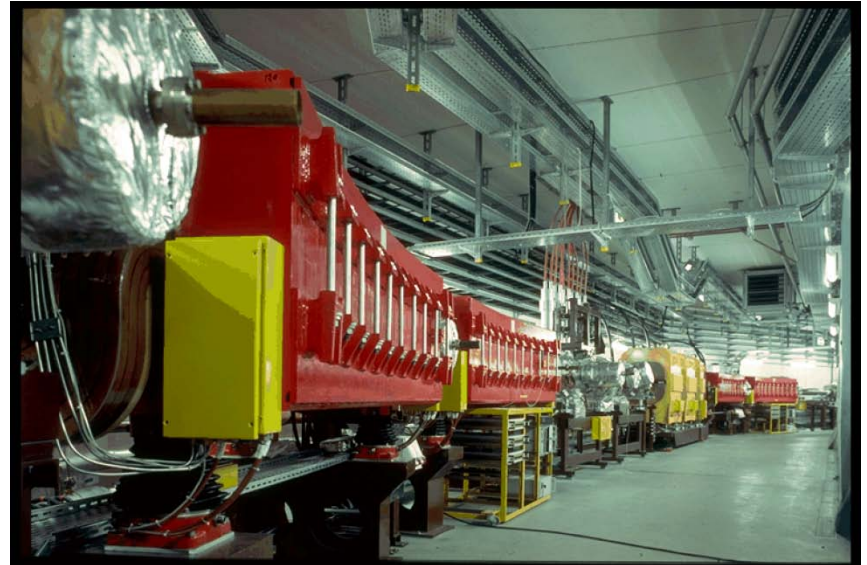
Fast radioactive beams – why?

- Production by the in-flight technique (vs ISOL)
 - Chemical blindness
 - Shortest half-lives
- Experimental issues
 - Doppler boosted
 - Kinematical focusing
 - **Resolution**
 - Small energy loss
 - Possible to use thick targets
 - **Ions, particles are hard to stop**
 - Can use many detectors
 - ion-by-ion tracking

Ranges in silicon	p	⁴⁰ Ar	²³⁸ U
10 MeV/u (ISOL + post-acc.)	700 μm	127 μm	107 μm
500 MeV/u (Relativistic in-flight)	63 cm	7.6 cm	1.8 cm

Fast radioactive beams – where?

- GSI since SIS (early 90's)
- Intermediate-energy RIBs (tens of MeV/u) since many years at GANIL, MSU, RIKEN
- Future (and current) facilities
 - RIBF@RIKEN
 - FRIB@MSU
 - FAIR-NuSTAR



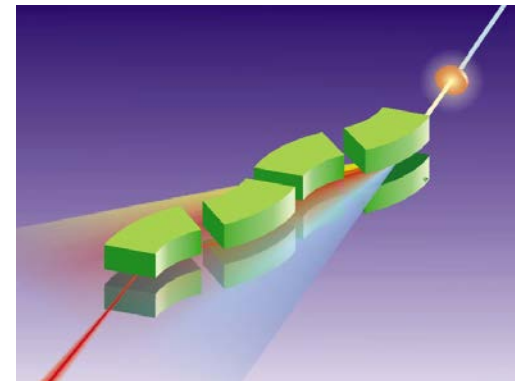
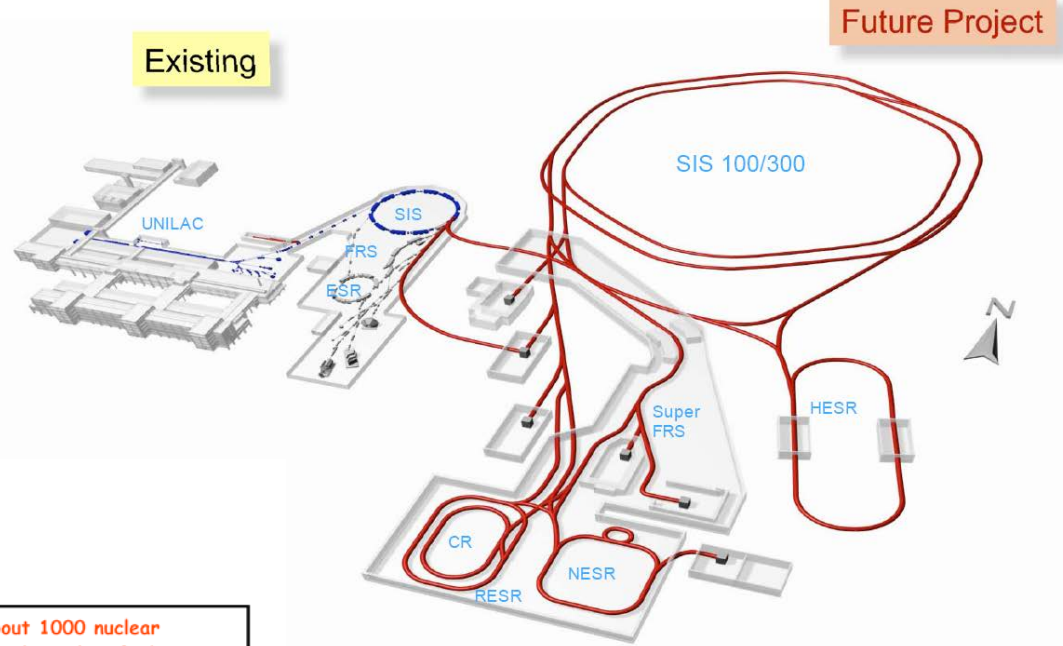
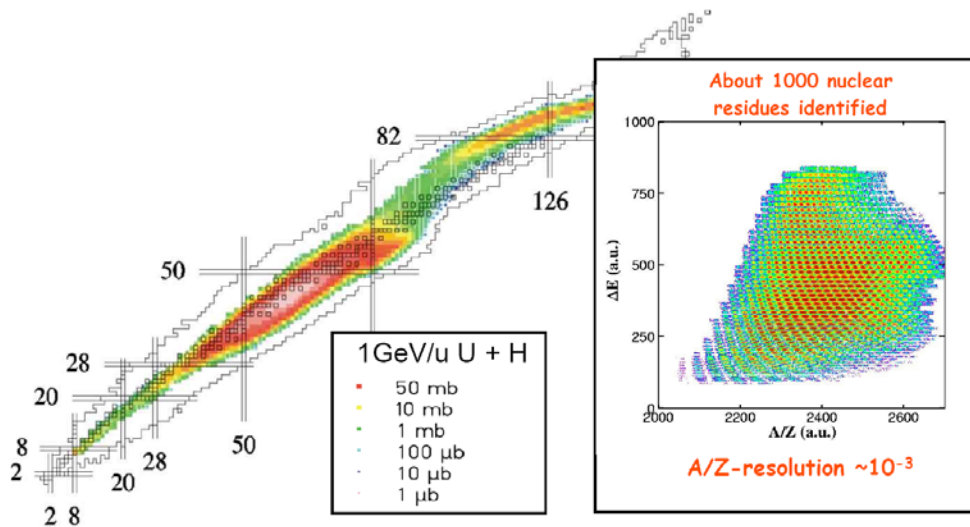
The future - FAIR

Next generation
RIBs facilities

The Future International Facility at GSI:
FAIR – Facility for Antiproton and Ion Research

1.0-1.5 GeV/u
all elements up to uranium

P. Armbruster et al.; Phys. Rev. Letters, Jan. 05



GSI



Nuclear reactions

- Relativistic energies *R3B*
- Cooled beams *EXL*
- High-res. spectroscopy *HISPEC*

Decay properties

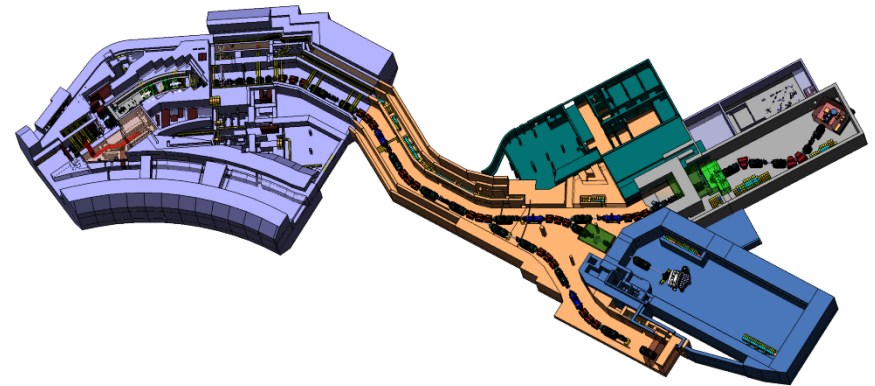
- Stopped beams *DESPEC*

Ground state properties

- Masses *MATS, ILIMA*
- Radii, momenta *LASPEC*

New tools

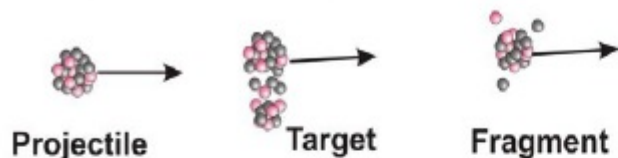
- Electron – RIB scattering *ELISe*
- p-bar – RIB collider *AIC*



super FRagment Separator

Production of exotic nuclei at relativistic energies

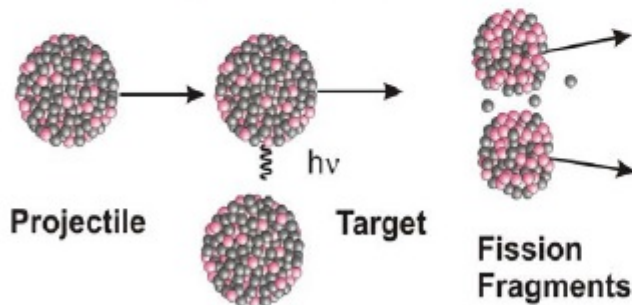
Projectile Fragmentation



Nucleon-nucleon collisions, abrasion, ablation

$$\vec{V}_f \approx \vec{V}_p$$

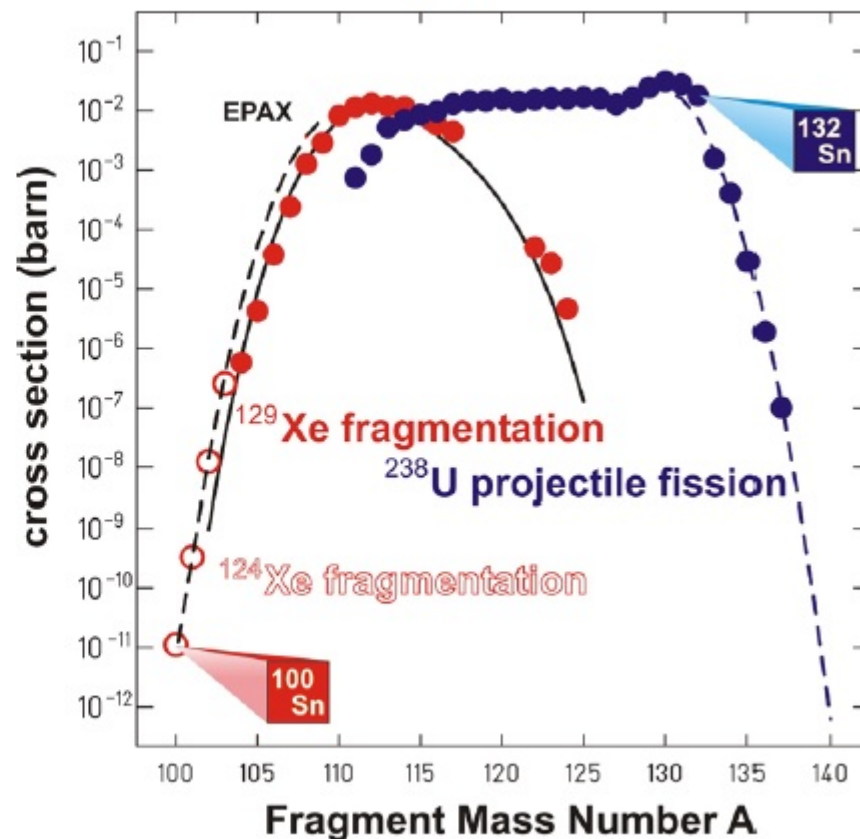
Projectile Fission



Electromagnetic excitation, fission in flight

$$\vec{V}_f \approx \vec{V}_p + \vec{V}_{fission}$$

Sn isotope production



K. Sümmerner

Production of exotic nuclei at relativistic energies

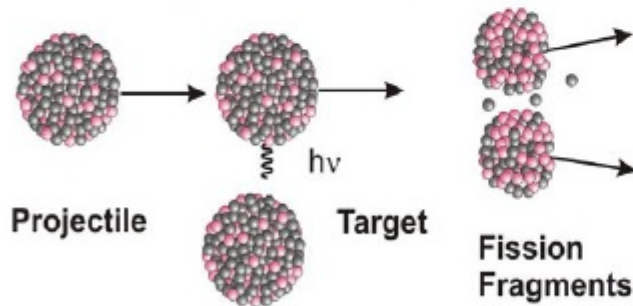
Projectile Fragmentation



Nucleon-nucleon collisions, abrasion, ablation

$$\vec{V}_f \approx \vec{V}_p$$

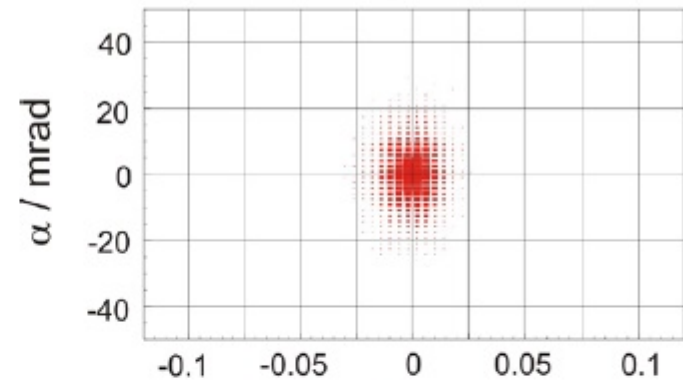
Projectile Fission



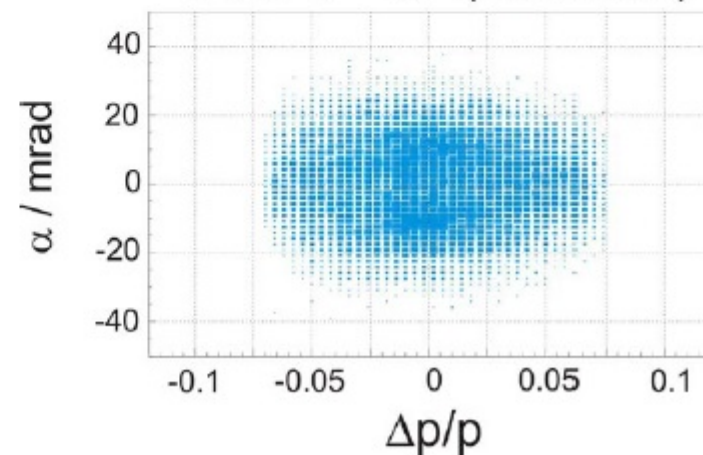
Electromagnetic excitation, fission in flight

