

*Zero-degree Auger Projectile Spectrometry  
in the New Experimental Storage Ring:  
Challenges and Prospects*



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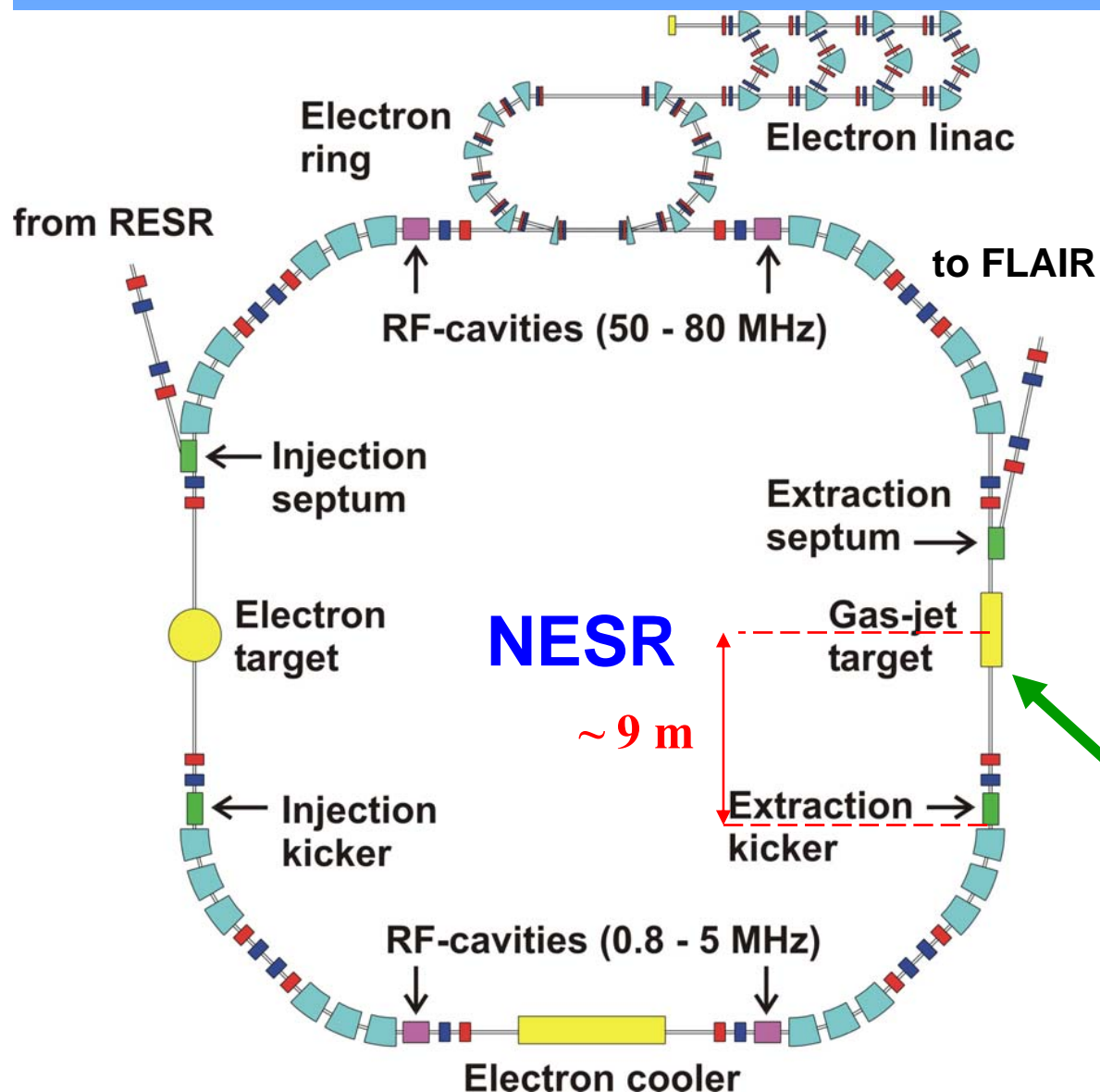
SPARC-Paris, February 12-15, 2007

*Motivation: High resolution electron spectrometry in the NESR*

- **Auger and conversion lines of high-Z few-electron ions**  
produced in collisions with atoms
  - **No measurements for ions!!!** (only neutral targets to date)
  - Charge state q-dependence
  - high precision Binding Energy determination of atomic levels
- **Physics of strong fields:**
  - Collision dynamics, state-selective cross section determination
  - Atomic Structure of high-Z few-electron ions (relativistic effects)

This will require **electron observation in the beam direction**  
known as **Zero-degree Auger Projectile Spectroscopy (ZAPS)**  
**Of relativistic electrons with lab energies up to about 0.5 MeV**  
**Combined with a spectrometer resolution of  $\Delta p/p \leq 10^{-4}$**   
**access to the natural line widths  $\Gamma$  should also be possible**

# The NESR: the ultimate fantasy for atomic collisions physics?



## NESR

- Circumference: 221.11 m
- Vacuum:  $\leq 10^{-11}$  mbar
- Ion energies: 4 - 740 MeV/u
- Ion beam species: H - U
- Radioactive beams: yes
- Ion charge states:  $A/q \leq 2.7$
- Number of ions:  $\sim 10^8$
- Beam particle current:  $\sim 10^{13}$  #/s
- Emittance: 0.1 - 1 mm-mrad
- Momentum  $\Delta p/p$  (cooled):  $\leq 10^{-4}$

### Gas-jet target (extrapolation from ESR)

- Areal Density:  $10^{12}$ - $10^{13}$  #/cm<sup>2</sup>
- Length: 1.4 - 4 mm
- Background pressure:  $10^{-9}$  mbar

*Two different electron spectroscopies*

Zero-degree Auger Projectile e- Spectroscopy (ZAPS)

- used successfully since the 1980's primarily at Tandems to measure Auger electrons from projectiles excited via capture, excitation or ionization and combinations
- emitter: low-Z HCI ions in the 0.1- 3 MeV/u collision regime
- electrostatic spectrometers:

lab electron energies  $\varepsilon = 0-6$  keV and  $R_\varepsilon = \Delta\varepsilon/\varepsilon \sim 0.1\%$

$$\frac{\Delta p}{p} = \frac{\gamma}{\gamma + 1} \frac{\Delta \varepsilon}{\varepsilon}$$

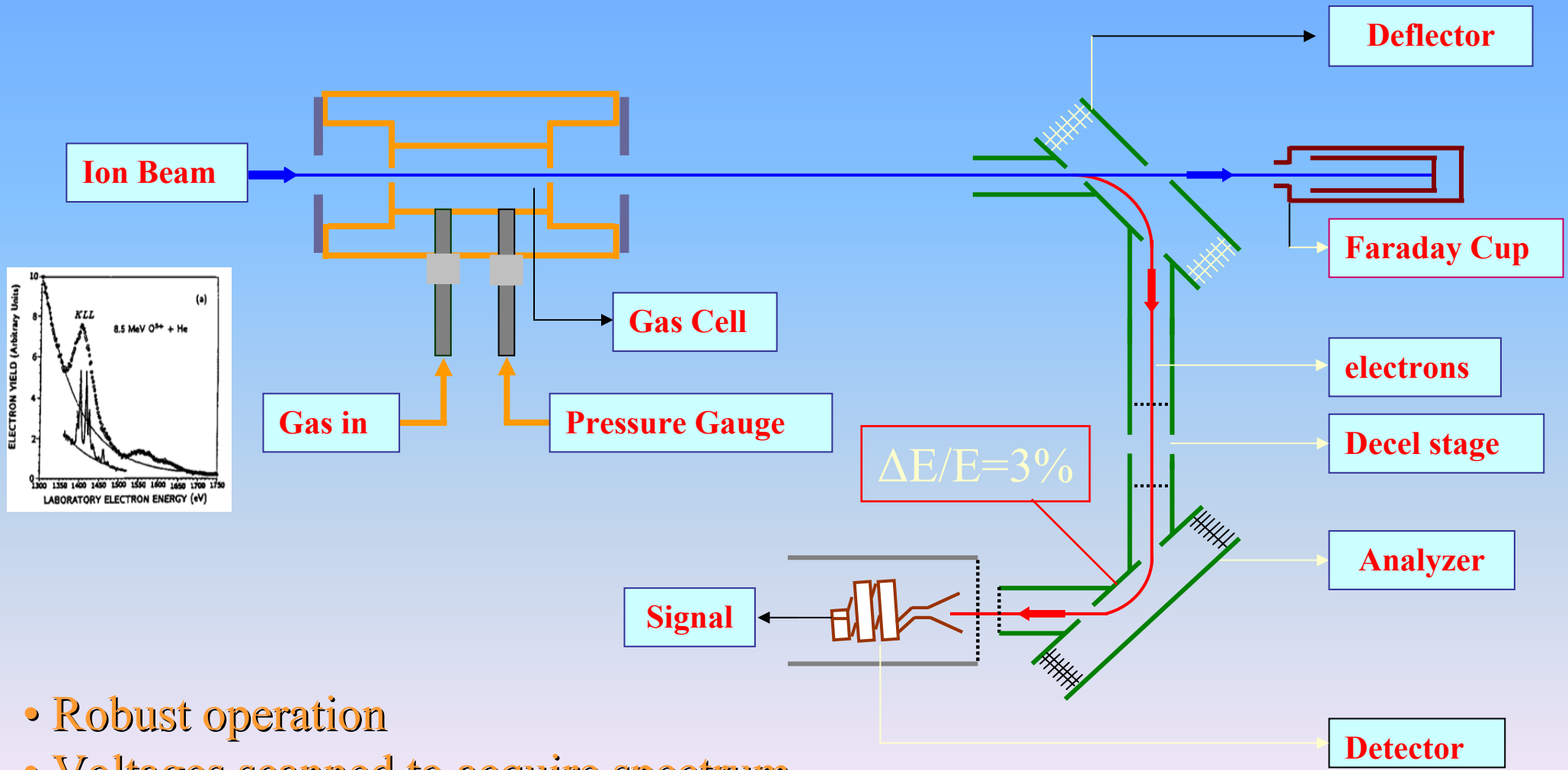
$\beta$ -ray spectroscopy

- used successfully in the 1950-1970's at high flux reactors or with radioactive sources to measure conversion and Auger electrons
- emitter: stationary high-Z neutral **activated** target atoms
- large radius (e.g. 50cm Uppsala, 100cm Chalk River, 50cm BILL) double focusing magnetic spectrometers:

lab electron energies  $\varepsilon \leq 3$  MeV and  $R_p = \Delta p/p \sim 0.01- 0.05\%$

The SPARC electron spectroscopy initiative  
is expected to combine both expertise

*Traditional  $0^0$   $e^-$  spectrometer  
two-stage  $45^0$  parallel plate with intermediate deceleration stage*