$$\vec{V}_f \approx \vec{V}_p + \vec{V}_{fission}$$

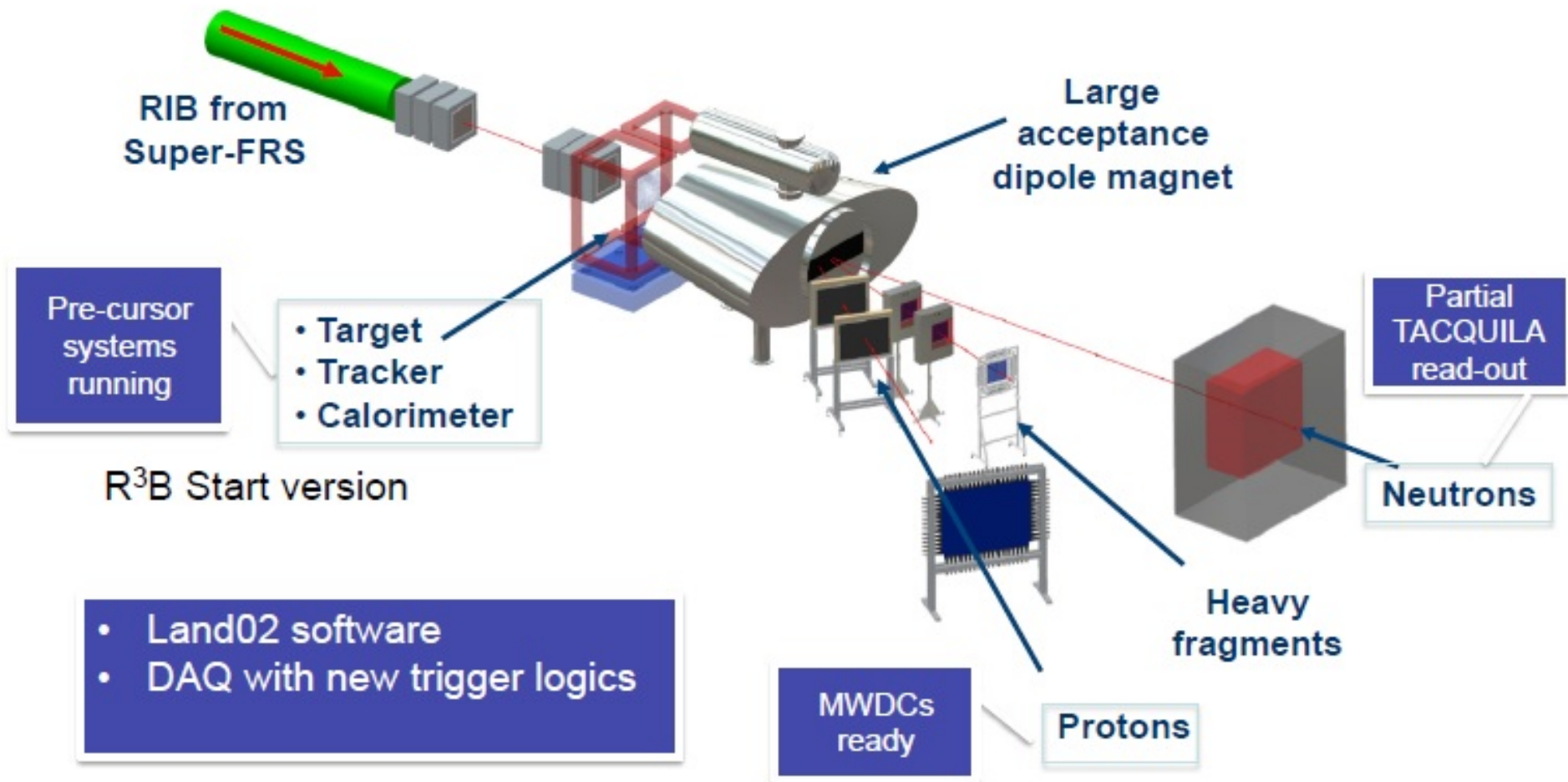
$^{124}\text{Xe} + \text{C} \rightarrow ^{100}\text{Sn}$ - (from Fragmentation)



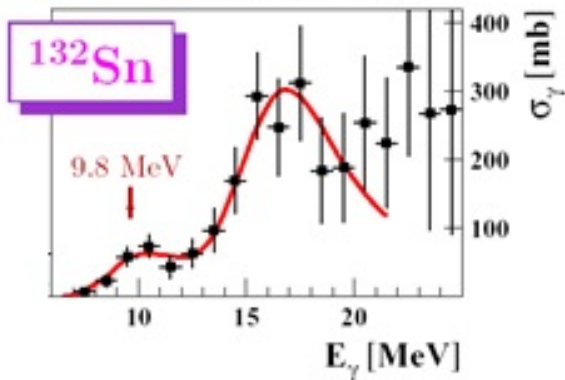
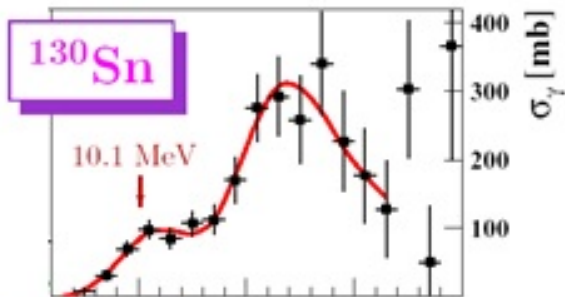
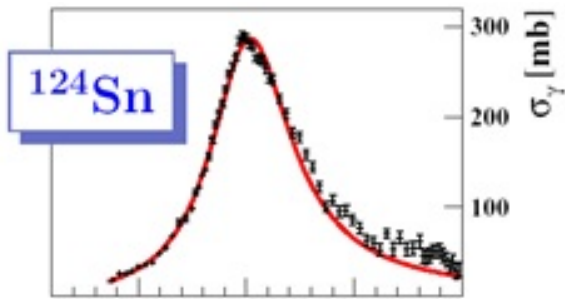
$^{238}\text{U} + \text{C} \rightarrow ^{132}\text{Sn}$ - (from Fission)



Reactions with Relativistic Radioactive Beams



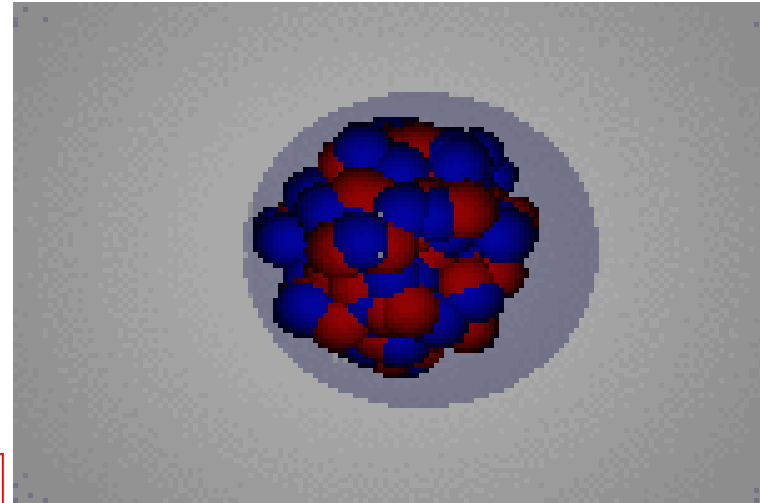
Pygmy resonances and the neutron skin



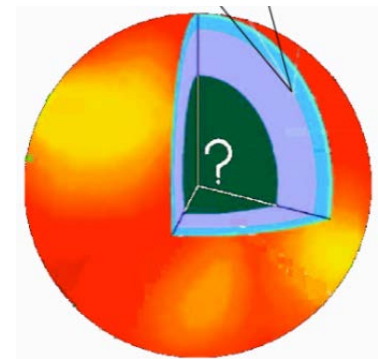
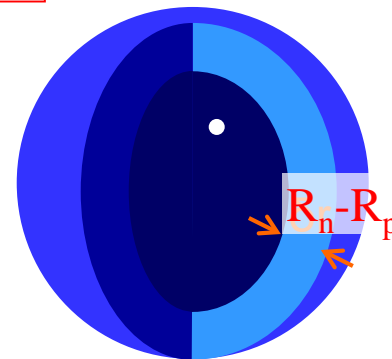
$R_n - R_p :$

$^{130}\text{Sn} : 0.23 \pm 0.04 \text{ fm}$

$^{132}\text{Sn} : 0.24 \pm 0.04 \text{ fm}$



neutron star



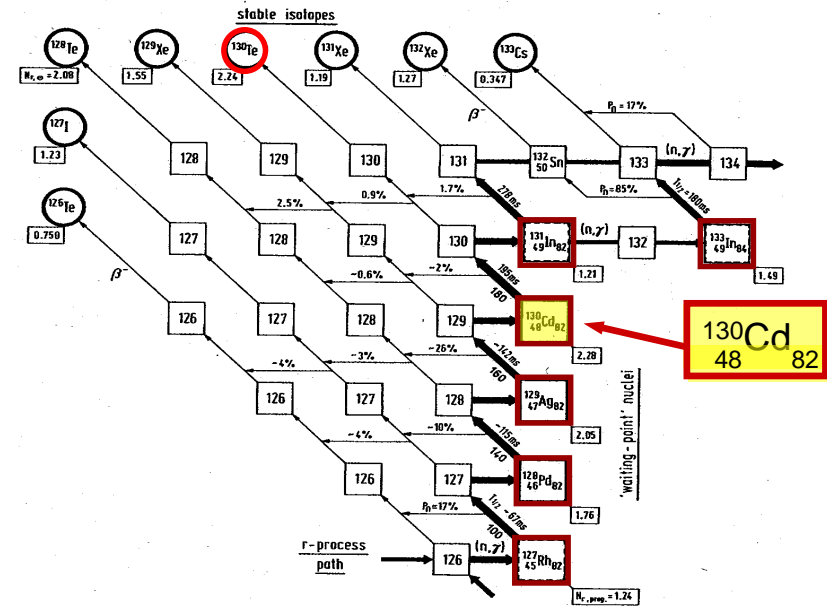
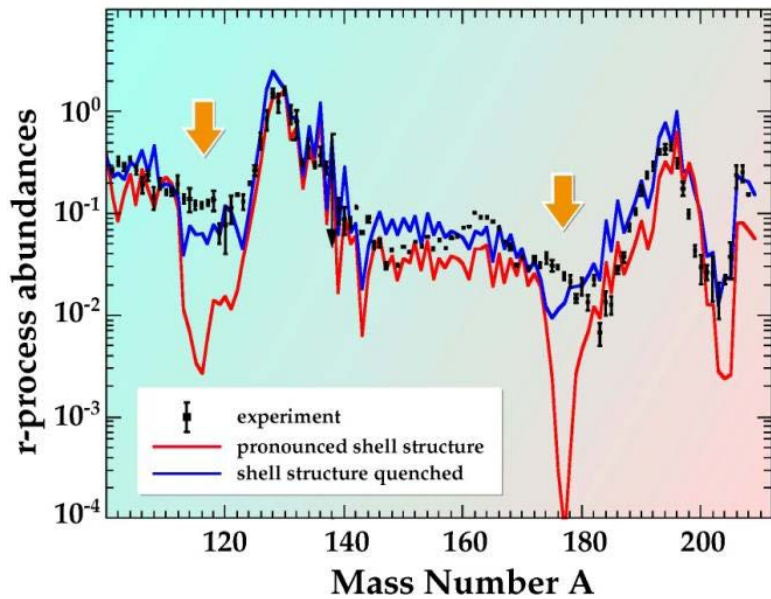
EoS

at LAND-GSI: $10^8/\text{s}$ ^{238}U beam (10 d)

at R3B-FAIR: 10^{12} ^{238}U beam (100 s)

P. Adrich et al., PRL 95 (2005) 132501 & A.Klimkiewicz et al., PRC 2008

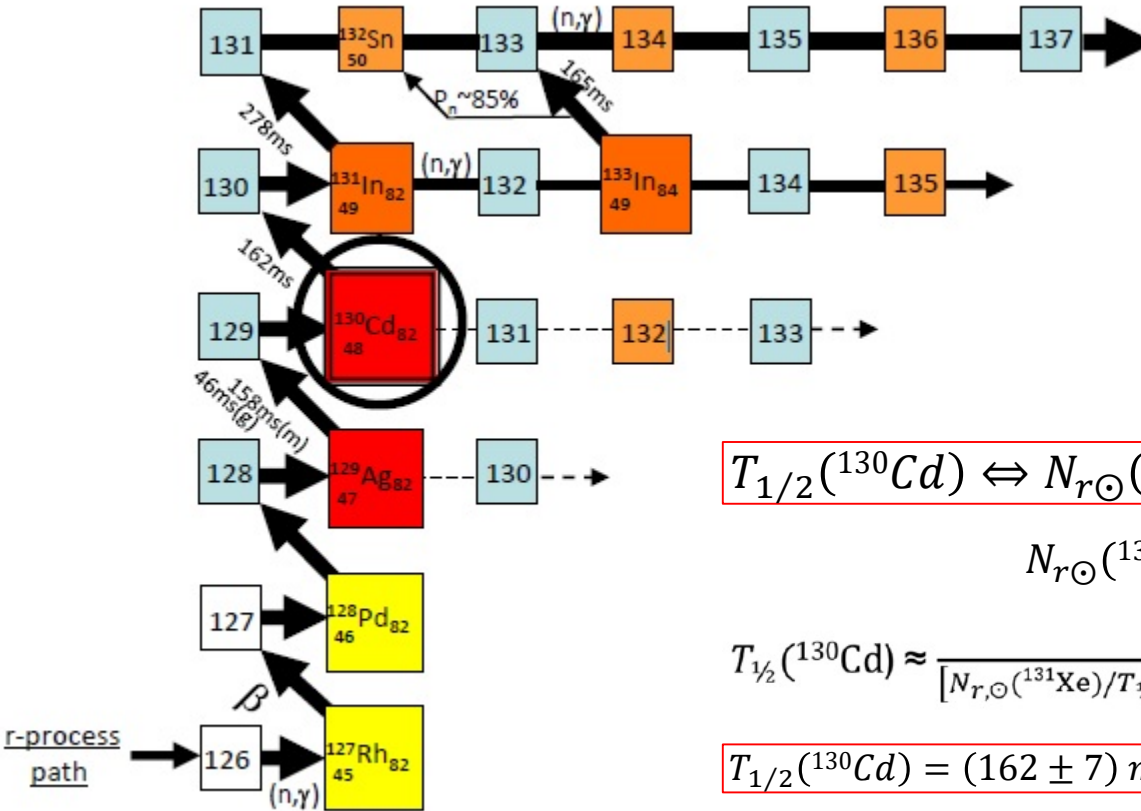
Waiting point at N=82



“..the calculated r-abundance ‘hole‘ in the $A \cong 120$ region reflects ... **the weakening of the shell strength** ... below ^{132}Sn “
 K-L Kratz

bottleneck at N=82 waiting point near stability?

^{130}Cd – the key isotope at the A=130 peak



$$T_{1/2}(^{130}\text{Cd}) \Leftrightarrow N_{r\odot}(^{130}\text{Te})$$

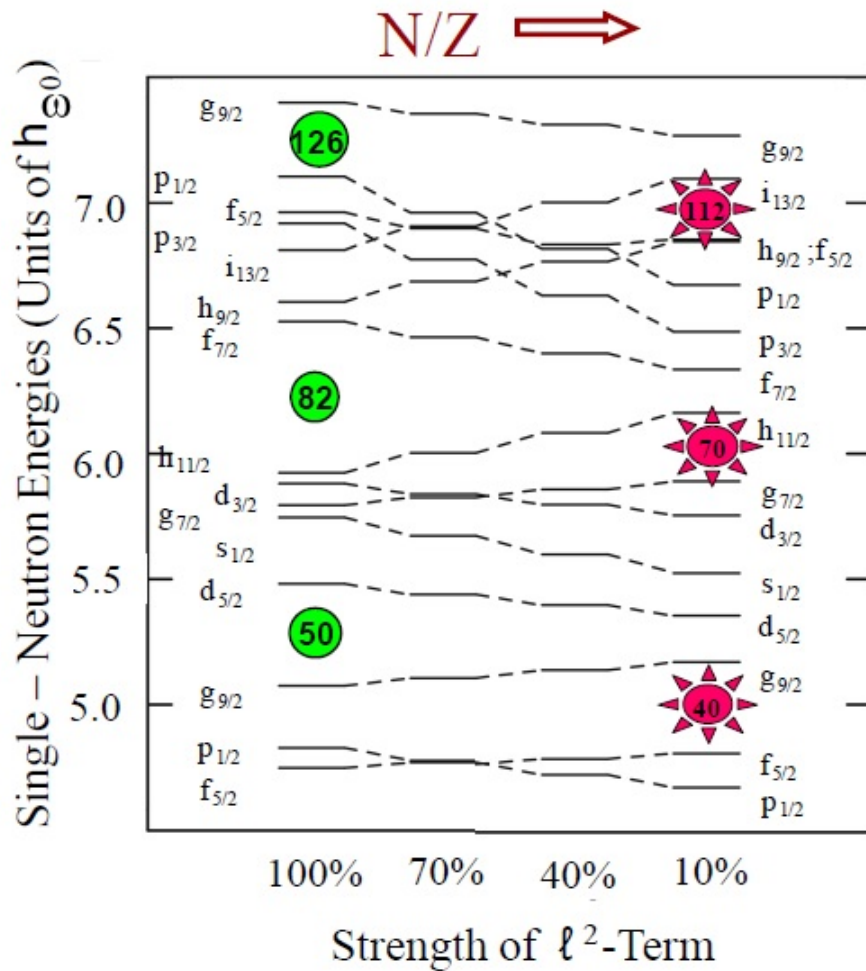
$N_{r\odot}(^{130}\text{Te}) \equiv$ solar system abundances ^{130}Te

$$T_{1/2}(^{130}\text{Cd}) \approx \frac{N_{r\odot}(^{130}\text{Te})}{[N_{r\odot}(^{131}\text{Xe})/T_{1/2}(^{131}\text{In})] + [1.1N_{r\odot}(^{132}\text{Xe})/T_{1/2}(^{133}\text{In})]} \approx 170\text{ms}$$

$$T_{1/2}(^{130}\text{Cd}) = (162 \pm 7)\text{ms}$$

climb up the N=82 ladder
A~130 “bottle neck”

Effects of N=82 “shell quenching”

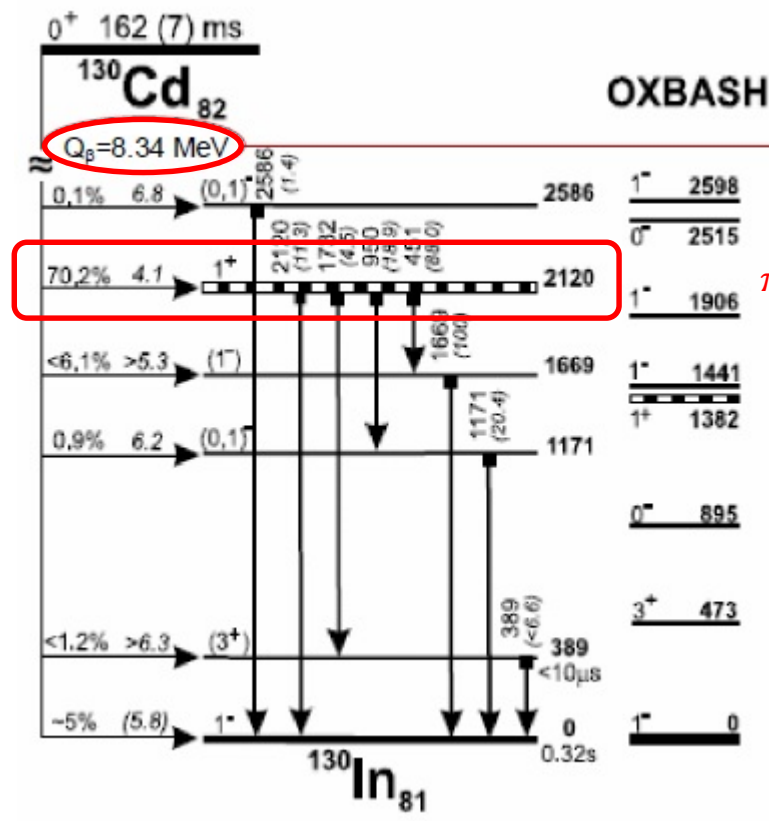
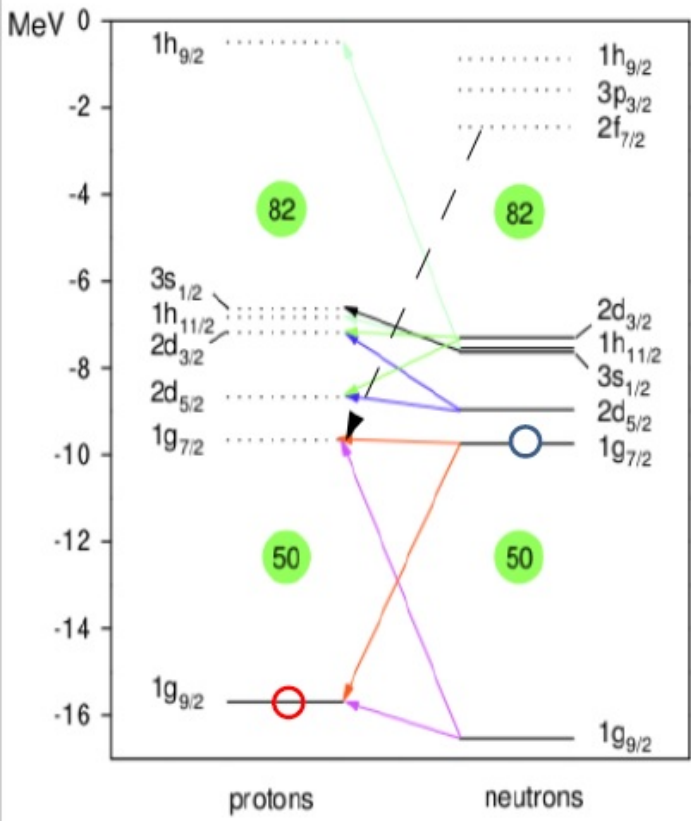


“Shell quenching”

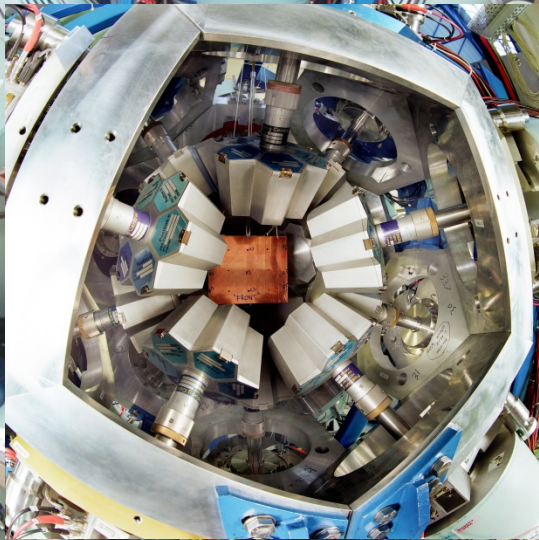
reduction of the spin-orbit coupling strength;
 caused by strong interaction between bound
 and continuum states;
 due to diffuseness of “neutron-skin” and its
 influence on the central potential

^{130}Cd decay spectroscopy

Large Q_β value
best reproduced
by mass models
with $N=82$ shell
quenching



$\pi g_{9/2} \otimes \nu g_{7/2}$



scintillator
(SC41)

ionization
chambers
(MUSIC41,42)

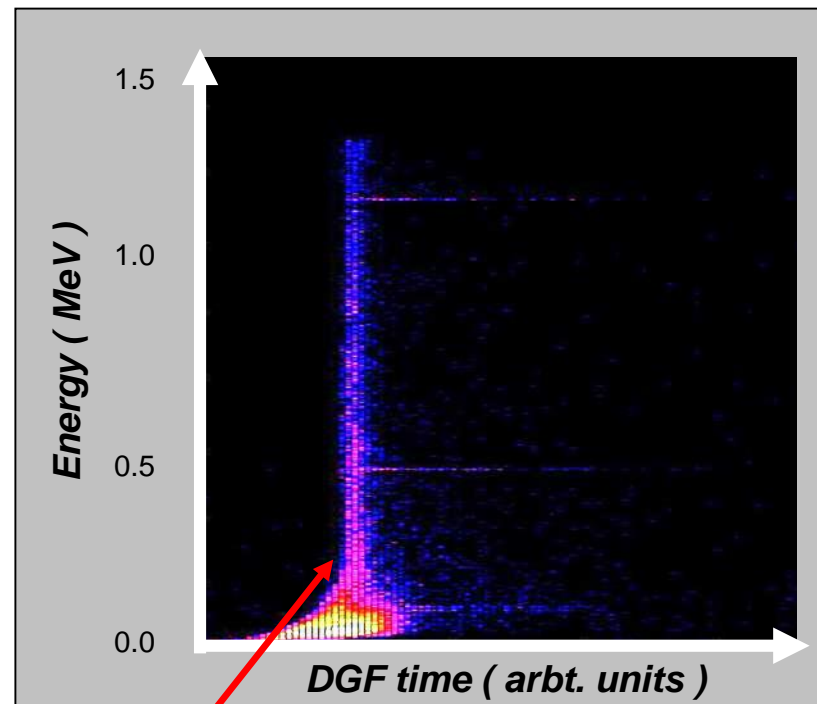
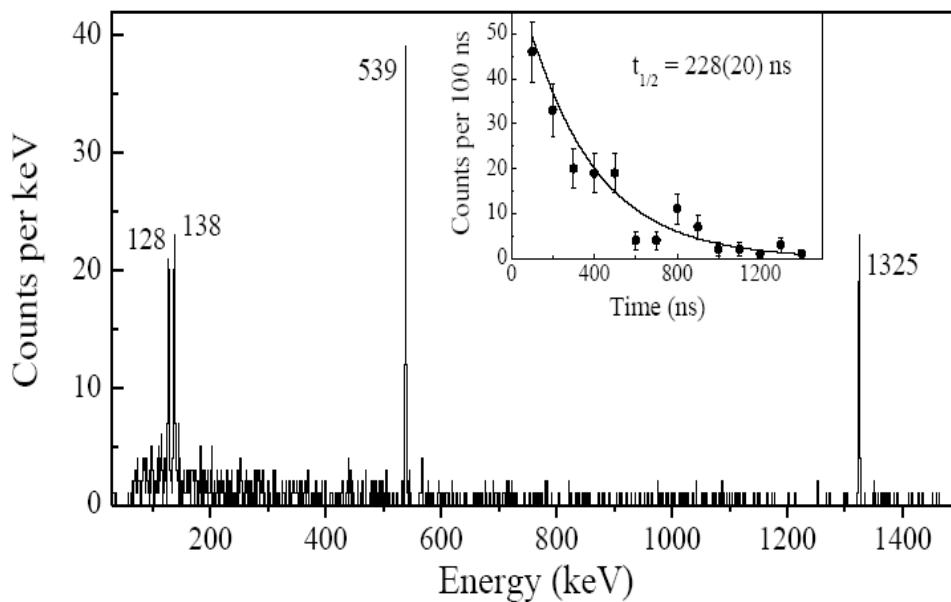
beam

degrader

multiwire
chambers
(MW41,MW42)