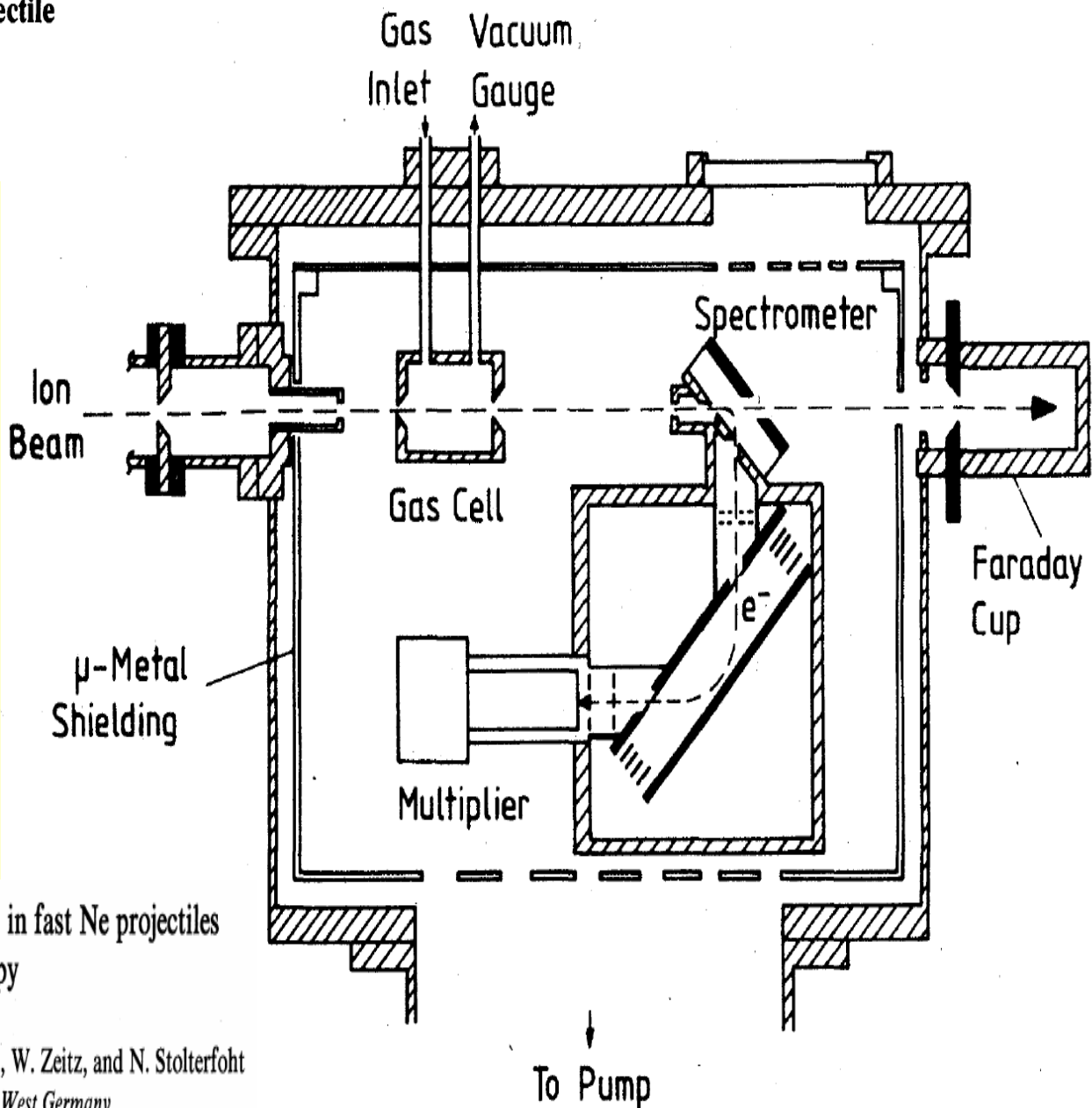
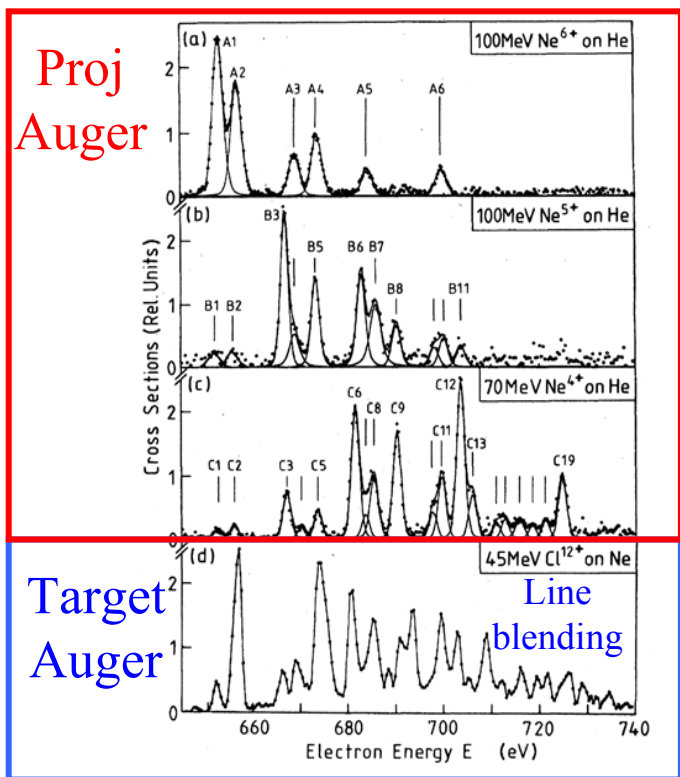


- Robust operation
- Voltages scanned to acquire spectrum
- High resolution  $\sim 0.1\%$  uses deceleration stage with fixed pass energy

**Selective production of Auger electrons from fast projectile ions studied by zero-degree Auger spectroscopy**

JPB16 (1983) 3965

A Itoh, T Schneider, G Schiwietz, Z Roller, H Platten, G Nolte, D Schneider and N Stolterfoht



**Selective production of Li-, Be-, and B-like K vacancy states in fast Ne projectiles studied by zero-degree Auger spectroscopy**

PRA31 (1985) 684

A. Itoh,\* D. Schneider, T. Schneider, T. J. M. Zouros, G. Nolte, G. Schiwietz, W. Zeitz, and N. Stolterfoht

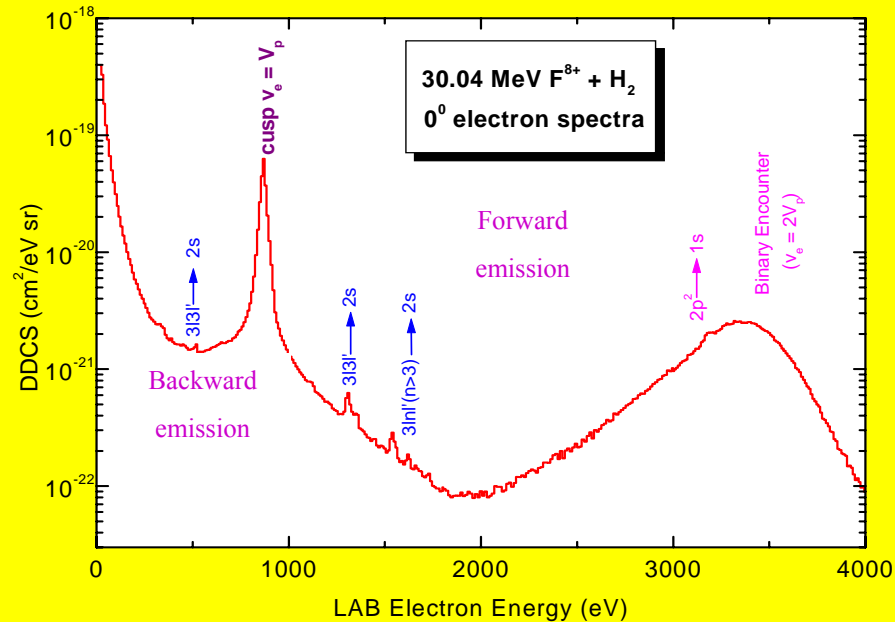
Hahn-Meitner-Institut für Kernforschung Berlin, D-1000 Berlin 39, West Germany

## *Advantages of $0^0$ Auger projectile spectrometry (ZAPS)*

- Only a **single pre-selected Projectile charge state** involved
  - considerable simplification of lines in spectrum
  - no line blending (mixture of different charge states)
  - “ion surgery” collisions with low-Z few-electron targets (He, H<sub>2</sub>)
  - very successful in isoelectronic studies
- **High resolution** technique with relatively **high overall efficiency**
  - $\Delta E/E \sim 0.1\%$ ,  $\Delta\theta \sim 1^\circ$ ,  $\Delta\Omega \sim 10^{-4}$  sr
  - Resolution good enough to resolve most **K-Auger** lines
  - Much **more efficient than comparable resolution crystal X-ray spectrometers** (for **low-Z ions** high Auger yield, no window absorption, large  $\Delta\Omega$ )
- Determination of **absolute** double differential **cross sections**
  - collisional energy dependence of well defined transition
- Deceleration stage provides useful variable resolution
  - low resolution or high resolution can be used as needed

Question: How can ZAPS be done effectively in the NESR?

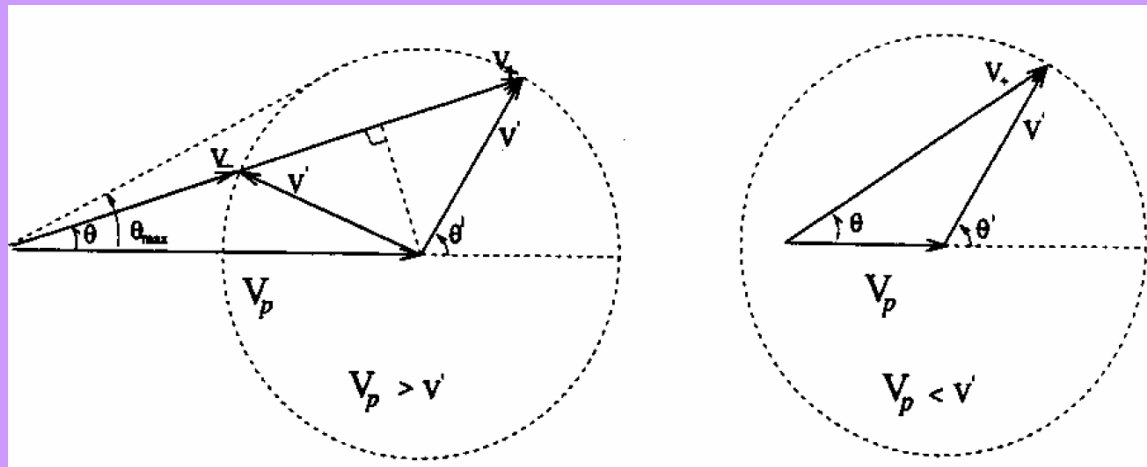
## Important features of $0^\circ$ Electron emission spectrum



1. Low energy Target e- continuum
2. Cusp e- at  $v = V_p$
3. High energy Projectile e- continuum
4. Broad Binary Encounter e- Peak
5. Kinematically shifted Auger Projectile e- lines

## Emission from moving source

### Kinematics: Laboratory energy Shifting and Doubling



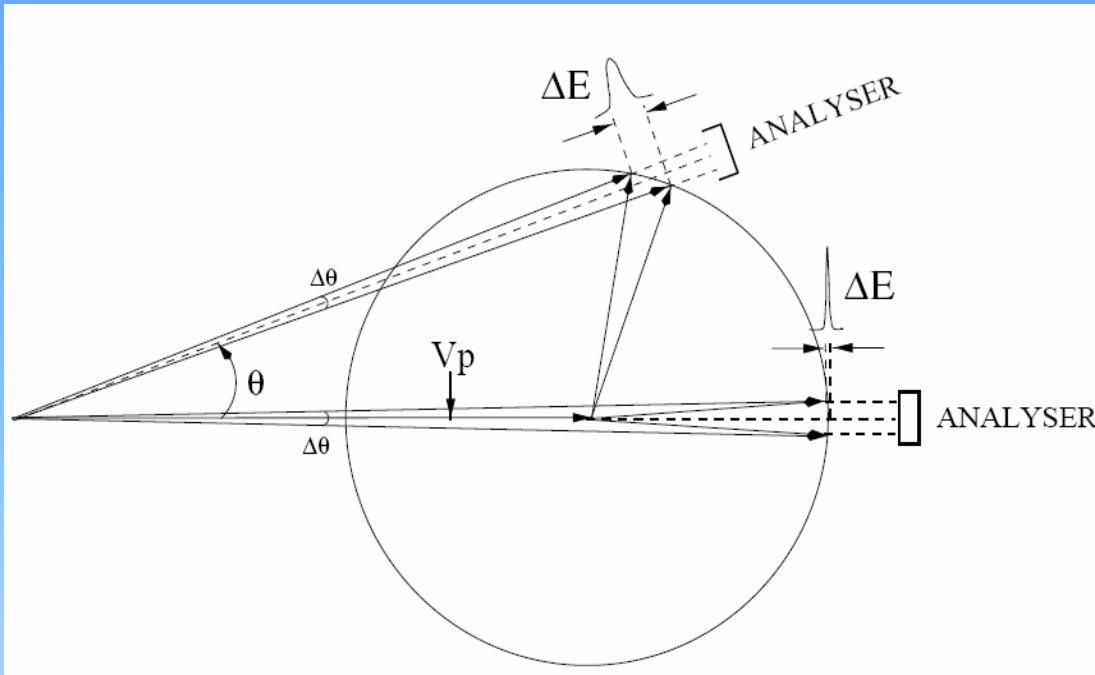
Projectile electrons are shifted energetically to higher and lower laboratory energies depending on whether they are emitted in the forward (+) or backward (-) direction.

For  $\theta = 0^\circ$ :

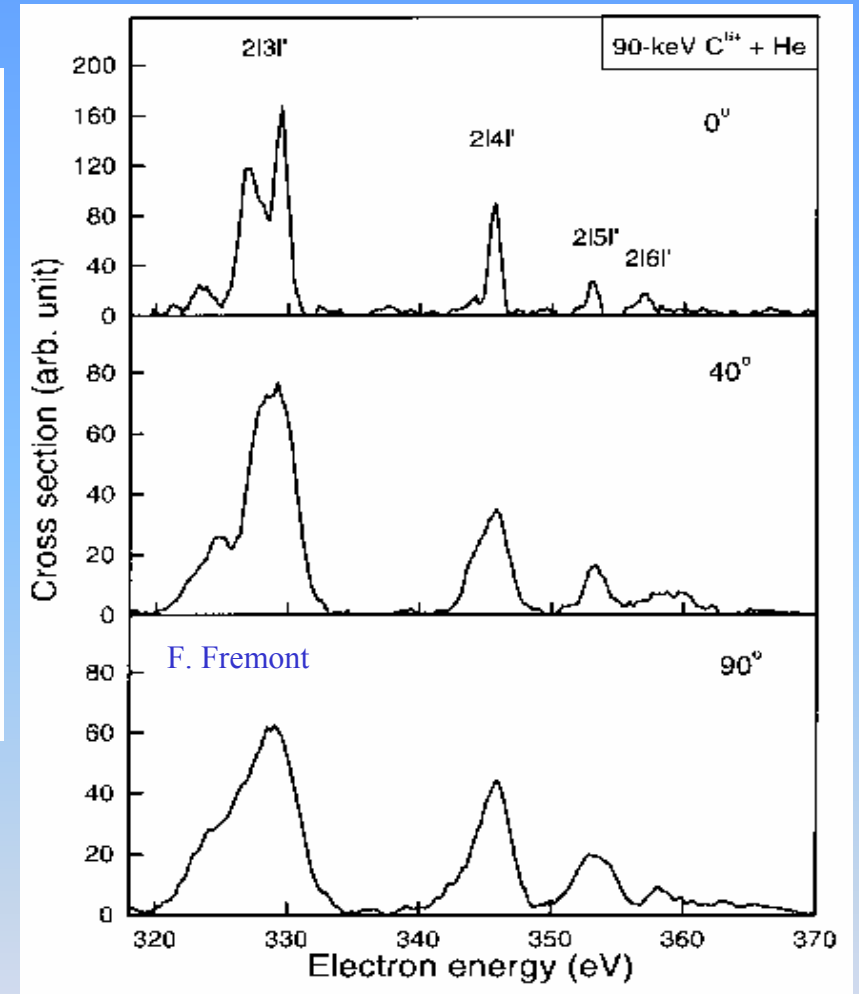
$$v_{\pm} = V_p \pm v' \quad (\text{classical})$$

$$\beta_{\pm} = \frac{\beta_p \pm \beta'}{1 + \beta_p \beta'} \quad (\text{relativistic})$$

# Kinematics: Instrumental Line Broadening I



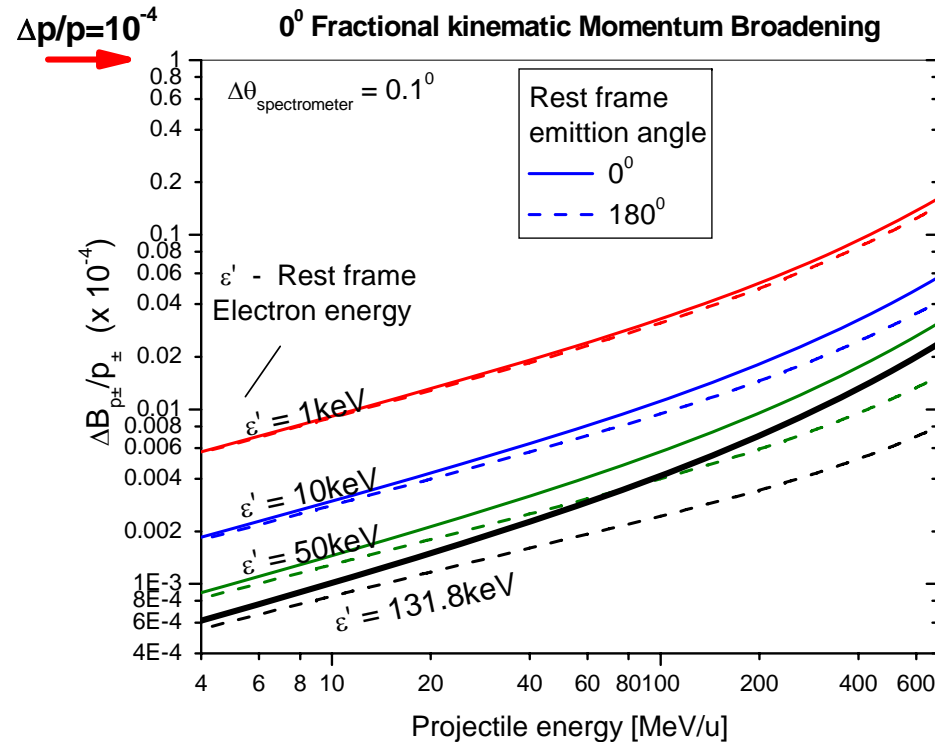
$$\Delta E_{\theta} = \begin{cases} |\varepsilon(\theta + \Delta\theta/2) - \varepsilon(\theta - \Delta\theta/2)| & \text{for } \theta > 0^{\circ} \\ |\varepsilon(\Delta\theta/2) - \varepsilon(0^{\circ})| & \text{for } \theta = 0^{\circ} \end{cases}$$



For  $\theta > 0^{\circ}$   $\Delta E_{\theta} \sim \Delta\theta$ , while for  $\theta = 0^{\circ}$   $\Delta E_{\theta} \sim \Delta\theta^2$   
 thus substantial gains in resolution can be attained  
 by going to  $\theta = 0^{\circ}$  observation angle

## Kinematics: Instrumental Line Broadening II

- At 0° the kinematic broadening
- grows with projectile velocity  $V_p$
  - grows approximately as  $\Delta\theta^2$
  - diminishes with Auger energy  $\epsilon'$



2nd order fractional momentum broadening:

$$\frac{\Delta B_{p_{\pm}}}{p_{\pm}} = \frac{\Delta\theta^2}{8} \frac{V_p}{v'} \quad (\text{classical})$$

$$\frac{\Delta B_{p_{\pm}}}{p_{\pm}} = \frac{\Delta\theta^2}{8} \frac{\beta_p}{\beta'} \frac{|1 \pm \beta_p \beta'|}{(1 - \beta_p^2)} \quad (\text{relativistic})$$

Range in spectrometer acceptance angle  $\Delta\theta$

$\Delta B_{p_{\pm}}/p_{\pm} = 10^{-4}$	$\epsilon' = 1 \text{ keV}$		$\epsilon' = 131 \text{ keV}$	
	4 MeV/u	740 MeV/u	4 MeV/u	740 MeV/u
$\Delta\theta$ (dgrs)	1.35	0.25	4.15	0.65
$\Delta\Omega$ ( $10^{-4}$ sr)	5.55	0.19	52.5	1.29

## *Comparison of Tandem - NESR Typical Beam and Target operational parameters*

	Target density n ( $10^{13}$ #/cm <sup>2</sup> )	Target Length L (cm)	n L ( $10^{13}$ #/cm <sup>2</sup> )	Charge State q	Beam Current $I_p = I_q / qe$ ( $10^{10}$ #/s)	$I_p n L$ ( $10^{25}$ #/cm <sup>2</sup> s)
<b>Tandem</b>	160 (50mTorr)	5	800	1-10	0.6-60	4.8-480
<b>NESR</b>	16* (5mTorr)	0.1-0.4	1.6 - 6.4	$A/q \leq 2.7$ <small>238U<sup>89-92+</sup></small>	1250 - 11000**	20 -700

\*projected from experience with ESR jet target

\*\* Assuming a constant  $10^8$  particles in the NESR over the 4-740 MeV/u energy range and for  $q=92$

### NESR advantages (+) vs disadvantages (-)

(+) High particle current, increases with collision energy

(-) jet target: smaller density and effective length (but note Grisenti talk on liquid targets)

Question: Possibility of cell target? Would increase rate by at least x500!