Decay spectroscopy probes shell closures

^{130}Cd : DGF-timing



Prompt γ -flash

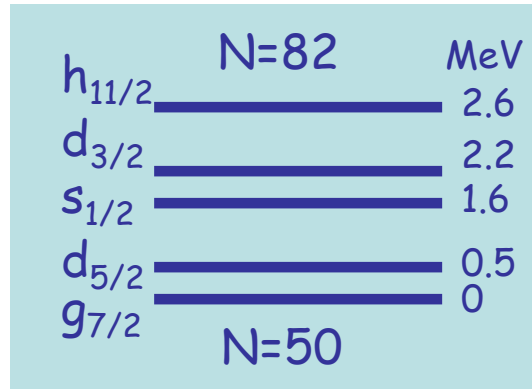
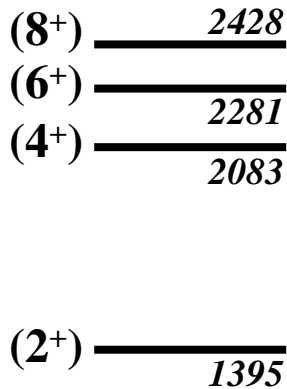
Decay time range: 20 ns ... 20 μs

$8^+(g_{9/2})^{-2}$ seniority isomers in ^{98}Cd and ^{130}Cd

Sn100 0.94 s 0+	Sn101 3 s 0+	Sn102 4.5 s 0+	Sn103 7 s 0+	Sn104 20.8 s 0+	Sn105 31 s 0+	Sn106 115 s 0+	Sn107 2.90 m (5/2+)	Sn108 10.50 m 0+	Sn109 18.8 m 5/2(+)	Sn110 411 s 0+	Sn111 35.3 m 7/2+	Sn112 0+	Sn113 115.09 d 1/2+	Sn114 0+	Sn115 1/2+	Sn116 0+	Sn117 1/2+	Sn118 0+	Sn119 1/2+	Sn120 0+	Sn121 17.06 h 3/2+	Sn122 0+	Sn123 129.2 d 11/2+	Sn124 0+	Sn125 9.64 d 11/2+	Sn126 1E+5 y 0+	Sn127 2.10 h (11/2-)	Sn128 59.07 m 0+	Sn129 2.15 m (3/2-)	Sn130 3.72 m 0+	Sn131 56.9 s (3/2-)	Sn132 39.7 s 0+
In99 EC	In100 7.8 s EC	In101 15.1 s EC	In102 21 s (6+)	In103 6 s (9/2+)	In104 1.30 m (6+)	In105 5.97 m (9/2+)	In106 6.2 m 7+	In107 32.4 m 9/2+	In108 58.9 m 7+	In109 4.1 h 9/2+	In110 4.9 h 7+	In111 2.8047 d 9/2+	In112 14.97 m 1+	In113 71.9 s 1+	In114 4.01E+5 y 1+	In115 14.15 s 9/2+	In116 43.3 m 9/2+	In117 5.9 s 9/2+	In118 2.4 m 1+	In119 3.38 s 9/2+	In120 33.1 s 9/2+	In121 5.9 s 9/2+	In122 15 s 1+	In123 5.9 s 9/2+	In124 1.11 s 3+	In125 2.36 s 9/2(+)	In126 1.60 s 3(-)	In127 1.59 s (9/2+)	In128 0.34 s (3+)	In129 0.61 s (9/2+)	In130 0.32 s 1(-)	In131 0.232 s (9/2+)
Cd98 9.2 s (5/2+)	Cd99 18 s (5/2+)	Cd100 49.1 s 0+	Cd101 1.36 m (5/2+)	Cd102 2.5 m 0+	Cd103 7.3 m (5/2+)	Cd104 57.7 m 0+	Cd105 35.5 m 5/2+	Cd106 461.6 d 0+	Cd107 4.59 h 5/2+	Cd108 0+	Cd109 461.6 d 5/2+	Cd110 12.49 s 0+	Cd111 12.89 s 1/2+	Cd112 24.13 s 0+	Cd113 2.9E+15 y 1/2+	Cd114 28.73 s 0+	Cd115 53.46 s 1/2+	Cd116 50.3 m 0+	Cd117 2.49 h 1/2+	Cd118 59.3 m 0+	Cd119 1.69 m 3/2+	Cd120 59.89 s 0+	Cd121 13.5 s (3/2+)	Cd122 2.34 s 0+	Cd123 2.10 s (3/2+)	Cd124 1.15 s 0+	Cd125 0.65 s (3/2+)	Cd126 0.596 s 0+	Cd127 0.37 s (3/2+)	Cd128 0.34 s 0+	Cd129 0.25 s (3/2+)	Cd130 0.20 s 0+

Cd98
9.2 s
0+
EC

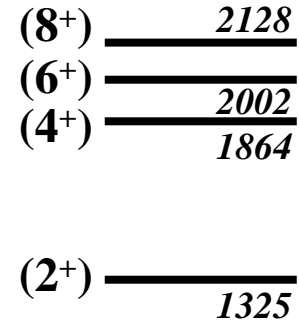
N=50
Z=48



participating neutron-orbitals

Cd130
0.20 s
0+
β-n

N=82
Z=48



two proton holes in the $g_{9/2}$ orbit

No dramatic shell quenching!



Surprising β -decay properties of ^{131}Cd

...just **ONE** neutron outside N=82 magic shell

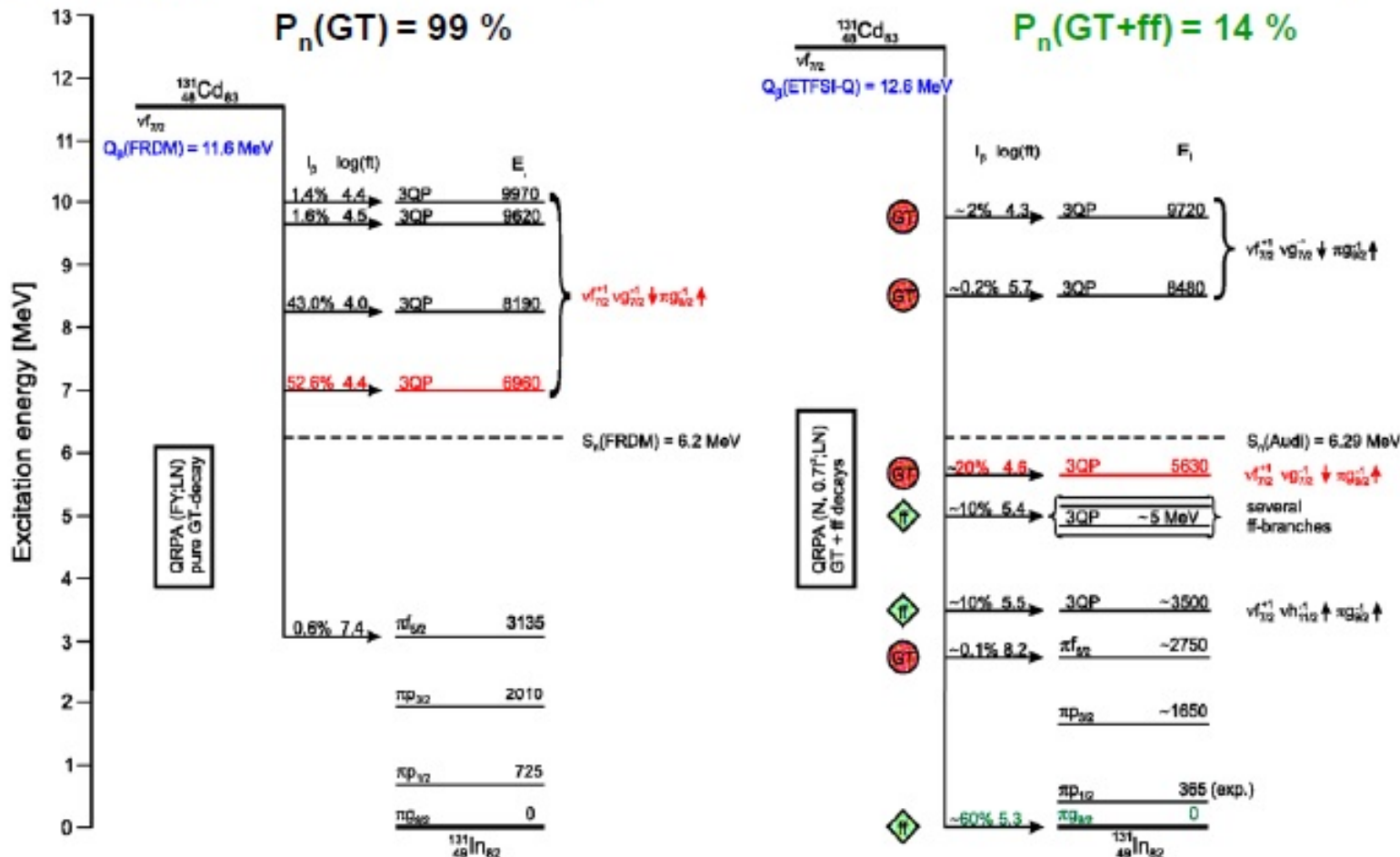
Experiment: $T_{1/2} = 68$ ms; $P_n = 3.4$ %

QRPA predictions: $T_{1/2}(\text{GT}) = 943$ ms;

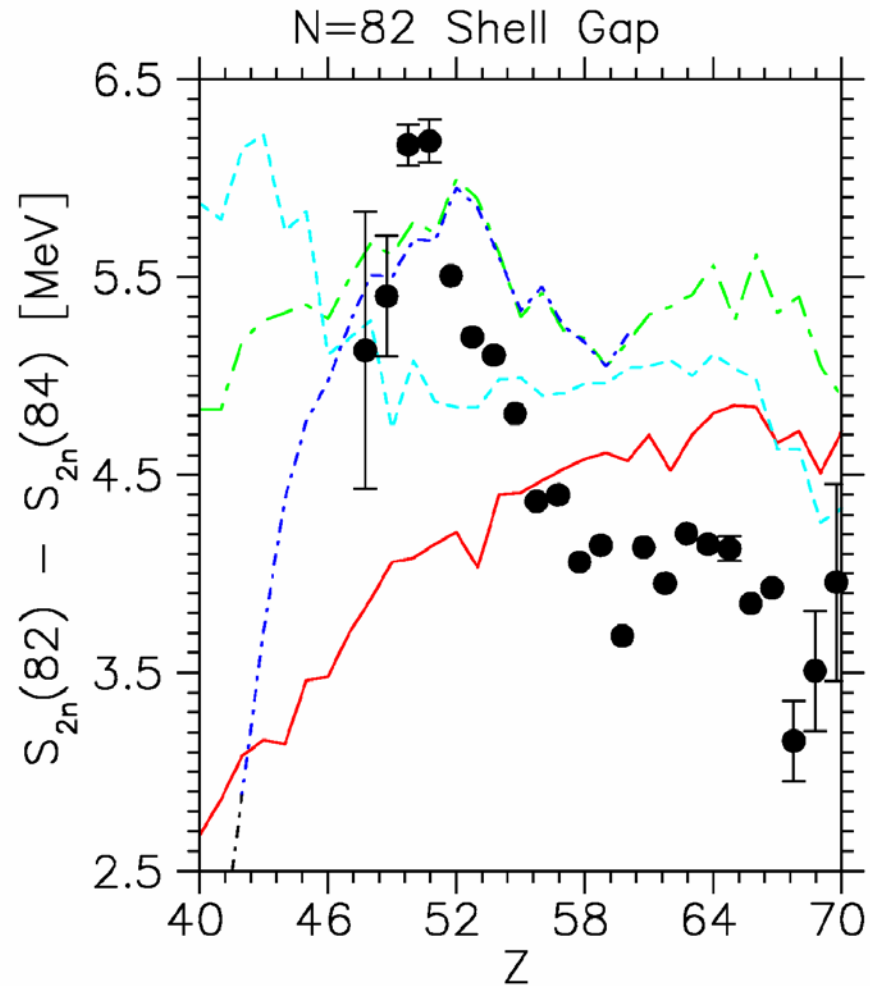
$P_n(\text{GT}) = 99$ %

$T_{1/2}(\text{GT+ff}) = 59$ ms;