A ONE-METER-RADIUS IRON-FREE DOUBLE-FOCUSING  $\pi\sqrt{2}$  SPECTROMETER

FOR  $\beta$ -RAY SPECTROSCOPY WITH A PRECISION OF 1:10<sup>5</sup>

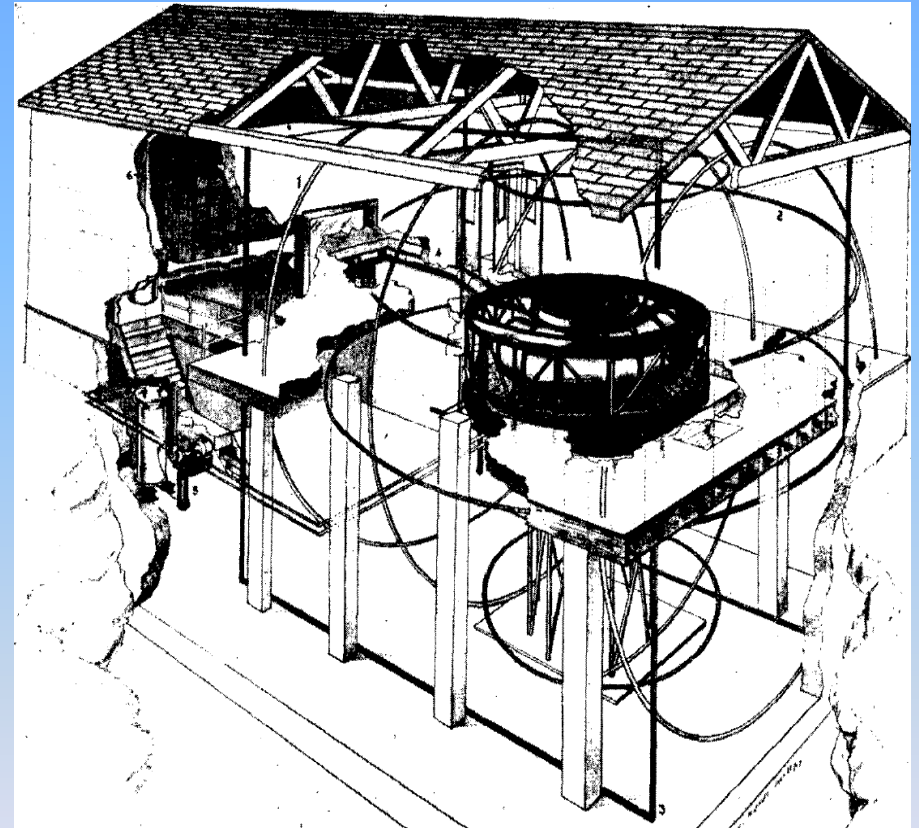
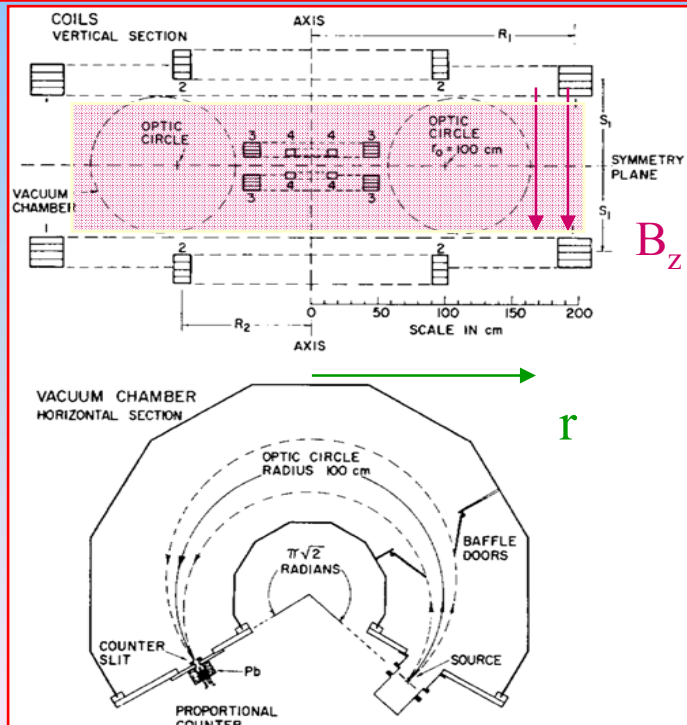
R. L. GRAHAM, G. T. EWAN and J. S. GEIGER

Physics Division, Atomic Energy of Canada Limited, Chalk River, Ontario

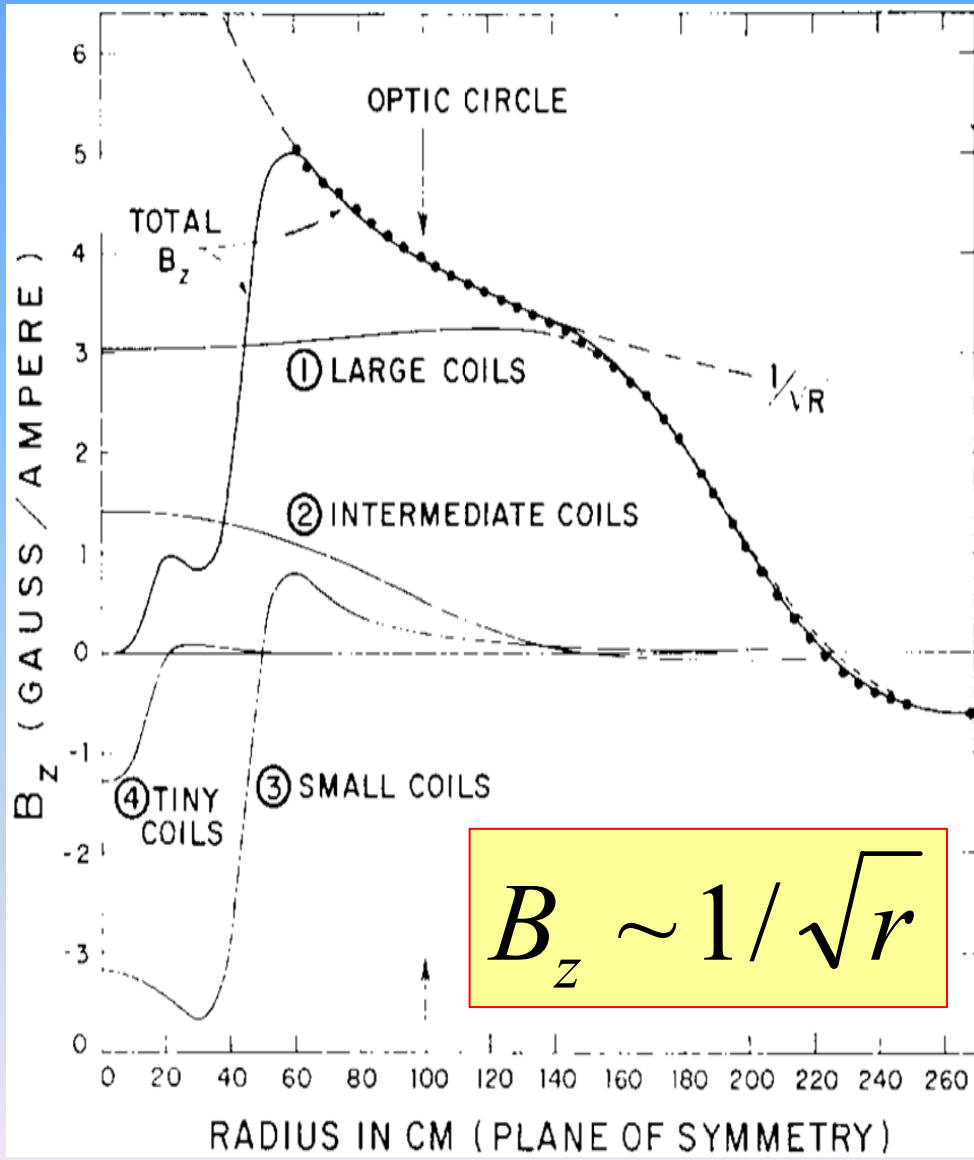
*$\beta$  – ray spectroscopy*

Parameters for optimum luminosity when using a counter<sup>a)</sup>

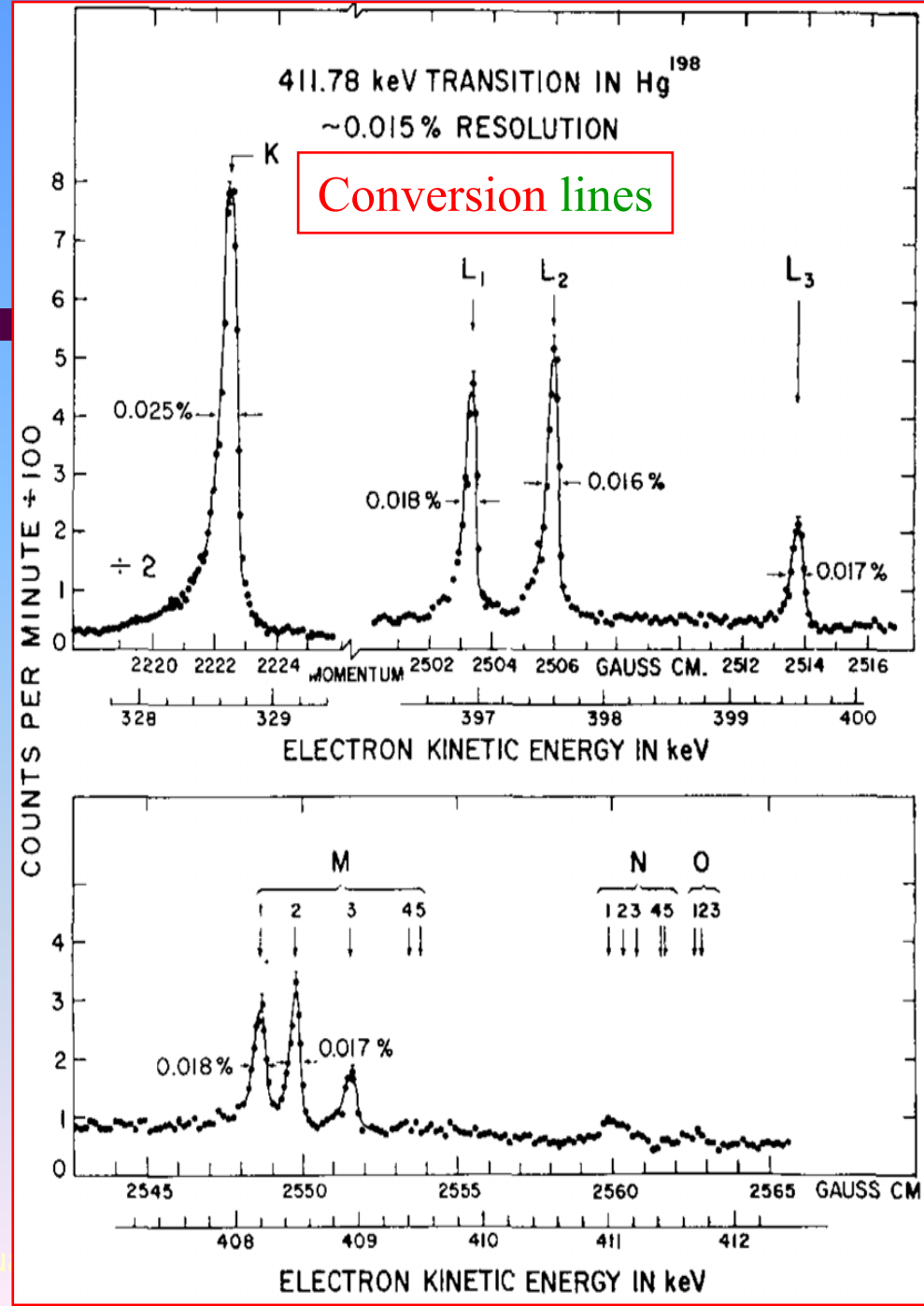
Resolution $R$ (%)	0.01	0.1	1.0
Source width $s = 2h_0$ (cm)	0.034	0.34	3.4
Source height $2t_0$ (cm)	3.7	11.6	37
Source area $A$ (cm <sup>2</sup> )	0.126	3.94	126
Transmission <sup>b)</sup> $\Omega/4\pi$	0.06	0.20	0.80
Luminosity $L = A \Omega/4\pi$ (cm <sup>2</sup> )	$7.6 \times 10^{-5}$	$7.9 \times 10^{-3}$	$10.1 \times 10^{-1}$
Counter slit width (cm)	0.034	0.34	3.4
Counter acceptance factor $F$	0.83	0.83	0.83
Effective luminosity $LF$ (cm <sup>2</sup> )	$6.3 \times 10^{-5}$	$6.6 \times 10^{-3}$	$8.4 \times 10^{-1}$



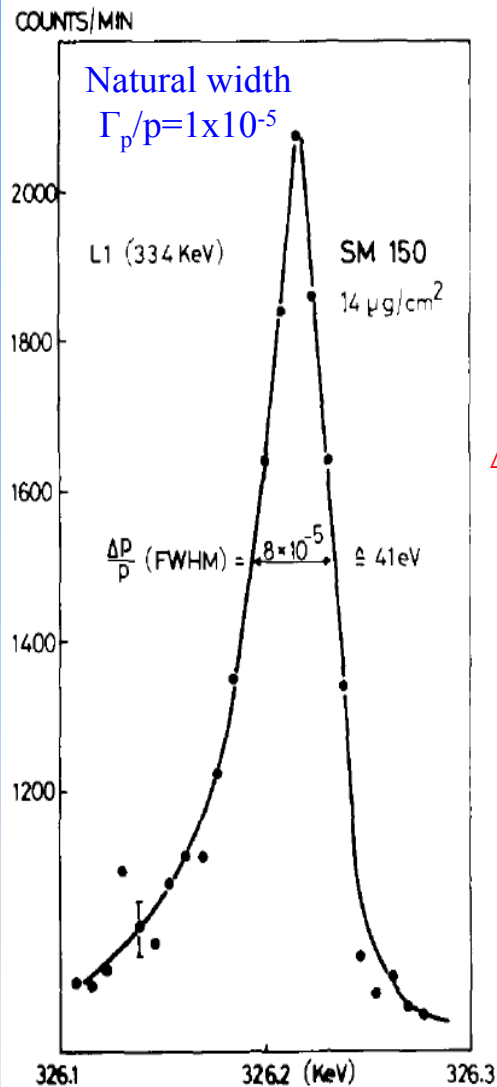
Dedicated machines requiring huge housing  
To ensure field uniformity and easy access



$$B_z \sim 1/\sqrt{r}$$



# 50cm 1/r BILL electron iron-core spectrometer



Slit widths:  
Entry =  $s_1$   
Exit =  $s_2$   
 $s_1 = s_2 = 0.2 \text{ mm}$   
  
 $\rho = 50 \text{ cm}$

Best resolution  
 $\Delta p/p = 7.6 \times 10^{-5}$

Fig. 18. Best resolution of BILL measured with the  $L_1$  line of the 334 keV transition in  $^{150}\text{Sm}$ .

NUCLEAR INSTRUMENTS AND METHODS 154 (1978) 127-149

## THE DOUBLE FOCUSING IRON-CORE ELECTRON-SPECTROMETER "BILL" FOR HIGH RESOLUTION ( $n, e^-$ ) MEASUREMENTS AT THE HIGH FLUX REACTOR IN GRENOBLE

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J. LARYSZ and T. von EGIDY\*

*Institut Laue-Langevin, 38042 Grenoble Cédex, France*

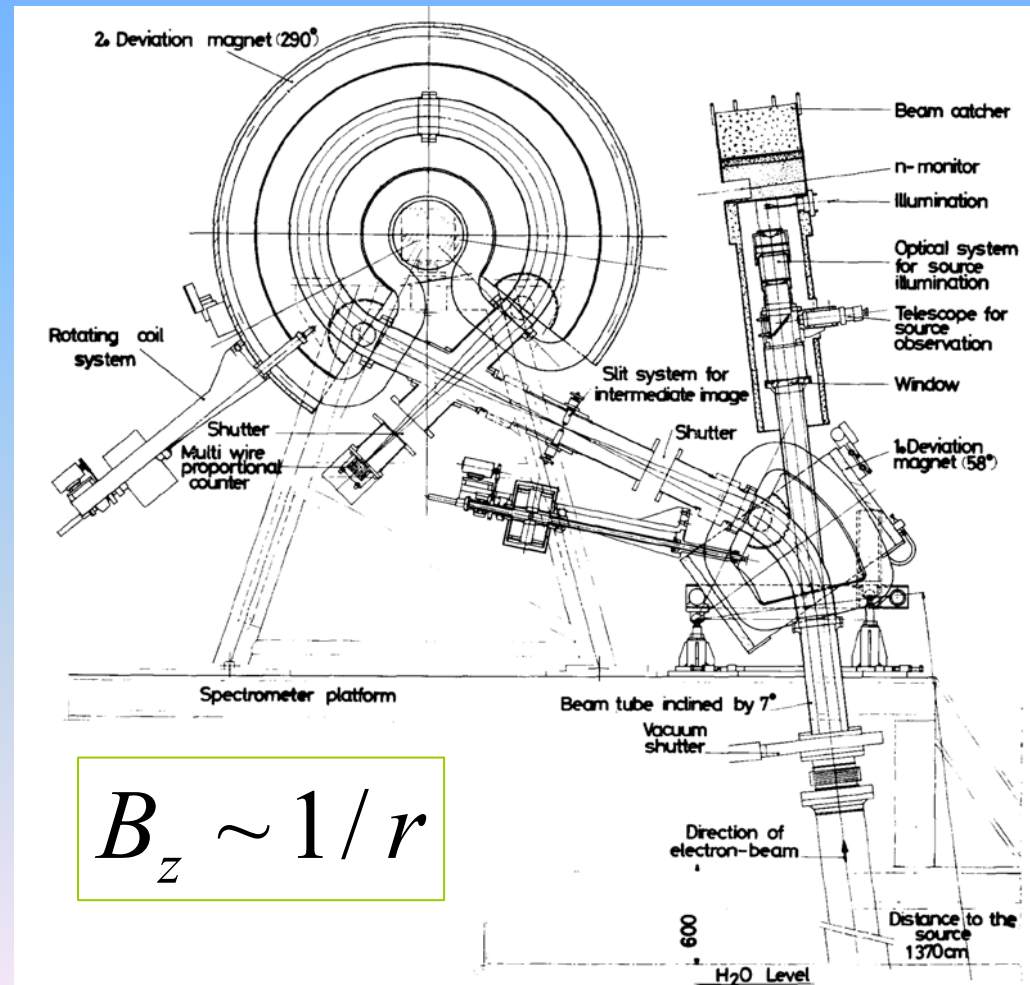
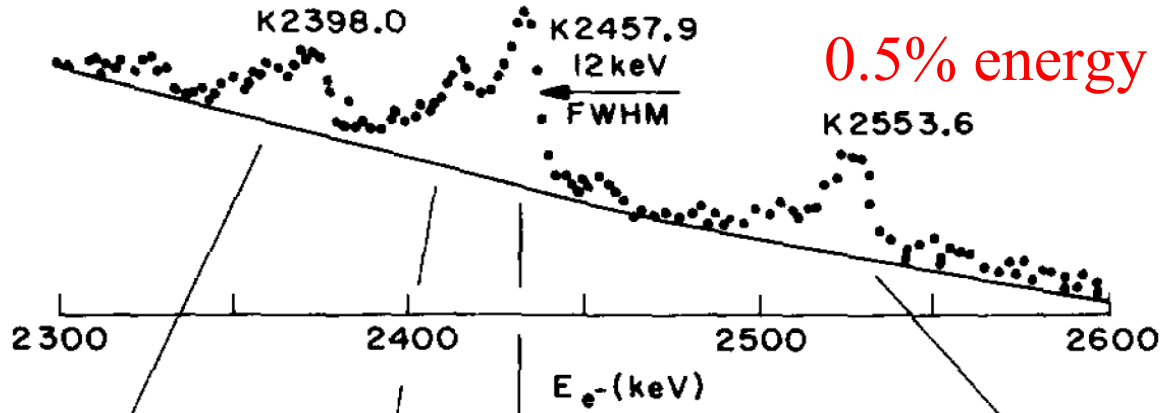
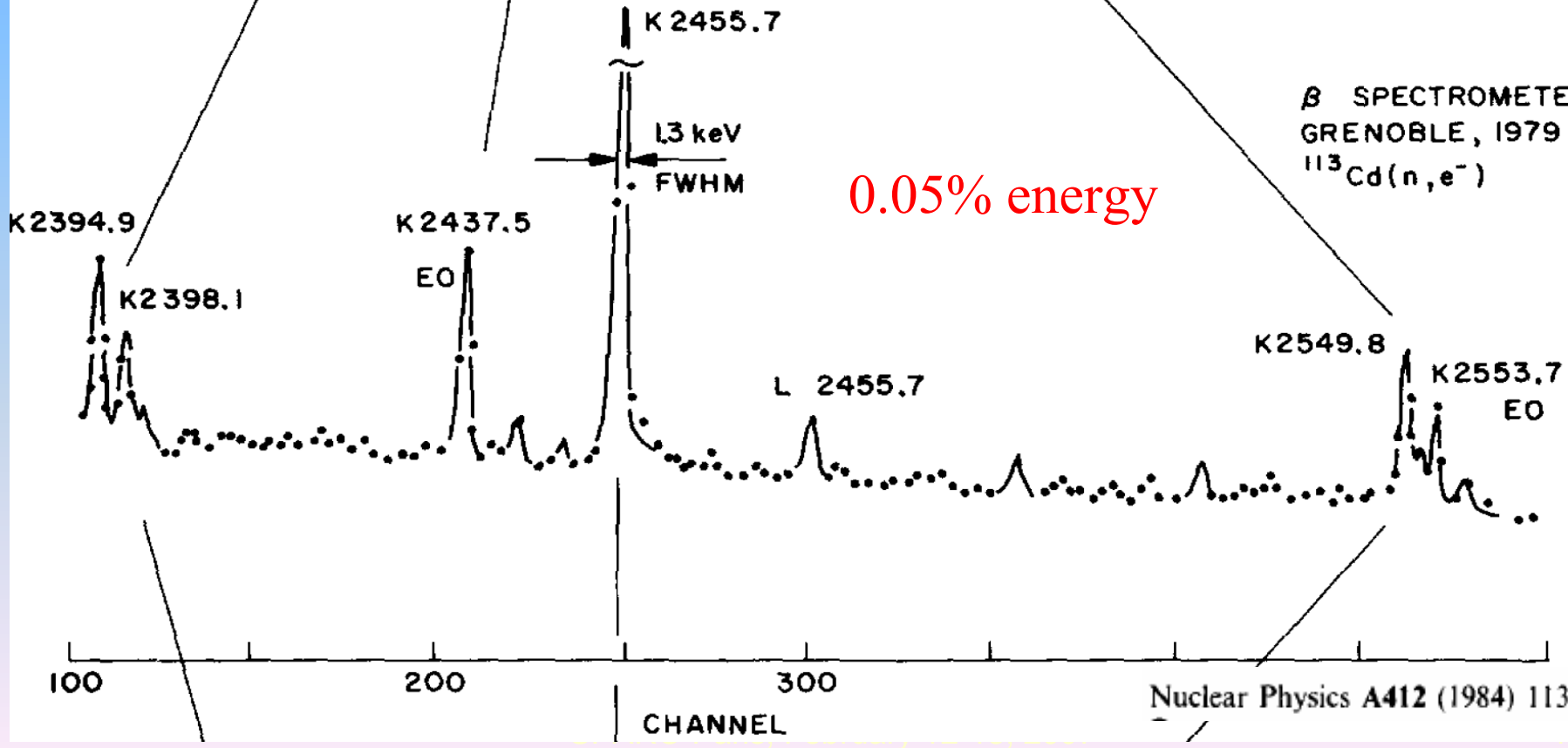


Fig. 2. Schematic view of the spectrometer BILL demonstrating the arrangement of the magnets.