$P_n(\text{GT+ff}) = 14$ %

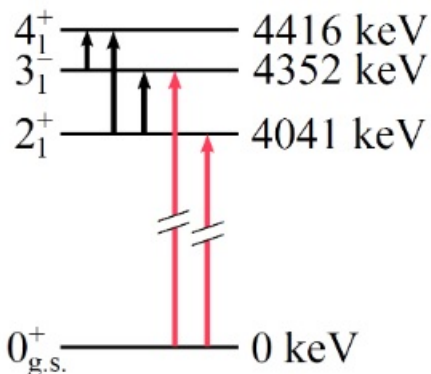
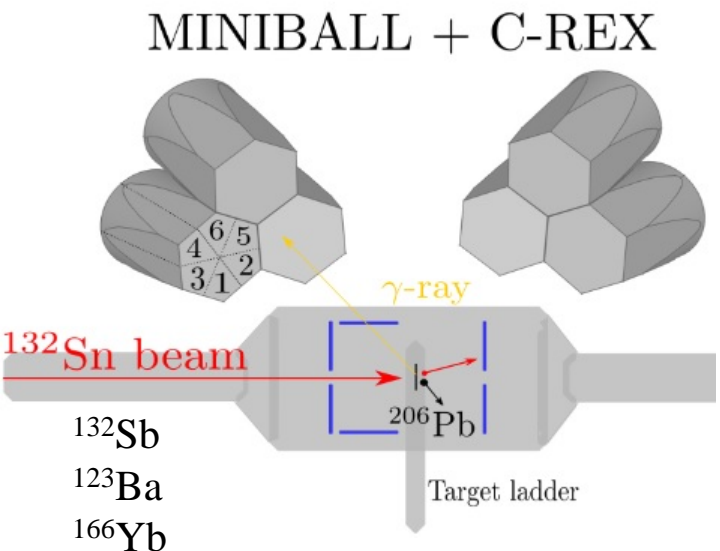


The N=82 shell gap

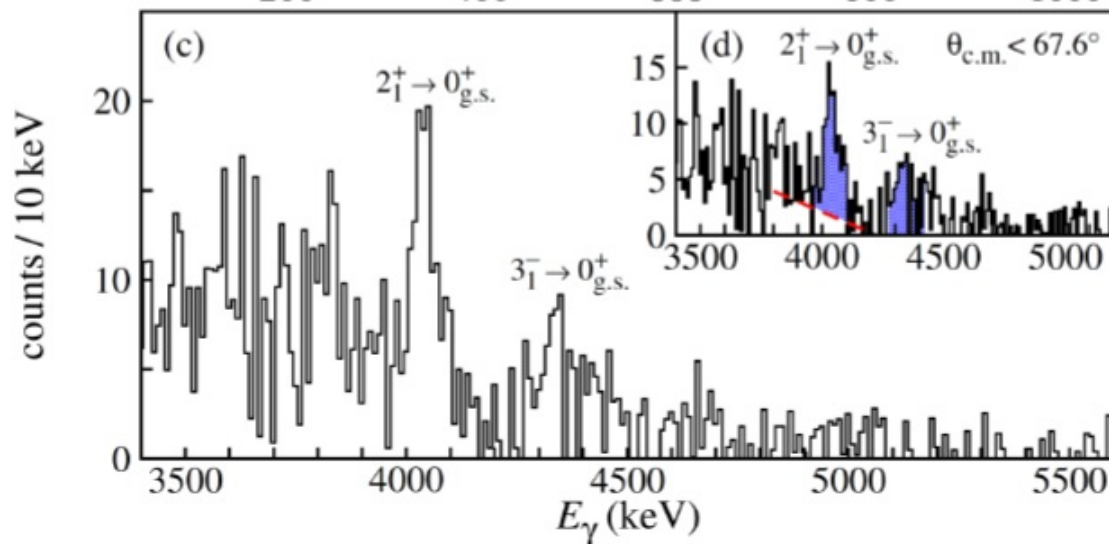
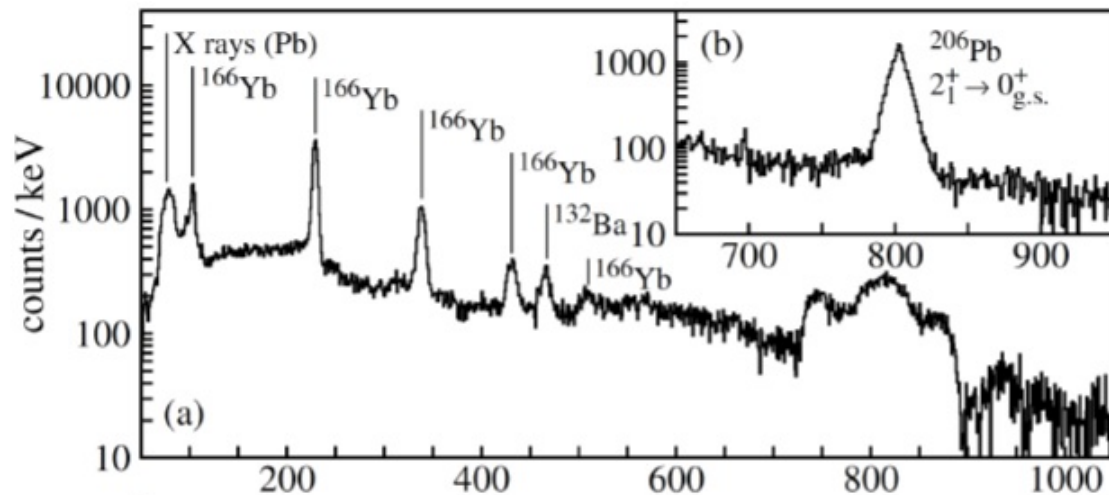


NONE of the mass models predicts the trend correctly!

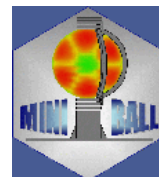
Enhanced quadrupole and octupole strength in doubly-magic ^{132}Sn



$B(E2; 0^+ \rightarrow 2^+)$	$0.087(19) e^2b^2$
$B(E3; 0^+ \rightarrow 3^-)$	$0.11(4) e^2b^3$
$B(E1; 2^+ \rightarrow 3^-)$	$9.1(31) \times 10^{-6} e^2b$



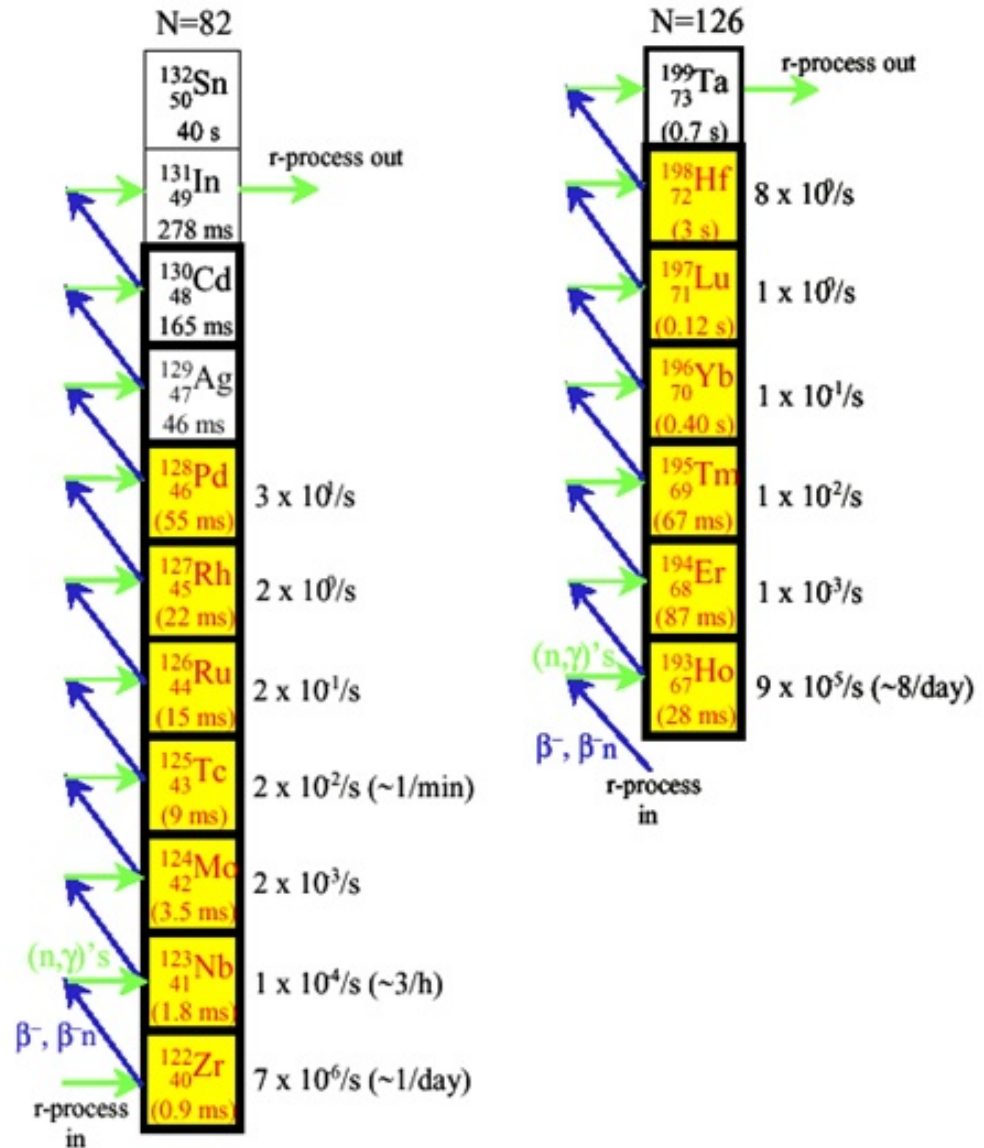
D. Rosiak et al., PRL 121, 252501 (2018)



Waiting point lifetimes with fragmentation facilities

Why fragmentation?

- Lifetime measurements can be done with beams from low energy facilities
- But, fragmentation facilities have advantages:
 - Use beams of mixed nuclides-- Identification on event by event basis
 - Greater reach toward dripline
 - FAIR, GSI, RIKEN, NSCL and RIA will cover a large part of r-process path

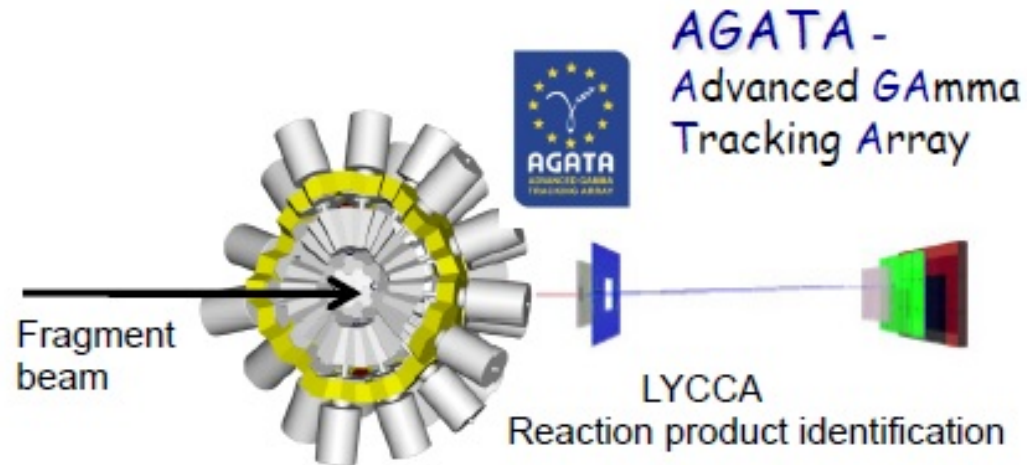
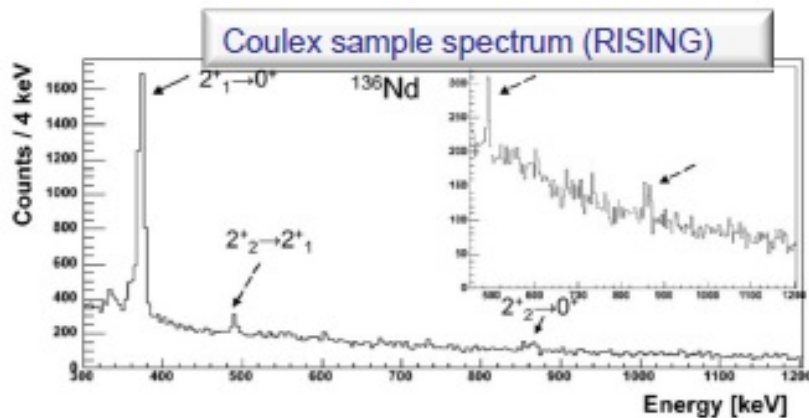