$\beta$  SPECTROMETER  
MUNICH, 1967  
 $^{113}\text{Cd}(n, e^-)$



$\beta$  SPECTROMETER  
GRENOBLE, 1979  
 $^{113}\text{Cd}(n, e^-)$



# Envisioned two-stage magnetic spectrometer

(original ESR proposal Rido Mann et al 1988 – GSI)

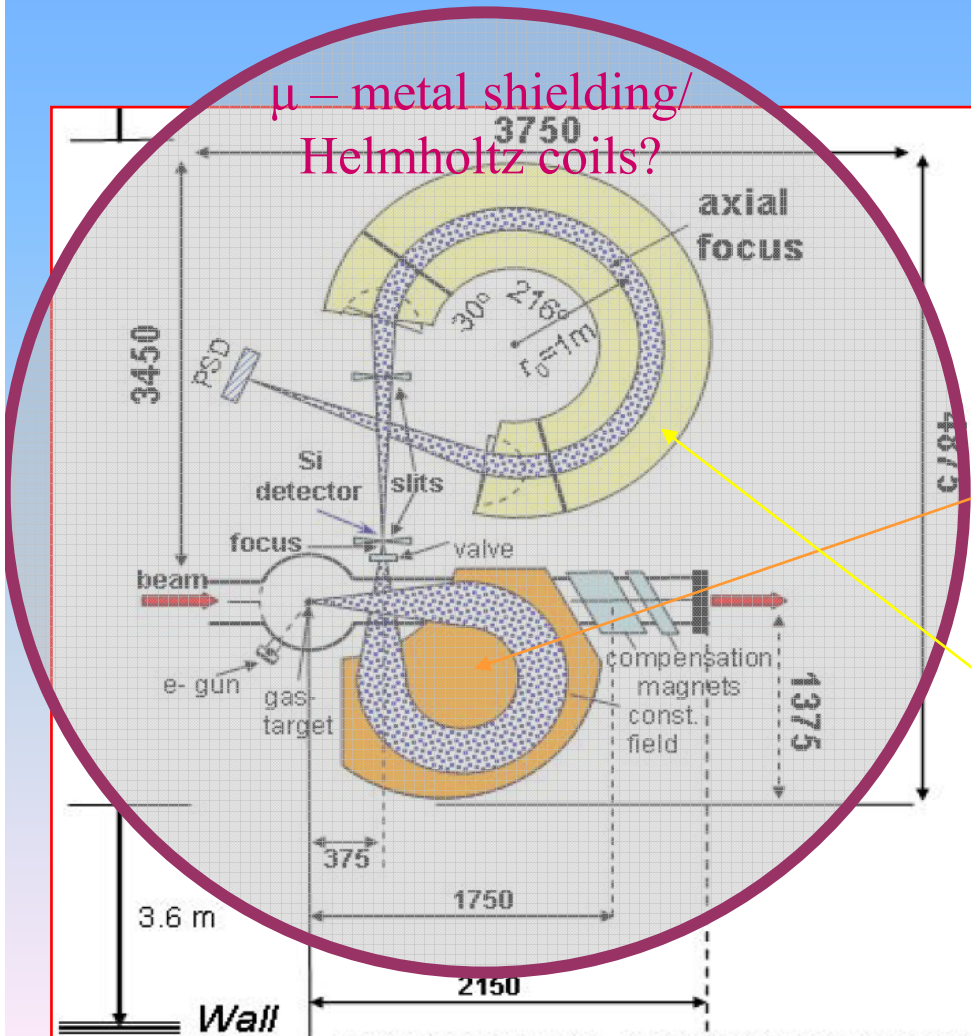
*Use in storage rings: The ultimate ZAPS?*



- Ultimate resolution:  $\delta p/p < 10^{-4}$
- Lab  $e^-$  energy: 3-500 keV

- 1<sup>st</sup> stage (deflector)
- low dispersion, low resolution
  - uniform field dipole
  - target spot size 1mm x 1mm
  - large angular acceptance 5-10<sup>0</sup>

- 2<sup>nd</sup> stage
- high dispersion, high resolution
  - $r^n B_z$ -field (n=1 BILL, n=2 Chalk River)
  - entry slit widths ~0.1-1 mm
  - small angular acceptance 0.1-0.5<sup>0</sup>
  - 2-D PSD



$\Gamma'=100\text{eV}$ ,  $\varepsilon'=80\text{keV}$ ,  $\Delta\theta=0.2^\circ$  (full-acceptance) *Resolution contributions*

Projectile  $\Delta p/p = 5 \times 10^{-5}$ , Spectrometer  $\Delta p/p = 1 \times 10^{-4}$

	4 MeV/u	200 MeV/u	740 MeV/u
Lab width $\Gamma$ (eV)	<b>118.895</b>	<b>258.73</b>	<b>476.01</b>
Kinematic Broadening (eV)	<b>0.0593</b>	<b>2.0875</b>	<b>15.267</b>
Projectile-energy Spread (eV)	<b>1.6310</b>	<b>21.805</b>	<b>58.677</b>
Spectrometer $\Delta E$ (eV)	<b>20.066</b>	<b>63.963</b>	<b>132.92</b>

Important technical issues to take into consideration/resolve:

- **Iron or Air -core magnet design?**
  - Air-core seems better but needs more space!
  - Iron -core: problem of magnetic field uniformity over 1 m radius?  
problem of Remanent magnetization?
- Solid angle considerations - small  $\Delta\theta \sim 0.1^\circ$  (to limit kinematic broadening for line width measurements) will severely limit count rate – PSD necessary
- Good design, optical alignment and slit/baffle controls will be critical
- High quality **non-magnetic** materials to be used in the entire target area
- Need for highest areal density target (liquid H<sub>2</sub>/He)?
- At the 10<sup>-5</sup> precision level
  - **Earth magnetic field annulment** ( $\mu$ -metal shielding/large Helmholtz coils?)
  - **Temperature stability to  $\sim 0.1^\circ$  C**

Anybody up to the challenge?

Come join the SPARC electron spectroscopy group!

[http://www.gsi.de/fair/experiments/sparc/electron-spectrometers\\_e.html](http://www.gsi.de/fair/experiments/sparc/electron-spectrometers_e.html)

## *Bibliography*

- N. Stolterfoht, Phys. Rep. **146** (1987) 315-424.
- T.J.M. Zouros and D.H. Lee, in *Accelerator -Based Atomic Physics Techniques and Applications*, ed. S. Shafroth and J.C. Austin, AIP, Chapter 13 (1997) p. 427-479.
- E.P. Benis et al. Phys. Rev. A **69** (2004) 052718.
- W. Mampe et al., NIM **154** (1978) 127-149.
- R.L. Graham, G.T. Ewan and J.S. Geiger, NIM **9** (1960) 245-286.

**For more information also check my home page:  
<http://www.physics.uoc.gr/~tzouros>**

**XX International Symposium  
on Ion-Atom collisions\* and  
*SPARC topical meeting on  
Electron spectrometry in the NESR***  
**August 1-4, 2007**  
**Agios Nikolaos, Crete, GREECE**

\*a satellite of XXV ICPEAC - Freiburg



$\Gamma'=10\text{eV}$ ,  $\varepsilon'=50\text{keV}$ ,  $\Delta\theta=0.2^\circ$  (full-acceptance)      *Resolution contributions*

Projectile  $\Delta p/p = 5 \times 10^{-5}$ , Spectrometer  $\Delta p/p = 1 \times 10^{-4}$

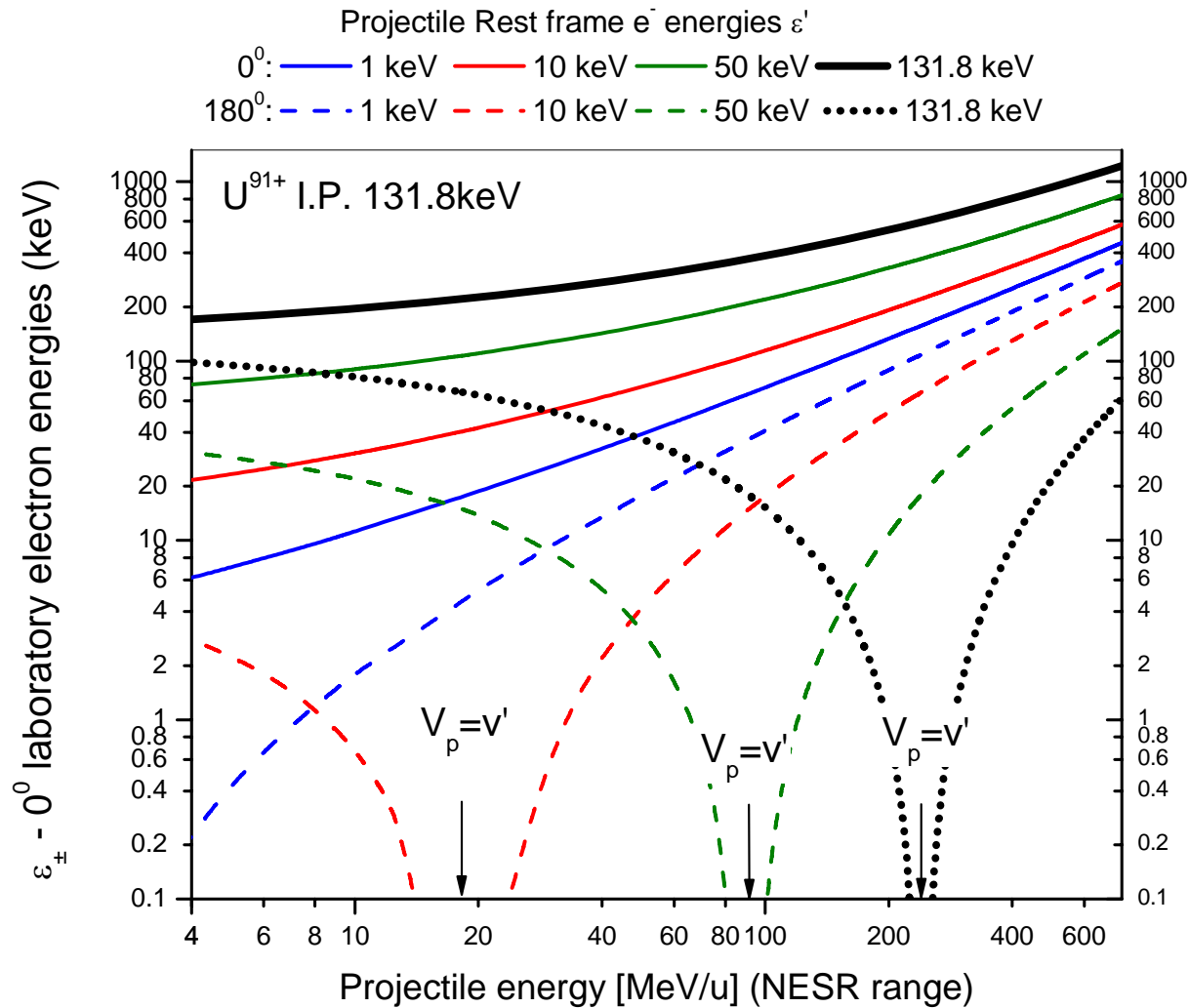
	4 MeV/u	200 MeV/u	740 MeV/u
Lab width $\Gamma$ (eV)	<b>12.291</b>	<b>28.857</b>	<b>54.048</b>
Kinematic Broadening (eV)	<b>0.04942</b>	<b>2.0248</b>	<b>15.347</b>
Projectile-energy Spread (eV)	<b>1.3143</b>	<b>18.963</b>	<b>51.949</b>
Spectrometer $\Delta E$ (eV)	<b>13.844</b>	<b>53.065</b>	<b>115.84</b>

$\Gamma'=10\text{eV}$ ,  $\varepsilon'=1\text{keV}$ ,  $\Delta\theta=0.2^\circ$  (full-acceptance)      *Resolution contributions*

Projectile  $\Delta p/p = 5 \times 10^{-5}$ , Spectrometer  $\Delta p/p = 1 \times 10^{-4}$

	4 MeV/u	200 MeV/u	740 MeV/u
Lab width $\Gamma$ (eV)	<b>24.894</b>	<b>122.53</b>	<b>256.45</b>
Kinematic Broadening (eV)	<b>0.0280</b>	<b>5.0441</b>	<b>47.754</b>
Projectile-energy Spread (eV)	<b>0.3678</b>	<b>11.124</b>	<b>34.056</b>
Spectrometer $\Delta E$ (eV)	<b>1.2258</b>	<b>23.851</b>	<b>69.620</b>

# Kinematics: energy shifting and doubling II

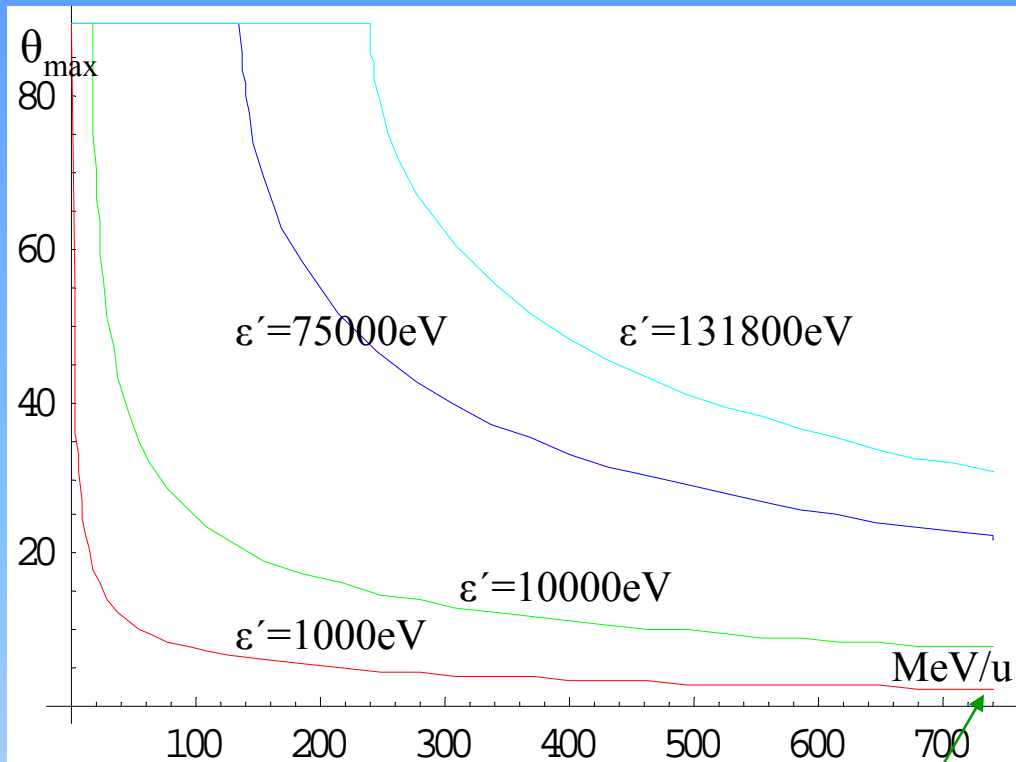


For He-like Uranium  
 K-Auger series energies  
 (75 -130 keV rest) frame  
 Range in Lab from about  
 100 - 1000 keV (+)  
 40 - 0 keV (-)  
 For projectile energies  
 4 -740 MeV/u

## Angular compression – “beaming”

For  $V_p / v' > 1$  ( $\beta_p / \beta' > 1$ ) there is  
a maximum lab observation angle  $\theta_{\max}$  :

$$\theta_{\max} = \text{Arc sin} \left[ \frac{\beta' \gamma'}{\beta_p \gamma_p} \right]$$

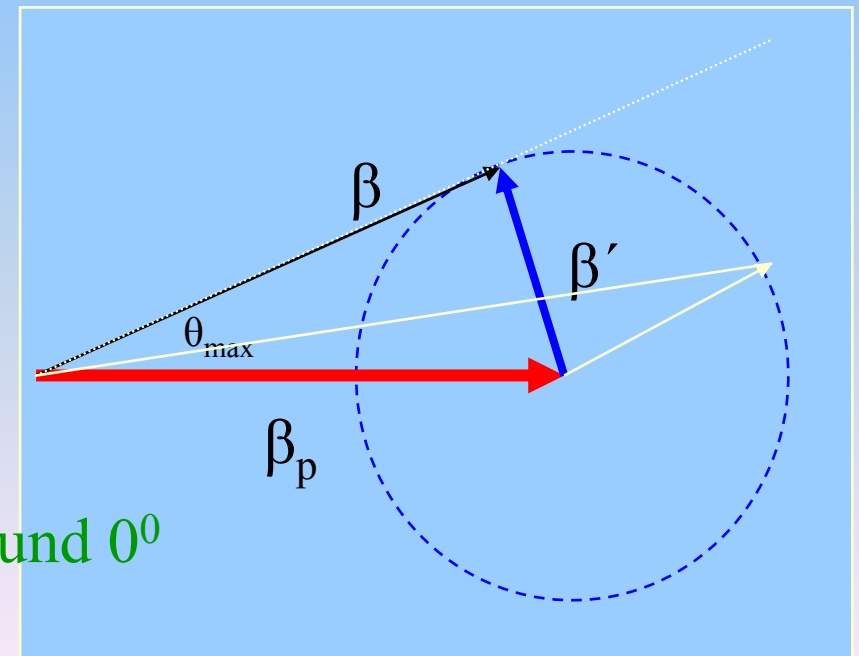


At 740 MeV/u we have:

For  $\epsilon' = 1000 \text{ eV}$ ,  $\theta_{\max} = 2.407^\circ$

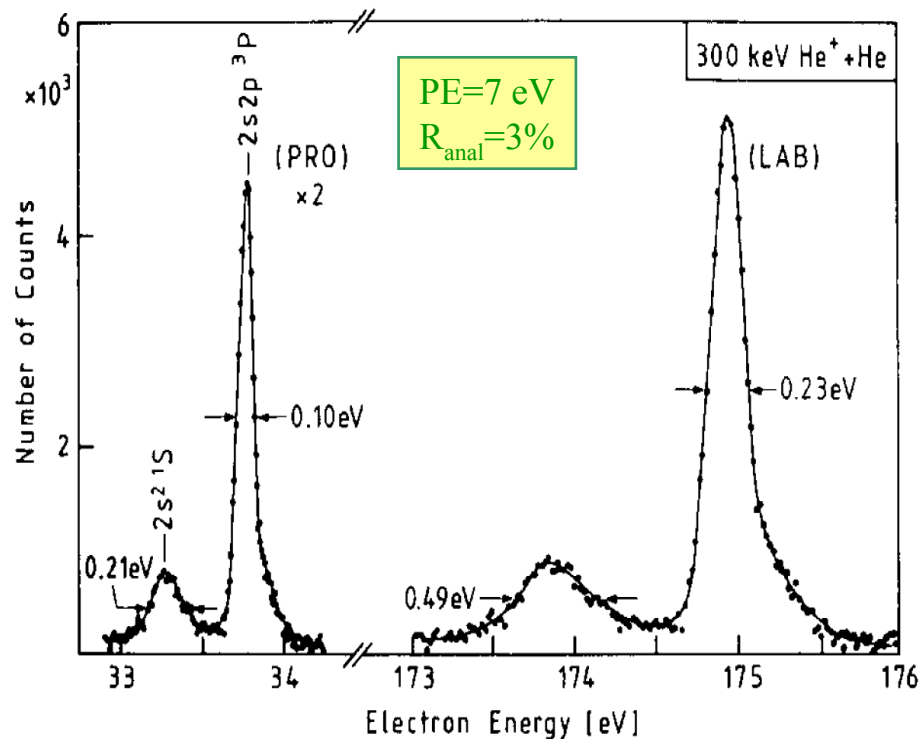
$\epsilon' = 10 \text{ eV}$ ,  $\theta_{\max} = 0.24^\circ$

Strong beaming for small electron energies  
Practically **total cross section** measured around  $0^\circ$   
All differential information averaged out



## *Kinematics: Line stretching and enhancing*

Natural Line widths  $\Gamma'$  (rest frame) are  
Changed to widths  $\Gamma_{\pm}$  in Lab frame



$0^0$  Kinematic Line Stretching or Compressing :

(- only for  $\beta_p > \beta'$ )

$$\Gamma_{p\pm} = |\Gamma_{p'}| \quad (\text{classical - no stretching})$$

$$\Gamma_{p\pm} = \gamma_p |\beta_p \beta' \pm 1| \Gamma_{p'} \quad (\text{relativistic - mild stretching})$$

Momentum Kinematic Line Enhancement

$$\left. \frac{d^2\sigma_{\pm}}{dp d\Omega} \right|_{\theta=0^0} = \frac{d^2\sigma'}{dp' d\Omega'} \quad (\text{classical - No!})$$

$$\left. \frac{d^2\sigma_{\pm}}{dp d\Omega} \right|_{\theta=0^0} = \frac{\gamma_{\pm}}{\gamma'} \frac{d^2\sigma'}{dp' d\Omega'} = \gamma_p (1 + \beta_p \beta') \frac{d^2\sigma'}{dp' d\Omega'} \quad (\text{relativistic})$$

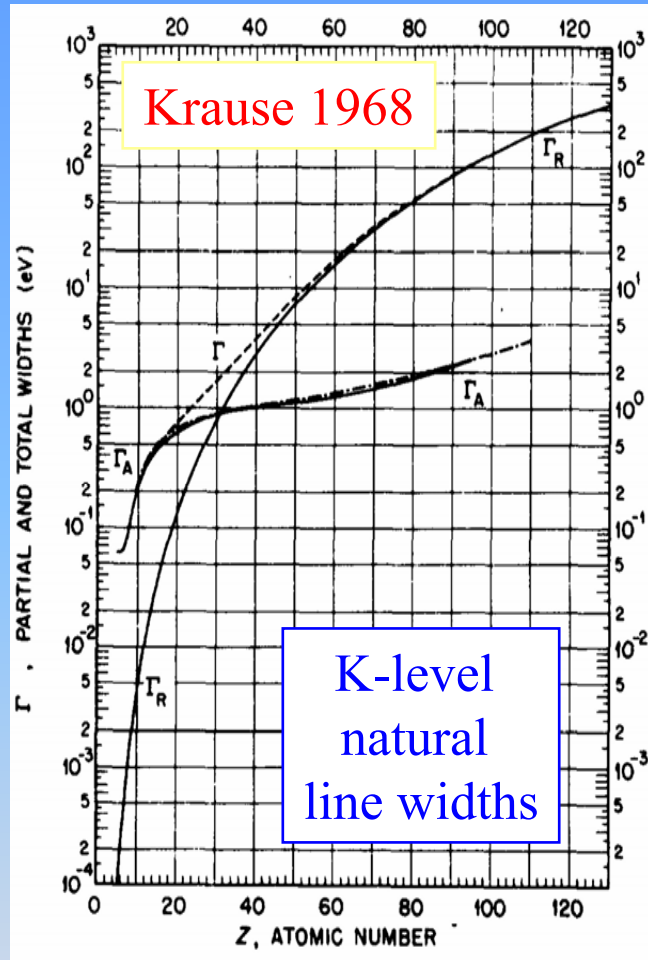
Energy Kinematic Line Enhancement

$$\left. \frac{d^2\sigma_{\pm}}{d\varepsilon d\Omega} \right|_{\theta=0^0} = \frac{V_p \pm v'}{v'} \frac{d^2\sigma'}{d\varepsilon' d\Omega'} \quad (\text{classical - yes!})$$