NuSTAR experiments

HISPEC: high resolution in-flight gamma-spectroscopy

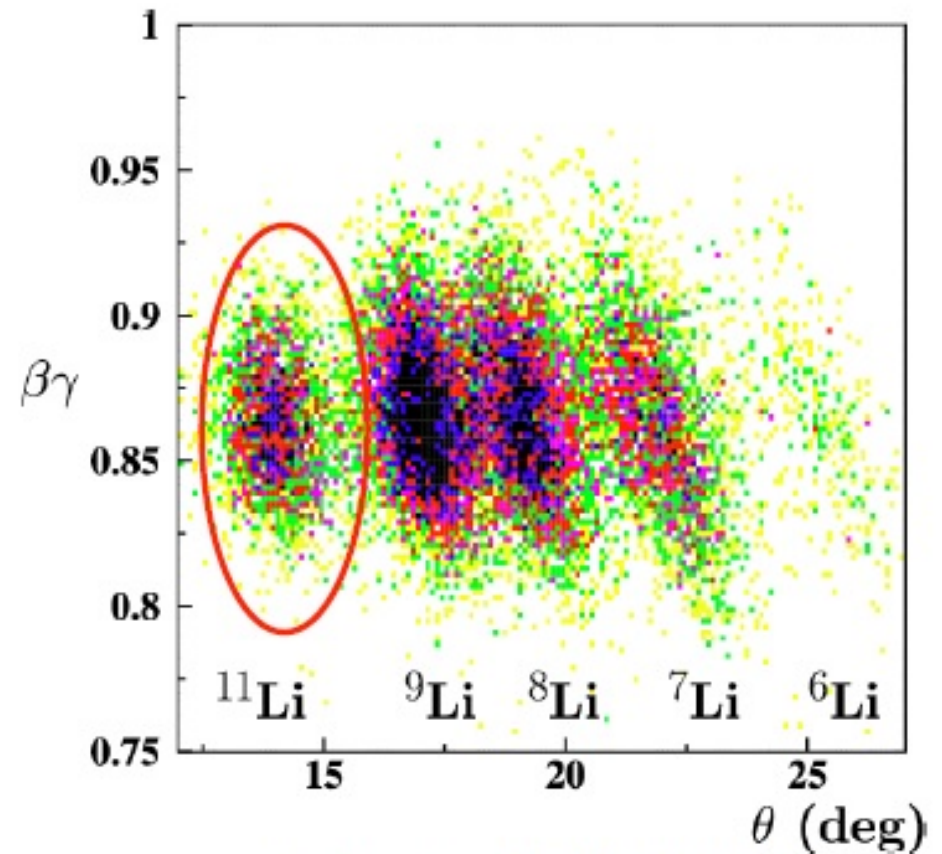
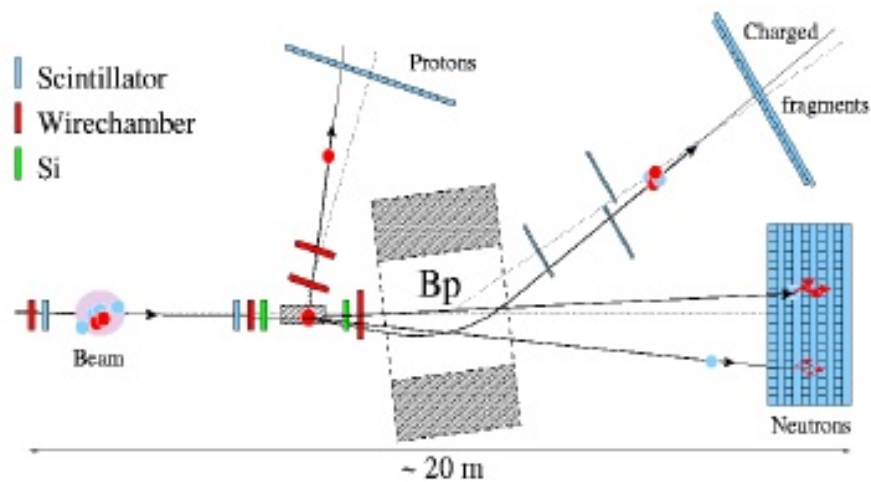
- Excited states in exotic nuclei
 - Inelastic excitation (EM, hadronic), secondary fragmentation, knock-out
- Lifetimes and g-factors of short-lived excited states
 - (evolution of shells and collectivity)
 - (phase-/shape-transitions & shape coexistence)
 - (isospin symmetry, pn-pairing)

- **Experiments possible with medium-energy beams (50-100 A MeV)**
- **Broad program for nuclei with A up to ~100**



Proton knockout from 304 MeV/u ^{14}Be

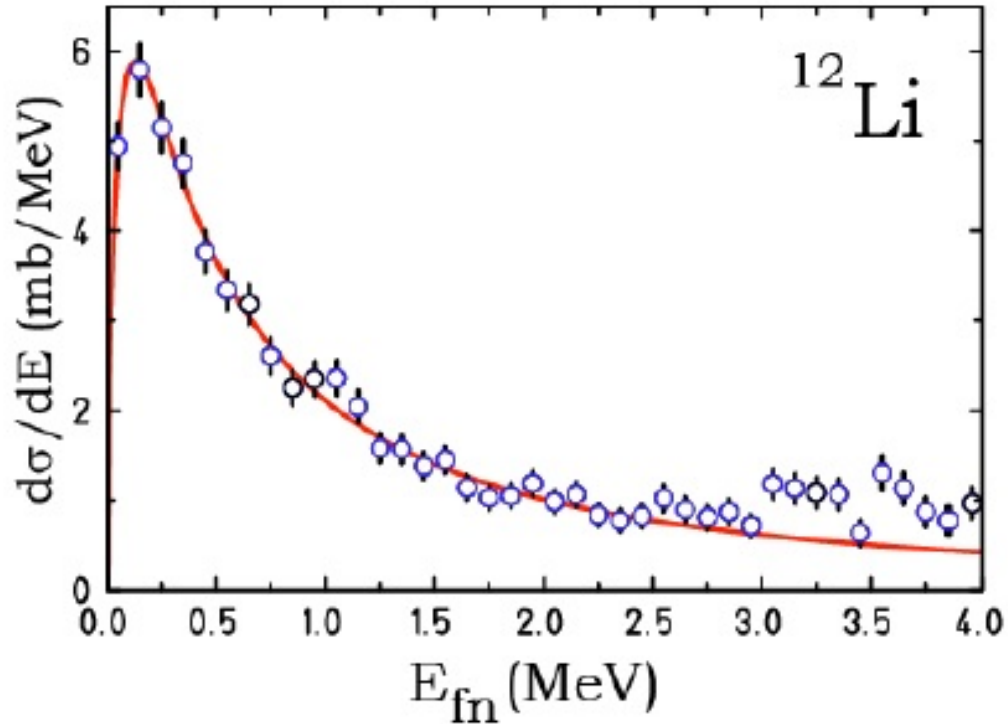
S245@GSI



Yu. Aksyutina *et al.* PLB 666(2008)430

Going beyond the dripline...

Breit Wigner Resonance



$$\varepsilon = 1.47(19)\text{MeV}$$

$$S_{2n} = 1.26(13)\text{ MeV}$$

$$E_{fn} = \left| \bar{P}_f + \bar{P}_n \right| - M_f - m_n$$

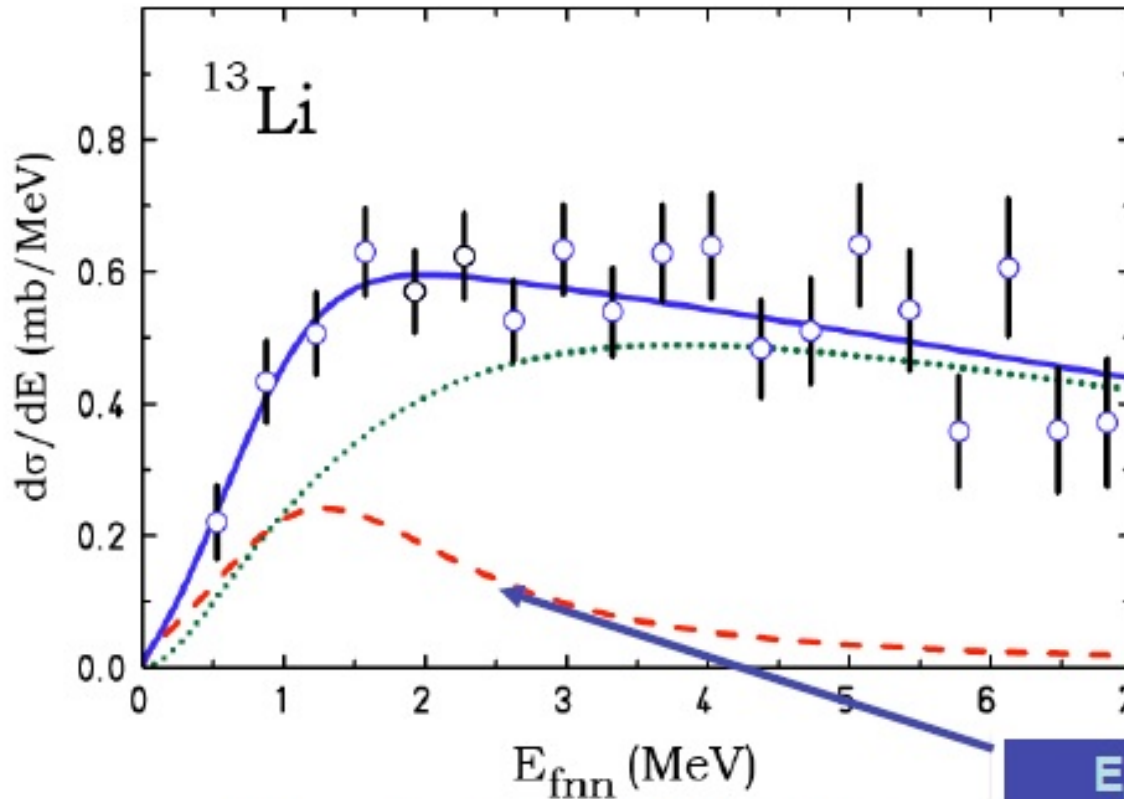
$$a_s = -13.7(1.6)\text{ fm}$$

scattering length (virtual s-state)

Yu. Aksyutina *et al.* PLB 666(2008)430

... and even further

Breit Wigner Resonance



$$\frac{d\sigma}{dE_{f2n}} = \frac{E_{f2n}^2}{(2.21S_{2n} + E_{f2n})^{7/2}}$$

C. Forssén *et al.* NPA 673 (2008) 143

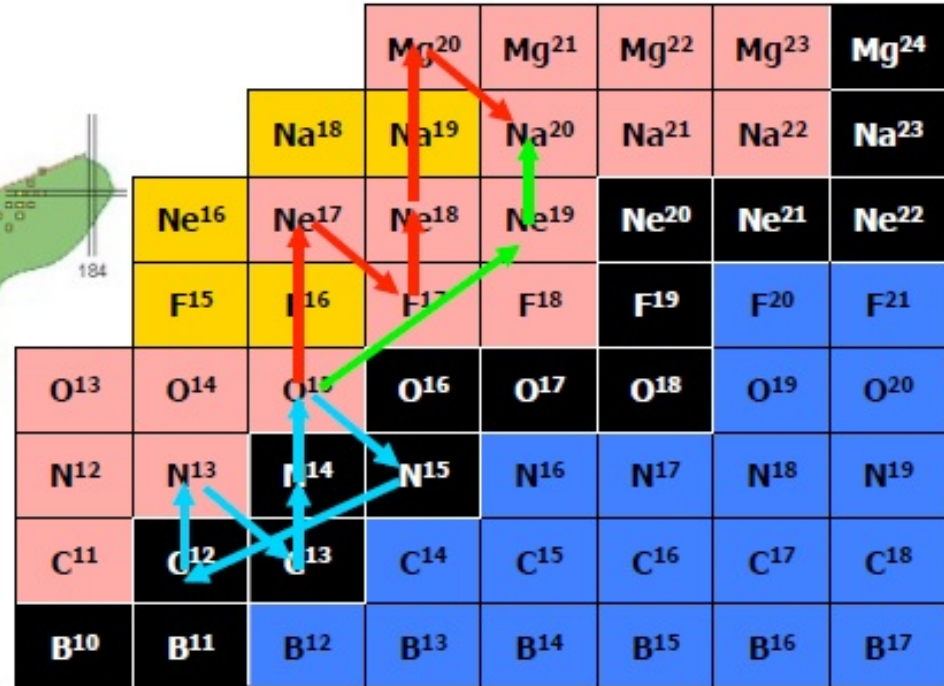
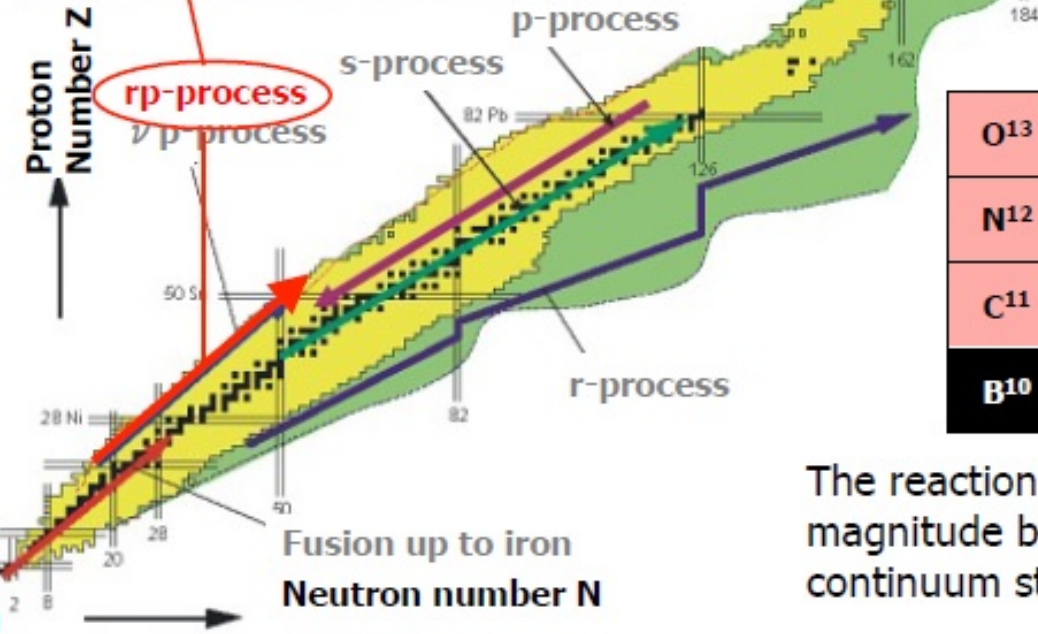
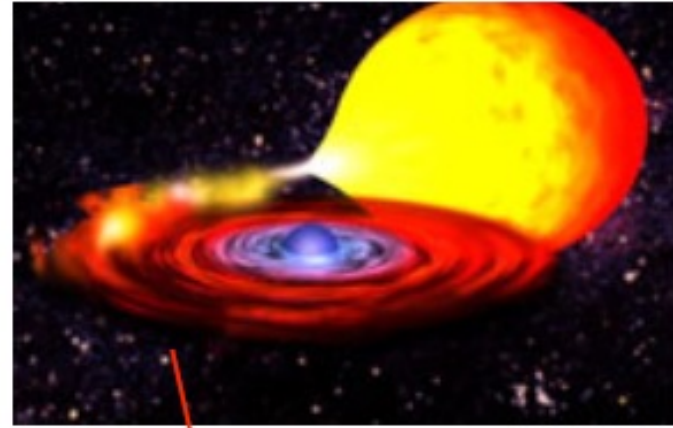
Yu. Aksytina *et al.* PLB 666(2008)430