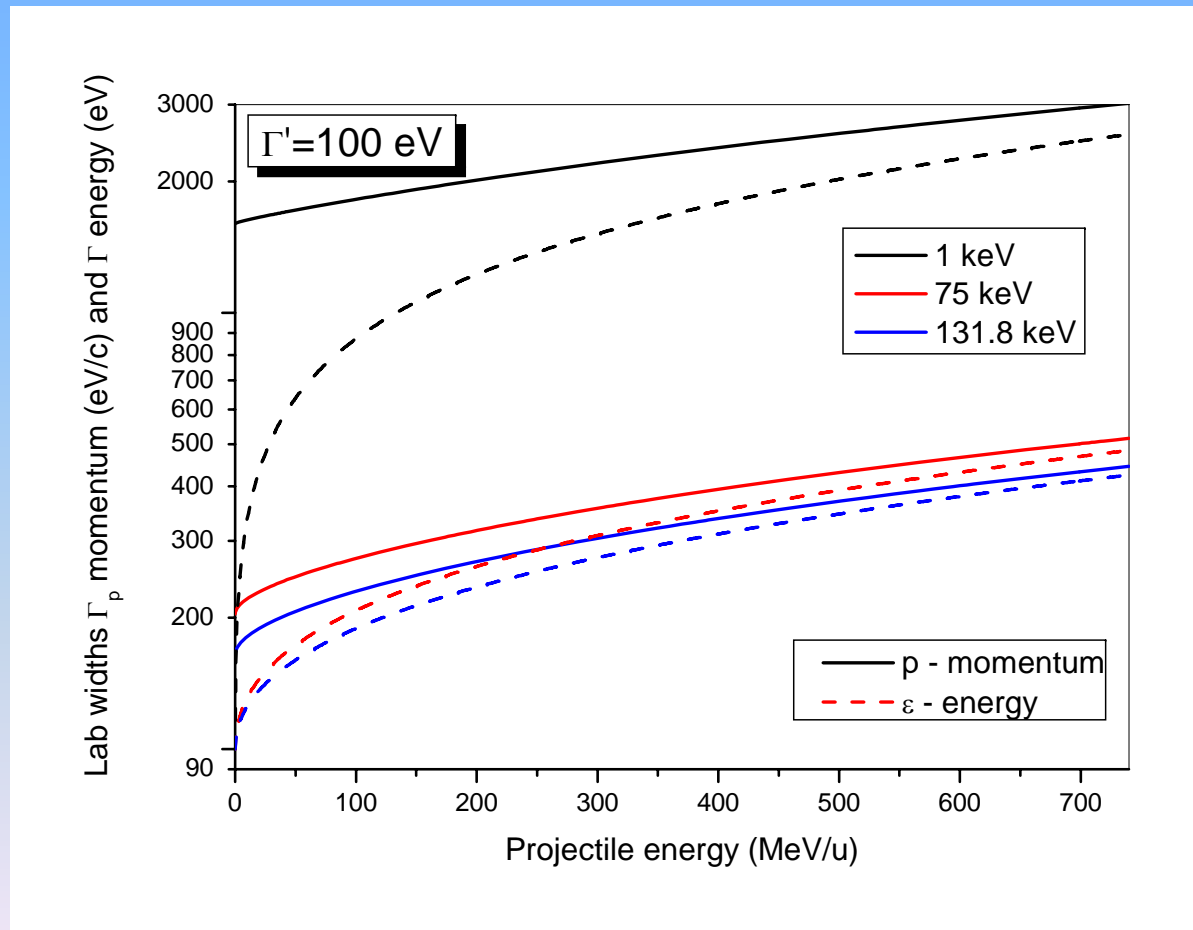
$$\left. \frac{d^2\sigma_{\pm}}{d\varepsilon d\Omega} \right|_{\theta=0^0} = \frac{p_{\pm}}{p'} \frac{d^2\sigma'}{d\varepsilon' d\Omega'} = \gamma_p \frac{(\beta_p \pm \beta')}{\beta'} \frac{d^2\sigma'}{d\varepsilon' d\Omega'} \quad (\text{relativistic})$$

Mild enhancement and stretching  
in momentum analysis!!

# Kinematic Widths and Enhancement factors



Momentum widths are stretched only weakly  
While energy widths a lot!

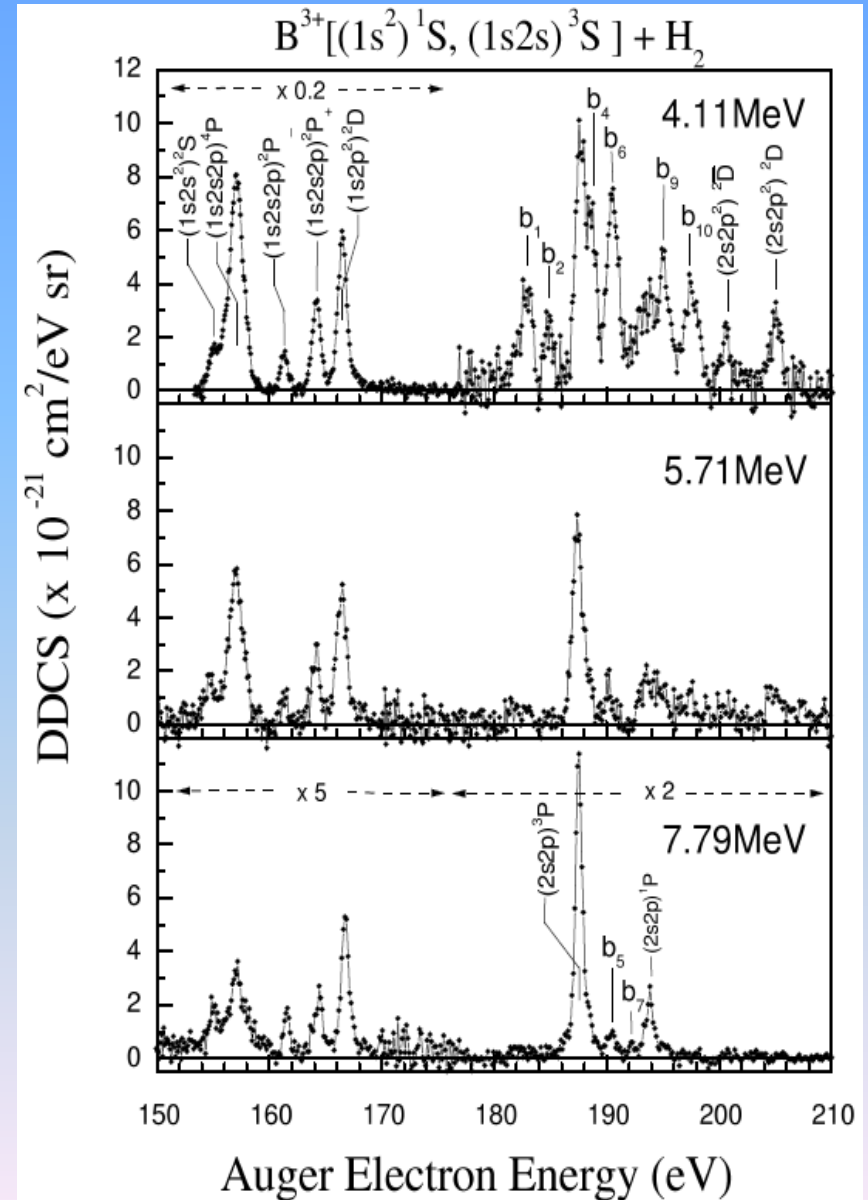
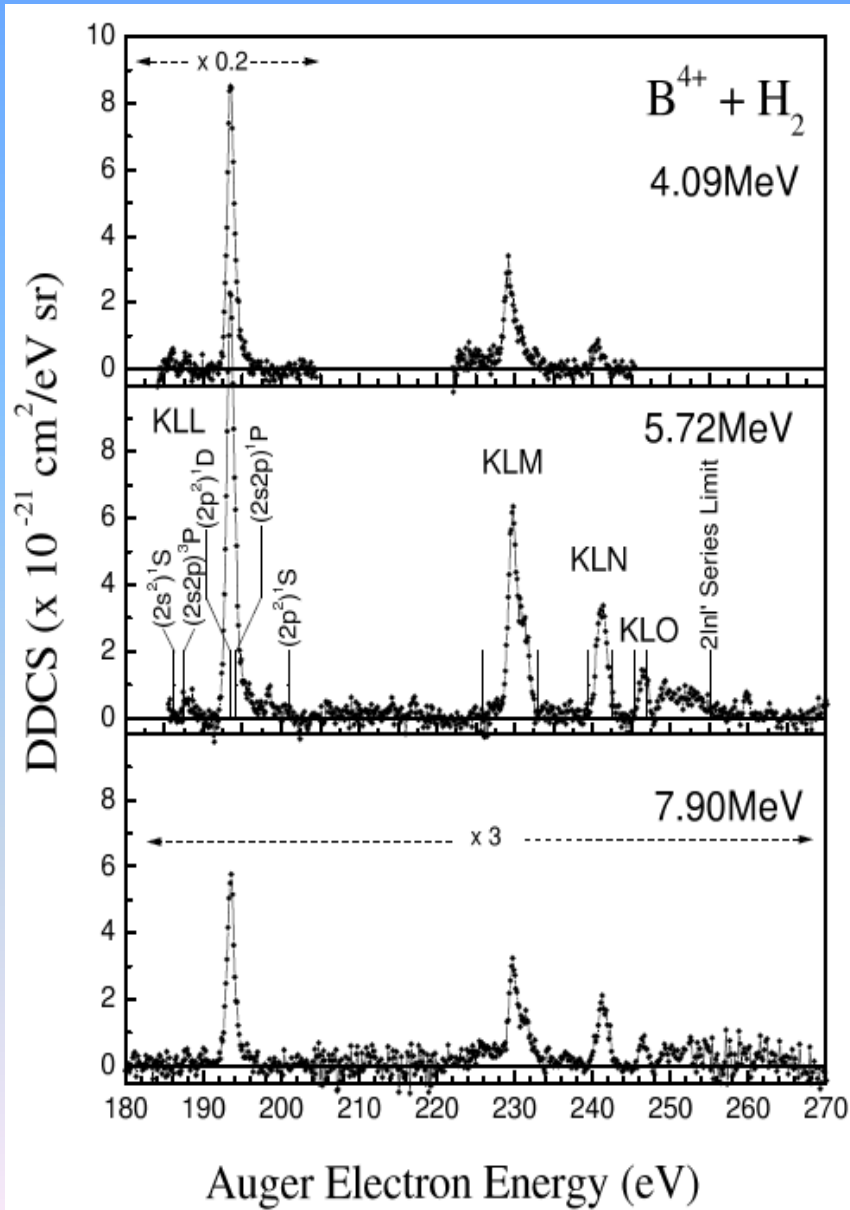


For **atoms**

K line-widths dominated by  
Radiative widths at high Z  
And Auger widths at low Z

**What about Highly Charged Ions?**

# Zero-degree Auger Projectile electron spectra



**Elastic resonant and nonresonant differential scattering of quasifree electrons  
from  $B^{4+}(1s)$  and  $B^{3+}(1s^2)$  ions**

E. P. Benis,<sup>1,\*</sup> T. J. M. Zouros,<sup>2,3,†</sup> T. W. Gorczyca,<sup>4</sup> A. D. González,<sup>5</sup> and P. Richard<sup>1</sup>

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