- 3-body resonance state ---
- correlated background ---
(from nn-correlations in initial bound-state wave function)

$^{15}\text{O}(2p,\gamma)^{17}\text{Ne}$ in Nuclear Astrophysics: X-ray bursts, rp-process, neutron stars

Cataclysmic binary systems (X-ray bursts):
 rp-process Görres et al., PRC 51 (1995) 392

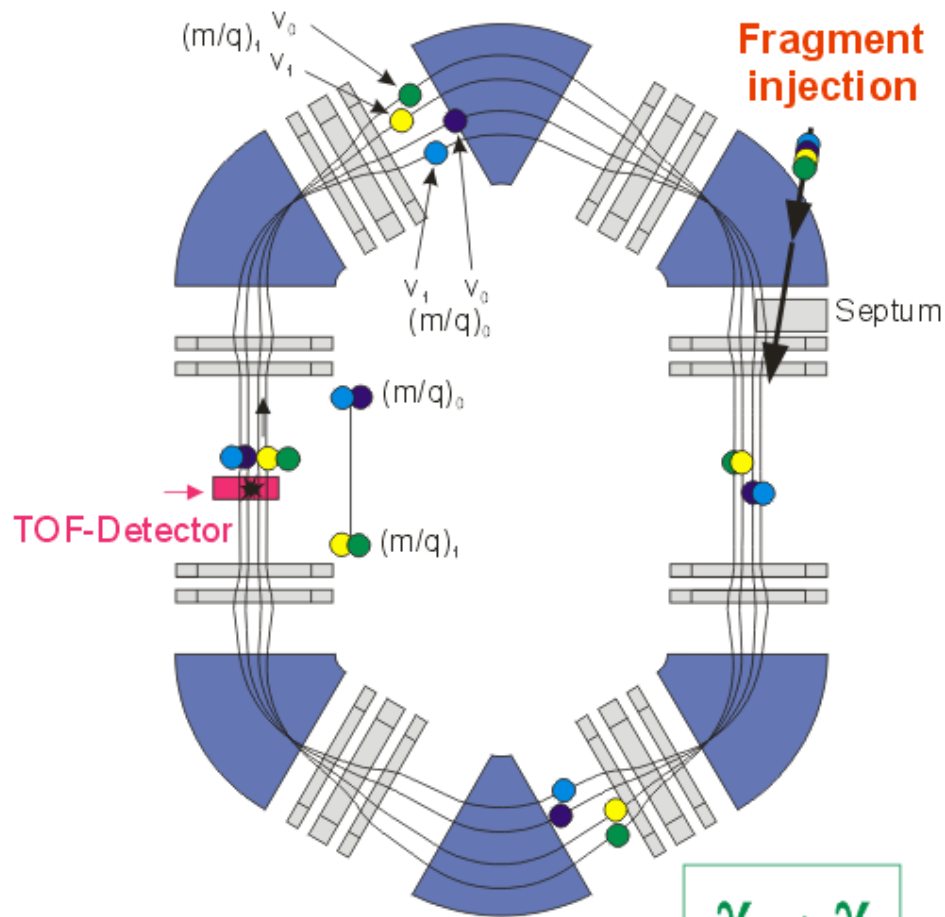
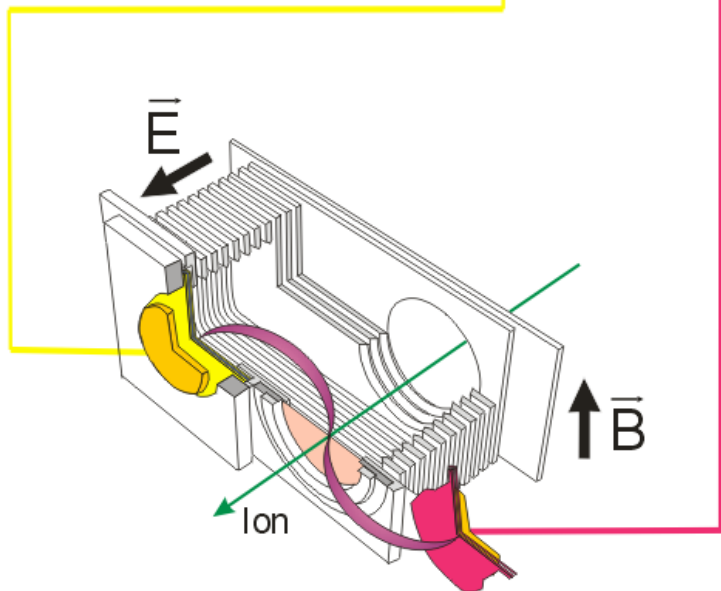
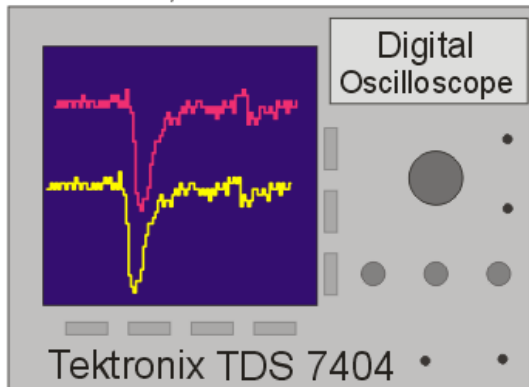
CNO cycle: ... $^{14}\text{N}(p,\gamma)^{15}\text{O}(\beta^-)^{15}\text{N}(p,\alpha)$
 ^{15}O is a waiting point for CNO-cycle breakup:
 Heavier elements: ... $^{15}\text{O}(\alpha,\gamma)^{19}\text{Ne}(p,\gamma)$...
 Alternative (rp): ... $^{15}\text{O}(2p,\gamma)^{17}\text{Ne}(\beta^-)^{17}\text{F}(p,\gamma)$



The reaction rate can be enhanced by a few orders of magnitude by taking into account the three-body continuum states. Grigorenko et al., PLB 641 (2006) 254

Isochronous mass spectroscopy in the ESR

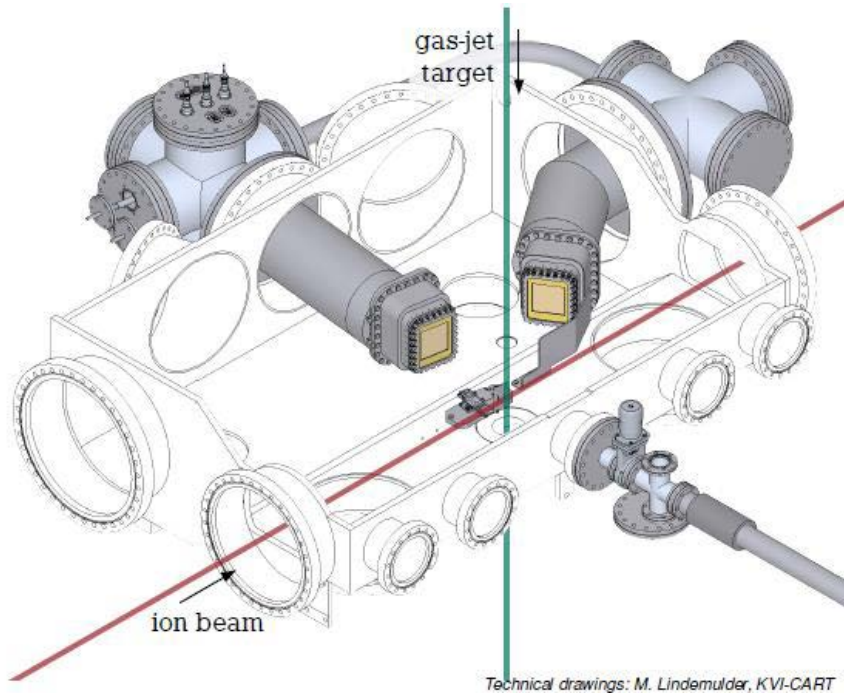
20GS/s , 10GS/s 2 channels



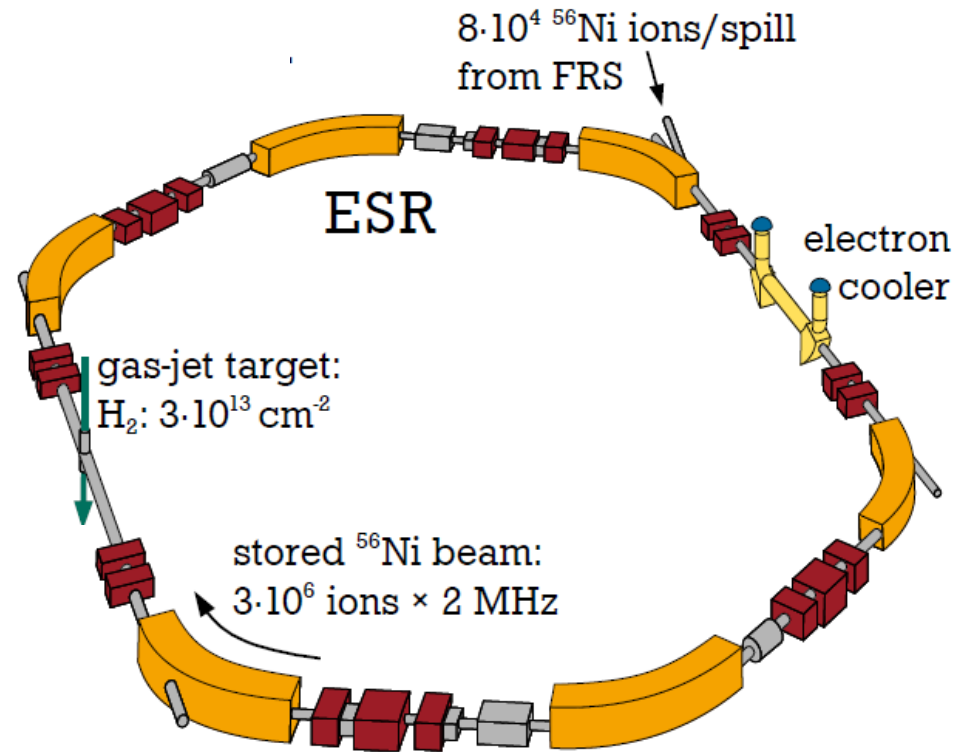
$$\gamma_t \rightarrow \gamma$$

$$\frac{\Delta f}{f} = -\frac{1}{\gamma_t^2} \frac{\Delta(m/q)}{m/q} + \frac{\Delta V}{V} \left(1 - \frac{\gamma^2}{\gamma_t^2}\right)$$

Scattering of RIB on hydrogen target



measure position and energy with Si detectors



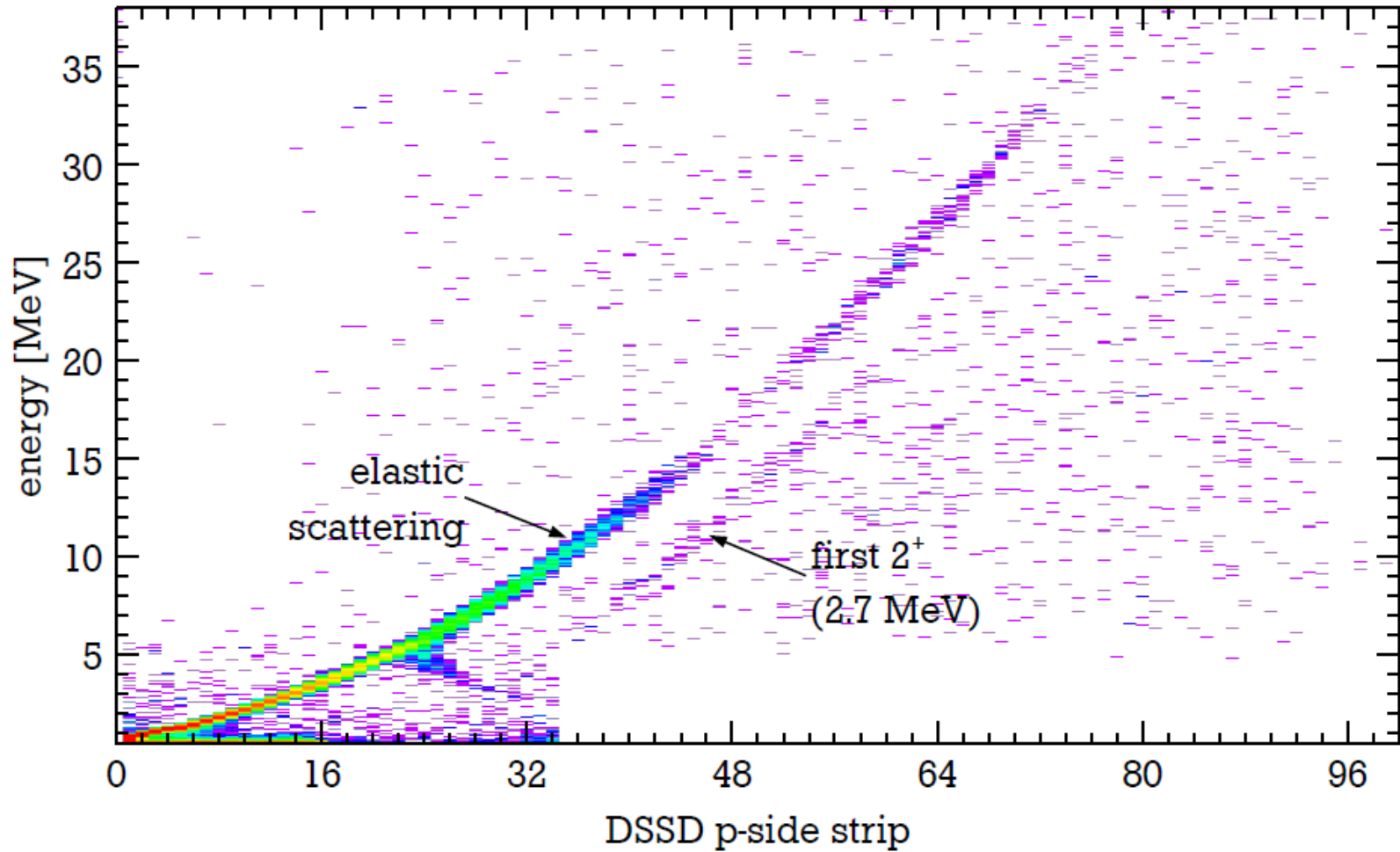
Picture: Phys. Scr. T156 (2013) 014016

➔ RMS radius of ^{56}Ni

luminosity: $2 \cdot 10^{26} \frac{\text{particles}}{\text{s cm}^2}$

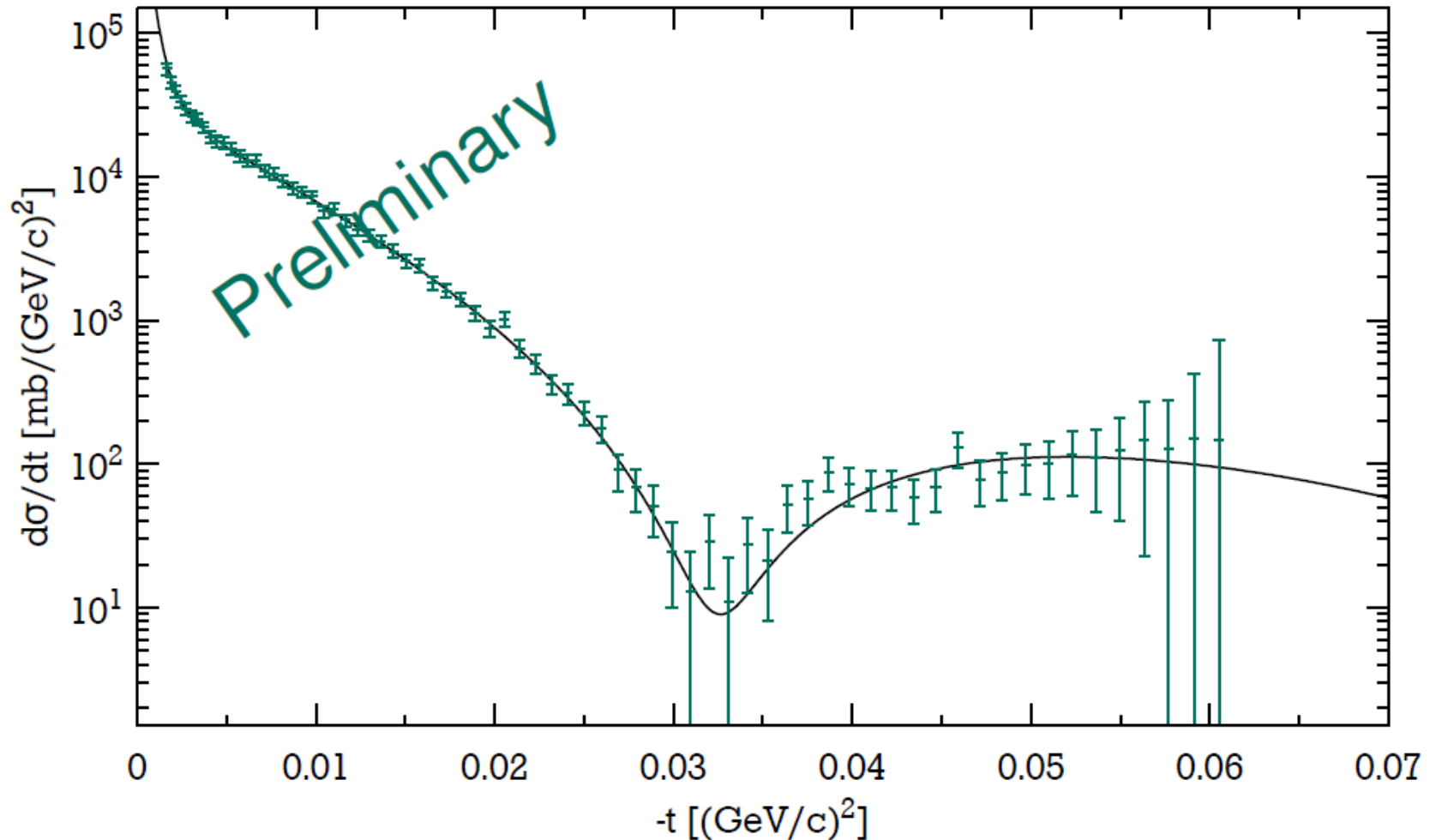
M. von Schmid, EXL, Phys. Scr. T116 (2015) 014005

$^{56}\text{Ni}(p,p)$ scattering distribution



$^{56}\text{Ni}(p,p)$ scattering distribution

Diffraction pattern (like for a wave after a single slit),
Extract radius of nucleus by fitting theory with parameters.



M. von Schmid, EXL, Phys. Scr. T116 (2015) 014005

Electron capture in hydrogen-like ions

Simple theoretical estimate:

R.B. Firestone, Table of Isotopes, 1996

Gamow-Teller transition $1^+ \rightarrow 0^+$

β^+ to EC branching ratio:

$$\lambda_{\beta^+}/\lambda_{EC} \text{ (neutral atom)} \approx 1$$

W. Bambynek et al., Rev. Mod. Phys. 49, 1977

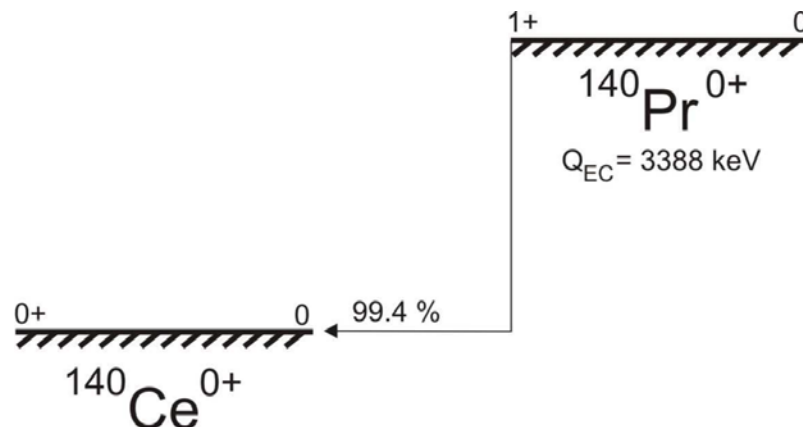
S-electron density at the nucleus:

$$|f_s(0)|^2 \propto 1/n^3$$

$$P_{EC}(\text{neutral atom}) \propto 2 \sum 1/n^3 = 2.4$$

$$P_K(H\text{-like}) \propto 1 * 1/1^3 = 1$$

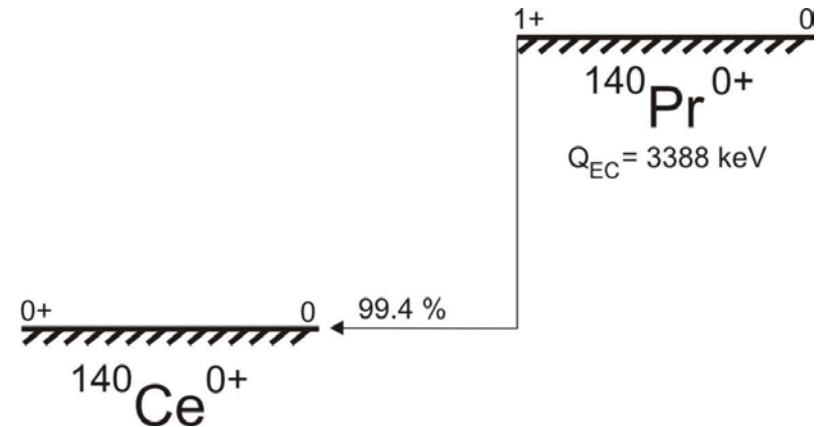
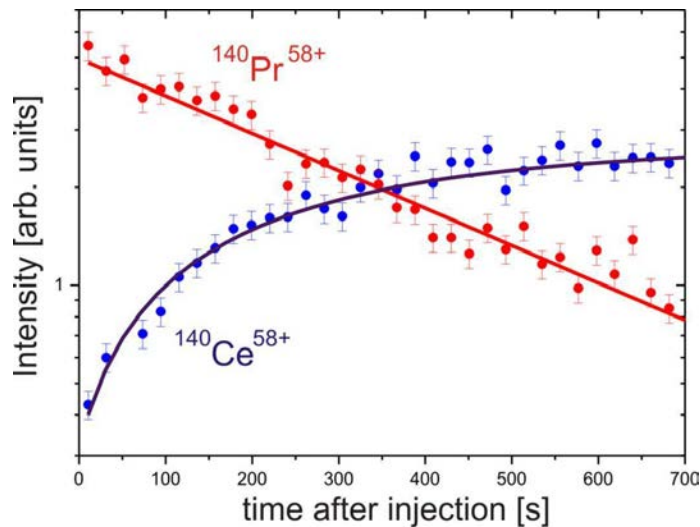
$$\lambda_{\beta^+}/\lambda_K \text{ (H-like)} \approx 2.4$$



conclusion:

H-like ions should have
41% longer half-life

Electron capture in hydrogen-like ions



G. Audi et al. NPA 729 (2003) 3

$$\lambda(\text{neutral}) = 0.0034(1) \text{ s}^{-1}$$

Decay of fully-ionized ^{140}Pr :

$$\lambda_{\beta^+} = 0.00172(7) \text{ s}^{-1}$$

Decay of H-like ^{140}Pr :

$$\lambda_K = 0.00213(19) \text{ s}^{-1}$$

J. Kurcewicz, N. Winckler et al.

Expectation:

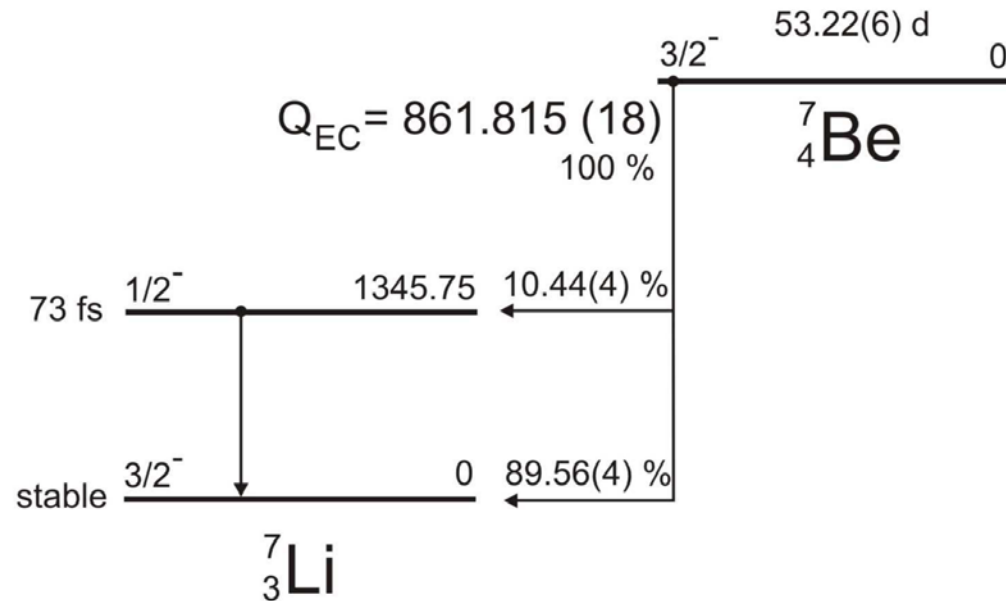
$$\lambda_{\beta^+}/\lambda_K \text{ (H-like)} \approx 2.4$$

Experiment:

$$\lambda_{\beta^+}/\lambda_K \text{ (H-like)} = 0.81 (8)$$

H-like ions decay ~20%
faster than neutral atom!!!

Electron capture in hydrogen-like ions



A.V. Gruzinov, J.N. Bahcall, Astroph. J. 490 (1997) 437

Ionization of ${}^7\text{Be}$ in the Sun can be $\sim 20\text{-}30\%$

Nucleosynthesis of heavier elements

